DISEÑO CINEMATICO DE MAQUINARIA 1984

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echa	Temario	Horario	Profesor
unio 25	Definiciones y Conceptos Básicos	9 a 13:30 h	. Dr.Jorge Angeles Alvaře
	Clasificación de Mecanismos	15 a 17:30 h.	
unio 26	Análisis de Mecanismos de Pares inferiores	9 a 13:30 h	Dr. Jorge Angeles Alvar
<i>v</i>	Introducción a la Robótica	15 a 17:30 h	
unio 27	Síntesis de Mecanísmos de Pares Inferiores	9 al3:30 h	Dr. Jorge Angeles Alvar
	Optimación de Mecanismos	15 a 17:30 h	Dr. Jorge Angeles Alvare
unio 28	Mecanismos de levas	9 a 13:30 h	Dr. Carlos López Cajún
		15 a 17:30 h	
unio 29	Mecanismos de engranes	9 a 13:30 h	M. en I. Angel A. Rojas
	Ejemplos de aplicación	15 a 17:30 h 🏅	5a1gado

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EVALUACION DEL PERSONAL DOCENTE

MANTENIMIENTO DEL INTERES. (COMUNICACION CON EL USO DE LOS ASISTENTES, AMENIDAD, FACILIDAD DE EXPRESION). AUDIO VISUALES DOMINIO DEL TEMA PUNTUALIDAD EFICIENCIA EN CURSO DISENO CINEMATICO DE MAQUINARIA AYUDAS FECHA: Del 25 al 29 de junio de 1984. CONFERENCISTA Ļ, þr. Jorge Angeles Alvarez 2 br. Carlos López Cajún 3. . Salgado Rojas 4 6 6 7 8 9 ESCALA DE EVALUACION : 1 a 10

 (\mathbf{I})

SU EVALUACION SINCERA NOS AYUDARA A MEJORAR LOS PROGRAMAS POSTERIORES QUE DISENAREMOS PARA USTED.	ORGANIZACION Y DESARROLLO DEL TEMA	GRADO DE PROFUNDIDAD Logrado en el tema	GRADO DE ACTUALIZACION Lograd o en el tema	UTILIDAD PRACTICA DEL Tema	
Definiciones y conceptos básicos		:	· · · · · · · · · · · · · · · · · · ·		
Clasificación de mecanismos		· · ·			
Análisis de mecanismos de pares inf			 _		
Introducción a la robótica			·		
Sintesis de mecanismos de pares int		· ·			
Optimización de mecanismos					
Mecnismos de levas					
Elemplos de aplicación		-			

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EVALUACION DEL CURSO

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	CONCEPTO	EVALUACION
1.	APLICACION INMEDIATA DE LOS CONCEPTOS EXPUESTOS	
2.	CLARIDAD CON QUE SE EXPUSIERON LOS TEMAS	
з.	GRADO DE ACTUALIZACION LOGRADO CON EL CURSO	
4.	CUMPLIMIENTO DE LOS OBJETIVOS DEL CURSO	
5	CONTINUIDAD EN LOS TEMAS DEL CURSO	
6,	CALIDAD DE LAS NOTAS DEL CURSO	
7.	GRADO DE MOTIVACION LOGRADO CON ELCURSO	

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ESCALA DE EVALUACION DE I A Ю

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1. ¿Qué le pareció el ambiente en la División de Educación Continua?

MUY AGRADABLE	AGRADABLE	DESAGRADABLE

2. Medio de comunicación por el que se enteró del curso:

PERIODICO EXCELSIOR AMUNCIO TITULADO DI VISION DE EDUCACION CONTINUA	PERIODICO NOVEDADES ANUNCIO TITULADO DI VISION DE EDUCACION CONTINUA	FOLLETO DEL CURSO

CARTEL MENSUAL	RADIO UNIVERSIDAD	COMUNICACION CARTA, TELEFONO, VERBAL, ETC.

REVISTAS TECNICAS	FOLLETO ANUAL	CARTELERA UNAM "LOS UNIVERSITARIOS HOY"	GACETA UNAM

3. Medio de transporte utilizado para venir al Palacio de Minería:

AUTOMOVIL PARTICULAR	METRO	OTRO MEDIO

- 4. ¿Qué cambios harfa usted en el programa para tratar de perfeccionar el curso?
- 5. ¿Recomendaría el curso a otras personas?



- 6. ¿Qué cursos le gustaría que ofreciera la División de Educación Continua?
- 7. La coordinación académica fue:

EXCELENTE	BUENA	REGULAR	MALA

 Si está interesado en tomar algún curso <u>intensivo</u> ¿Cuál es el horario más conveniente para usted?

LUNES A VIERNES PE 9 A 13 H. Y DE 14 A 18 H. (CON COMIDAS)	LUNES A VIERNES DE 17 A 21 H.	LUNES, MIEROOLES Y VIERNES DE 18 A 21 H.	MARTES Y JUEVES DE 18 A 21 H.	

VIERNES DE 17 A 21 H. SABADOS DE 9 A 14 H.	VIERNES DE 17 A 21 H. SABADOS DE '9 A 13 Y DE 14 a 18 H.	OTRO

9. ¿Qué servicios adicionales desearía que tuviese la División de Educación Continua, para los asistentes?

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10. Otras sugerencias:



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DIVISION DE EDUCACION CONTINUA FACULTAD DE INGENIERIA U.N.A.M.

DISERO CINEMATICO DE MAQUINARIA

SOFTWARE PARA EL ANALISIS DIGITAL DE SISTEMAS MECANICOS

> JORGE ANGELES ALVAREZ MANUEL CALLEJAS CASTRO



.UNIO, 1984

"SOFTWARE PARA EL ANALISIS DIGITAL DE SISTEMAS MECANICOS"

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Manuel Callejdo Castro Ayudante de Profesor Departamento de Ingeniería Mecánica Facultad de Ingeniería, UNAM.

Abatract

Subprograms for digital computer are presented that allow, by their coupling within a main program, the analysis of multiple-loop linkages. Each subprograms analyzes kinematically a single-loop-single-degree-of-freedom linkage. The subprogram developed thus far analyze plans ERRR, RPRR and KRRP linkages. As an example, the analysis of the driving linkage of a machanical shovel is included.

Resumen

Se presentan subprogramas de computadora digital que permiten, mediante su acoplamiento en un programa principal, el análisis de mecanimmos de relabones rígidos de malla múltiple. Cada subprograma analiza cinemáticamente un mecanismo plano du grado de libertad simple y de una solo malla. Los subprogramas hasta ahora desarrollados apalizan mecanismos RRRR, RPRR y RRP. Como ejemplo, se incluye el análisia del mecanimmo accionador de una pala mecánica.

Introducción

El análisis cinemático de sistemas mecánicos mediante computadore digital cobra importan cia en el proceso de diseño o de rediseño de tales sistemas, pues a través de este gnálisis es posible evaluar su operación sin necesidad de construirlo, En efecto, mediante este análipie puedo deperminarse la evolución de variables tales como máxima aceleración (lineal o angular) de portes críticas, la ventaja mecánica de la transmission, o bien detectorse situsciones adversas tales como interforoncias. Adicionalmente, cabe señalar que el enflísis mediante computadora digital permite examinat la operación de sistemas mecánicos con todo el detalle deses do y con tanta precisión como sea nuceario, sin tener que construir costosos prototipos. con el ahorro consecuente en economía y en cuan to a tiempo.

Los sixtemas mecánicos a los que es aplica ble el "software" aquí presentado consisten de acoplamientos de mecanismos planos de grado de libertad simple y de une sola malla.

Existen programas tales como el KAPKA o el IMP(I) que sirven para analizar cinemática y dinámicamente sistemas mecánicos de mallas miltiples y de múltiple grado de libertad. Inclusive, en (2) se consigna ampliamente el software disponible para el propósito mencionado

El objeto que se persigue al desarrollar el softwars aquí presentado es múltiple. Por un lado, orientarlo hacia la aplicación de técnicas interactivas de graficación. Por otro lado, desarrollarlo en forma modular, lo que permiti rá su utilización en forma más eficiente, pues así el diseñador puede conter con una programo reca de la que puede seleccionar solu los subprogramas que él requiera, sin tener que utili zar memoria de computadora que tendría ociusa, en caso de recurrir a un programa de propósito general. Finalmente, el desarrollo de un soltware propio es desemble, pues esto contribuye a la creación de una tennología propia.

Descripción del soltware

Los subprogramas que se describen a contipuación, RESCII (RESpuesta Cinemática), RESCI2 y RESCI3, sirven para el snálisis de mecanismos de los tipos RRRR, RPRR y RRRP, respectivamente. Estos mecanismos se muestran en las Figs. I-3. Cada uno de esos subprogramas se describa a continuación.

RESCII. Algoritmo de cálculo

Se mupone que se conoce perfectimente los valores a_1, \ldots, a_5 y a de la Fig 1, esí como la excitación del mecanismo, $\Psi = \phi(t)$; así, se cono cen también $\psi(t)$, $\psi(t)$ y derivadas de orden superior de esta función. El análisis comprende la obtención de la tespucata cónomítica del mecanismo, que incluye $\psi(t)$, ψ

El ángulo ; se obtiene de la ecusción de Freudenstein (3):

 $k_1 - k_2 \cos \phi + k_2 \cos \phi + \cos (\phi - \psi) = 0 \tag{1}$

donde

$$k_{1} = \frac{a_{3}^{2} - a_{1}^{2} - a_{2}^{2} - a_{4}^{2}}{2a_{2}a_{4}^{2}}, \quad k_{2} = \frac{a_{1}}{a_{2}}, \quad k_{3} = \frac{a_{1}}{a_{4}}$$
(2)

Escribiendo 1a ec. (1) en la forma

y Sustituyendo las siguientes identidades trigo • nométricas:

$$\operatorname{sen}_{\varphi} = \frac{2t \tan(\varphi/2)}{1 + \tan^2(\varphi/2)}, \quad \operatorname{cus}_{\varphi} = \frac{1 - \tan^2(\varphi/2)}{1 + \tan^2(\varphi/2)}$$
(3)

1e ec. (1') se transforms en

donde

$$A=k_1+k_2-(1-k_3)\cos\psi \qquad (5a)$$

Așî la ec. (4) es cuadrătica en tan(¢/2), y su selución es, sencillamențe,

$$ran(\phi_{1,2}/2) = \frac{-B^{\pm}\sqrt{B^{2}-4AC}}{2A}$$
 (6a)

de dande

$$s_{3,2} = 2 \tan^{-1} \left(\frac{-B \pm \sqrt{B^2 - 4AC}}{2A} \right)$$
 (6b)

Le solución (ba) e la cuadrática (4) es contrecta desde el punto de vista algebraico. Sin embargo, numéritamente puede causar dificultades catastróficas, como se apunte en $\{4\}$. En efecto, si $B^2 >> 4AC$, la primera raíz (tomando el signo positivo en la solución (ba)) se anula a causa del error de redondeo. Para evitar esta situación, se calcula en este caso la primera raíz como

$$\tan(\phi_1/2) = \frac{B + (B^2 - 4AC \operatorname{sgn}(B))}{2\lambda}$$
(7a)

y la segunda, como

$$\tan(\phi_2/2) = \frac{C}{A \tan(\phi_1/2)}$$
 (7b)

Las due raices anteriores corresponden a las posiciones conjugadas del mecanismo. En un punto muerto del colabón de entrada, claramente las dos posiciones conjugadas se reducen a una sola, lo cual succede cuando se anula el radical. Si este radical no se anula para ningún valor de \$, al aslabón de entrada no tiene ninguna posición de punto muerto, esto es, constituye una manivela. En estas condiciones, conviene en pri mer lugar determinar ei este eslabón es del tipo balancín o del Lipo manivala. Esto se puede determinar del criterio de Grasholf (5); pero como en este cuso interesa conocor, siempre que se trate de un balancin, sus configuraciones extreman, se procede de la siguiente forma: Ya que el mocanismo su encuentra en una configu ración extrama al anularse el redical, se debe determinar para qué valores de 9 surede esto, es decir, es recesario despejar \$ de la ecuación

$$\tau(\psi) = B^2 - 4AC = 0 \tag{8}$$

Al sustituit los valores (5a-t) en la ec. (8) ae tiene

$$\cos^2\psi + 2b\cos\psi + c = 0$$
 (9)
double

$$b = \frac{k_2^{+k_1k_2}}{k_3^2}$$
(10a)

$$c = \frac{k_1^2 - k_2^2 - 1}{k_1^2}$$
(10b)

Las raíces de la uc. (9) son, entoures,

$$\cos \psi_{1,2} = b \pm \sqrt{b^2 - c}$$
 (11)

Haciendo la sustitución de b y c en la ecuación (11) en términos de k_1, k_2 y k_1 y en seguida sustituyendo éstas por a_1, a_2, a_3 y a_4 se tiene 2 a 2

$$\cos \psi_{1,2} = \frac{a_1^2 + a_2^2 - (a_1 + a_2)^2}{2a_1 a_2}$$
(12)

que, como se ve, da lugar a rafcas reales. Se tiene, entonces, las siguientes situaciones posibles

- i) El valur absoluto de ambas ruíces es menor que l.
- Sólo una raíz tiene valor absoluto senor que 1.
- iii) Ambas raices rienen velor absoluto mayor que 1.

En el caso i) el estabón de entrada es del tipo balancín, estando dadas sus 2 configuraciones extremas por las raíces (12). En el caso i)) este estabón también es balancín y, supaniendo que la primera raíz tenga valor absoluto menor que i, las configuraciones extremas están dadas por esta raíz y son simétricas, esto es, $\frac{4}{7} = -\frac{4}{1}$. En el caso iii), el estabón ve del tipo manivela. Una vez determinadas las configuraciones extremas, se definen $\frac{1}{7}$ máx

(para emlabones del tipo bulancin) de acuerdo etn los valores que adquiera :"(ϕ) en esas con figuraciones, según el criterio conocido de la segunda derivada, esto es, ϕ es minima donde $\psi'(\phi)$ se anula y $\psi''(\phi) > 0$; ψ es máxim, donde $\psi'(\phi)$ se anula y $\psi''(\phi) < 0$.

Los variables 9(t) y 3(t) se obtienen de las fórmulas

$$\dot{\phi}(z) = \frac{d\phi}{d\psi} \dot{\phi} \qquad (13)$$

 $\ddot{\bullet}(\mathbf{t}) = \frac{d^2 \mathbf{e}^{+2}}{d \psi^2} + \frac{d z}{d \psi} \ddot{\psi} \qquad (14)$

donde las derivadas con respecto a y se calculan de la ecuación de Freudenstrin (1) que define, bajo lus condiciones de existencia del caso (6), a e como función implícita de t. En

(Sb)

$$f(\phi,\phi) = 0$$
 (15)

de donde

$$\frac{d\phi}{d\psi} = -\frac{\partial f/\partial \psi}{\partial f/\partial \phi} = -\frac{N}{D}$$
(16)

$$\frac{d^2\phi}{d\phi^2} = -\frac{1}{D} \left(\frac{\partial N}{\partial \phi} + \frac{d\phi}{\partial \phi} - \frac{\partial D}{\partial \phi} \right) \frac{d\phi}{d\phi} - \frac{\partial}{\partial \phi} \frac{\partial}{\partial \phi} = \frac{\partial}{\partial \phi} \frac{\partial}{\partial \phi}$$
(17)

Las variables $\theta(t)$, $\dot{\theta}(t) \neq \theta(t)$ so calculan de una ecuación semejante a la de Freudenstein (η) :

$$L_1 + L_2 \cos \theta + L_3 \cos \theta - \cos (\theta - \theta) = 0$$
 (18)

El ángulo 8(#) se calcula en forma análoga a como se calculó e(#) de las ecs. (7a y b), obteniëndose asī dos posiciones conjugadas. Las variables 0(t) y 0(t) se obtienen, análoga mente, como

$$\hat{\theta} = \frac{d\Phi}{d\psi} \hat{\psi}$$
(19)

y

Y

$$\ddot{\theta} = \frac{d^2\theta}{d\psi^2} \dot{\psi}^2 + \frac{d\theta}{d\psi} \ddot{\psi}.$$
 (20)

donde las derivadas parcisles con respecto a # se obtienen de la ec. (12) que define anâlogamente a 8 como función implícite de #. En efec to, escribiendo la ec. (18) como

se obtuene

$$\frac{d\theta}{d\psi} = -\frac{\partial h/\partial \psi}{\partial h/\partial \theta} = -\frac{N}{D}$$
(22)

$$\frac{d^2\theta}{d\phi^2} = \frac{1}{D} \left(\frac{\partial N^4}{\partial \phi} + \frac{d\theta}{d\phi} \frac{\partial D^4}{\partial \theta} \right) = \frac{1}{\partial \phi} \frac{N^4}{D^4}$$
(23)

Las variables x(t), y(t) as calculan de le geometrís de la Fig 1, sencillaments como

calculándose sus derivadas por derivación directa de estas ecuaciones, Las fórmulas corres pondientes son fácilmente obtenibles y por la falta de espacio no se incluyen. Finalmente, el ángulo de transmisión se calcula sencilleecote de -

RESCI2. Algoritmo de câlculo

Se supone que un conocan perfectamente los valores $a_1 \neq a_2$, as como la excitación $a(t) \neq sus derivadas. Se desea determinar <math>\theta(t)$, $\dot{\theta}(t)$, $\theta(t)$, $\phi(t)$, $\dot{\phi}(t)$ y $\ddot{\phi}(t)$. De la Pig 2 se obtim-0 e $a^2 = a_1^2 + a_2^2 + 2a_1a_2come$ (26)

de donde

В

$$\frac{de \ donde}{\phi} = \cos^{-1} \left[\frac{a^2 - a_1^2 - a_2^2}{2a_1 a_2} \right]$$
adicionalmente.
(27)

$$\theta = \tan^{-1} \left(\frac{a_2 \sin \phi}{a_2 \cos \phi + a_1} \right)$$
(28)

Las derivadas se calculan derivando directamente las relaciones:

As1, we obtione

$$\dot{\phi} = \frac{2\dot{a}}{a_2 \operatorname{men}(\phi - \theta)}$$
(30)

$$\partial = \frac{-\dot{\theta}}{g \tan(\phi - \theta)}$$
(31)

$$= \frac{-(s-e_2\cos(\phi-\theta))}{e^2 \sigma \tan(\phi-\theta) \sin^2(\phi-\theta)} = \frac{1}{e^2 \sin(\phi-\theta)} = (32)$$

$$\tilde{\theta} = \frac{a_2 \cos(\phi-\theta) \left(1 + \cos^2(\phi-\theta)\right) - a_2}{a_2 - \frac{2}{a + \sin^2(\phi-\theta)}} = \frac{B}{a \tan(\phi-\theta)}$$
(33)

RESCIJ. Algoritmo de câlculo

Se supone que se conocen perfectamente las dimensiones a1, a2, a3 y a4, así como el ingulo a del mecanismo RRRP de la Fig 3, además de la excitación e(t) y sus derivadas. Se deses calcular s(t), s(t), T(t), \$(t), \$(t), \$(t), $x(t), \dot{x}(t), \dot{x}(t), \dot{y}(t), \dot{y}(t), \dot{y}(t)$. De la geometrís de la Fig 3,

De la ec. (34b),

$$= 4 \times n^{-1} \left[\frac{-4 \cdot 2^{4} \times n + 1}{3 \cdot 3} \right]$$
 (35)

Si $a_1 + a_2 \ge a_3$, la monivela oscila de ψ_1 a ♦.. donde

$$e_1 = \sin^{-1}\left(\frac{a_3 - a_1}{a_2}\right)$$
 (364)

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$$\psi_2 = -(\psi_1 + 160^*)$$
 (36b)

Derivando (34b) con respecto al tiempo,

Durivando (34a) con respecto al ticapo y susrituyendo (37a) en la expresión así obtenid∎

$$a = - (a \pm en\psi + a \pm en\phi) \frac{a_2 \cos \phi}{a_3 \cos \phi}$$
(37b)

Derivando las expresiones (374) y (376) con respecto al tiempo se obtiene

$$\frac{a_{2}^{2} \cos^{4}\psi \exp(-a_{2}a_{3}^{2} \sin\psi \cos^{4}\phi)}{a_{3}\cos^{4}\phi} \frac{1}{\psi^{2}} \frac{a_{2}\cos\psi}{a_{3}\cos\phi} \frac{1}{\psi} (38a)$$

$$\frac{a_{2}^{2} \sin\psi \exp(a_{3}^{2}\cos^{2}\phi) - a_{2}^{2}a_{3}(\cos^{2}\psi) - a_{2}^{2}a_{3}(\cos^{2}\psi)$$

Las variables x(t), y(t) se calculan de la geometría de la Fig 3, sencillamente como

$$y = *_{x} sen(y + a_{x} sen(y + a))$$
(395)

Calculándose sus derivadas por derivación directa de estas acuaciones. Las fórmulas cofrespondientes son facilmente obtenibles y por falta de espacio no se incluyen.

• Ejumplo:

² Determine $\phi(t)$, $\dot{\phi}(t)$, $\ddot{\phi}(t)$ y $\mu(t)$ del macaniumo accionador de la pala macanica de la Fig 4. El modelo cinemático y la definición de las variables anteriores se indican en la Fig 5. La entrada del macanismo es a(t)=3+0.5 uen² st. Se acoplaron las subrutines RESCI1 y RESC12 ²⁷ en un programa principal y se obtuviron las cur vas de la Fig 6.









Fig 5, Modelo cinemático del ______ mistema macánico de la Fig 4.

Referenciae

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- Angeles J., <u>Análisis y Síntesis Cinemáticos</u> <u>de Sistemas Mecánicos</u>, Limusa, S.A., C. de México, 1978, pp. 53-60.



Fig 6. Respuesta cinemática del elstema mecánico de la Fig 4.



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DIVISION DE EDUCACION CONTINUA FACULTAD DE INGENIERIA U.N.A.M.

DISEND CINEMATICO DE MAQUINARIA

AN ALGEBRAIC FORMULATION OF GRASHOF'S MOBILITY CRITERIA WITH APPLICATION TO LINKAGE OPTIMIZATION USING GRADIENT-DEPENDENT METHODS

> JORGE ANGELES MANUEL CALLEJAS

JUNIO, 1984

AN ALGEBRAIC FORMULATION OF GRASHOF'S MOBILITY CRITERIA WITH APPLICATION TO LINKAGE OPTIMIZATION USING GRADIENT-DEPENDENT METHODS

Jorge Angeles, Professor¹ (ASME Member 79) Manuel Callejas, Research Assistant¹

Abstract

Sets of constraints on analytic functions of linkage parameters, equivalent to those of Grashof's mobility criteria are derived. These inequalities represent necessary and sufficient mobility conditions for the input and the output links of planar 4R linkages, as well as for its coupler link. The application of the foregoing constraint relations to linkage optimization using gradient -dependent methods is shown with an example that is fully solved resorting to Newton-Raphson's method and an interior penalty function.

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Introduction

Grashof's mobility criteria for RRRR plane linkages establish the conditions on the relative magnitudes of the links for the existence of double-crank, crank-rocker and double-rocker linkages. The proof of such criteria can be seen in either [1] or [2]. A recent reassessment of such criteria was given by Paul [3], who proved necessity and sufficiency of those. Paul also showed that two types of double-crank linkages should be distinguished, namely those with fully-revolving couplers and those with oscillating ones. The mobility of the coupler link is also analyzed in the present paper. Indeed, necessary and sufficient conditions for a fully-revolving and for an oscillating coupler are derived. Litvin [4] has established general conditions in algebraic form for the existence of cranks in closed kinematic chains of any type. He does not provide, however, specific formulae for specific types of linkages.

Grashof's conditions for the existence of the aforementioned types of linkages take the form of inequalities in which the lengths of the different links appear linearly. In synthesizing RRRR plane linkages for function generation, however, a system of linear equations not on the linkage lengths, but on a different set of parameters, arising from Freudenstein's equations, is to be solved. Given the RRRR plane linkage of Fig 1, let

$$k_1 = \frac{a_1^2 + a_2^2 - a_3^2 + a_4^2}{2a_2a_4}, \quad k_2 = \frac{a_1}{a_2}, \quad k_3 = \frac{a_1}{a_4}$$
(1)

The problem of linkage synthesis for function generation consists of finding a set of values $\{a_1, a_2, a_3, a_4\}$ for the link lengths that produce a prescribed set of input-output pairs $\{(\psi_i, \phi_i)\}_{i=1}^{n}$, where ψ and ϕ are the

input and the output angles, respectively.

Freudenstein's equation [5;6, p 297] allows one to compute the set $\{k_1, k_2, k_3\}$ for the prescribed input-putput values. This equation is the following:

$$k_1 + k_2 \cos \phi_i - k_3 \cos \phi_i = \cos(\phi_i - \phi_i), i = 1,..., n$$
 (2)

Since the problem contains three unknowns, three input-output values can be prescribed, eq (2) thus leading to a linear algebraic system of the form

$$Ak = b \tag{3}$$

where A is a 3x3 matrix, k and b being the 3-dimensional vectors given next:

$$A = \begin{cases} 1 & \cos\phi_{1} & -\cos\psi_{1} \\ 1 & \cos\phi_{2} & -\cos\psi_{2} \\ 1 & \cos\phi_{3} & -\cos\psi_{3} \end{cases}, \quad k = \begin{cases} k_{1} \\ k_{2} \\ k_{3} \end{cases}, \quad b = \begin{cases} \cos(\psi_{1} - \phi_{1}) \\ \cos(\psi_{2} - \phi_{2}) \\ \cos(\psi_{3} - \phi_{3}) \end{cases}$$
(4)

By inverting the system of equations (3) one can obtain the unique set of values $\{k_1, k_2, k_3\}$ that solve the proposed problem. Given the simple structure of matrix A and the low number of equations, matrix A can be inverted² explicitly. In fact, formulae are available in the literature [36, p 298] for the computation of k from (3). Once the value of k has been computed, the lengths a_i can be computed by inversion of eqs (1) as:

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[&]quot;Here it is assumed that the synthesis problem leading to eqs (1) is well posed, i e cases rendering either matrix A singular or vector b zero are discarded.

$$a_1 = 1, a_2 = \frac{1}{k_2}, a_3 = \frac{1}{k_2}, a_3 = \frac{1}{k_2} + \frac{\left[k_2^2 + k_3^2 + k_2^2 + k_3^2 - 2k_1k_2k_3\right]}{k_2k_3}, a_4 = \frac{1}{k_3}$$
 (5)

.

In the above discussion nothing prevents k_2 and k_3 from resulting negative, thus producing negative values for either a_2 or a_4 , which situation is next dealt with; a₁ is arbitrarily chosen unity, for a scaling of all lengths by the same factor does not alter the input-output relationship, whereas a₃ can always be made positive by a proper choice of the sign of the square root. Negative values of either a_2 or a_4 indicate that the angle ϕ or, correspondingly, ϕ , should be measured to an extension of eitherthe input or the output link as shown in Fig 2.1f it is necessary to verify Grashof's conditions within the synthesis process, then ${
m a_2}^-$ and ${
m a_4}^-$ should be computed using the absolute values of k_2 and k_3 . Introducing such absolute values, however, removes the smoothness of one side of Grashof's inequalities, which might be undesirable if the foregoing computations are to be performed within an optimization procedure requiring the computation of the gradient of that side of the inequalities. An alternative approach [7] consists of determining lower and upper bounds for the length of the coupler link if either the input or the output link, or both, is to be a crank. For RRRR plane linkages, however, this approach leads to Grashof's unsmooth inequalities. Waldron [8] and Waldron and Stevensen [9] have proposed alternate approaches based on graphical methods, whereas Gupta has proposed sufficient algebraic smooth inequalities [10, 11] guaranteeing the existence of an input link of the crank type. Necessary and sufficient algebraic smooth inequalities have been proposed for the existence of an input crank [12].

In what follows, a set of smooth inequalities is obtained, that is necessary and sufficient to produce an input, an output or a coupler link of either type, crank or rocker. This set of inequalities is meant to be adjoined to optimization programs to produce linkages whose links be of a given type, while meeting prescribed functional requirements for the classical problems of linkage synthesis, namely rigid-body guidance, path generation and function generation. By incorporating further requirements on the transmission angle, additionally, as discussed in [13], the aforementioned problems of linkage synthesis can be treated as nonlinear programming problems, which can be readily solved using standard optimization packages, normally available in any program library.

Derivation of mobility conditions

First a set of mobility conditions for the input link is derived. To do this, indices are dropped from Freudenstein's equation for simplicity. Next, the following identities are introduced:

$$\cos\phi = \frac{1 - \tan^2(\phi/2)}{1 + \tan^2(\phi/2)}, \sin\phi = \frac{2 \tan(\phi/2)}{1 + \tan^2(\phi/2)}$$
 (6)

Freudenstein's equation is thus transformed into

A
$$\tan^2(\phi/2) = 2B \tan(\phi/2) + C = 0$$
 (7)

where

$$A = k_1 - k_2 + (1 - k_3) \cos \psi$$
 (8a)

$$C = k_1 + k_2 - (1+k_3) \cos \psi$$
 (8c)

Eq (7), then, defines a quadratic equation in $tan(\phi/2)$ for a given value of ψ . If none of its coefficients vanishes, this equation produces two values of ϕ , given by

$$\phi_{1,2} = 2 \tan^{-1} \left[\frac{B \pm (B^2 - AC)^{1/2}}{A} \right]$$
 (9a)

The two values of $\phi(\psi)$ given above can be: i) both complex, in which case the input link is of the rocker type, the corresponding value of ψ lying outside of the mobility range of this link, ii) both real and distinct, thus corresponding to the two conjugate configurations of the linkage, or iii) both real and identical to each other, in which case ψ attains an extremum value, i.e. the output link attains a dead-point position.

Now, if $A(\psi)=0$ and $B(\psi)\neq0$, the set $\{k_1,k_2,k_3\}$ must observe the following relationship:

 $|k_1 - k_2| \le |1 - k_3|$

In this case, the left-hand side of eq (7) degenerates into a line. Viewed this line as a particular case of the general quadratic function, its *two intersections* with the tan $(\phi/2)$ - axis can be thought of lying at tan $(\phi/2) = C/2B$ and at infinity-which thus produces the two following values for corresponding conjugate configurations:

 $\phi_1 = \tan^{-1} (C/2B), \phi_2 = \pi$ (9b) The two values of ϕ appearing in (9b) can also be derived formally by taking the limit, as $A \rightarrow 0$, of both values appearing in (9a).

. .

Should both A and B vanish simultaneously, then C would necessarily vanish, as well, in which case, for this particular value of ψ , eq (7) would hold identically, for any value of ϕ . The zeroing of B implies, from eq (8b), $\psi = 0$ or π , which then yields $k_1 = k_3$ and $k_2 = 1$, or $k_1 = -k_3$ and $k_2 = -1$, respectively, in view of eqs (8a & c). In either case $a_2^2 = a_1^2$ and $a_4^2 = a_3^2$, i.e.the linkage is of the change-point type and either at $\psi = 0$ or $\psi = \pi$, the input and the fixed links are coincident, as well as the coupler and the output links. The linkage thus degenerates into a two-link open chain, for which ϕ can attain any real value, as predicted by eq (7).

Furthermore, if the discriminant of eq (7), i.e. the radical of expression (9a) is negative for all real values of ψ , then the four lengths a_1 , a_2 , $a_3 \& a_4$ do not define a closed quadrilateron.

In establishing Grashof's mobility criteria, it is always assumed that the linkage lengths define a quadrilateron, that is to say, each length is smaller than the sum of the remaining ones. Such closure condition is next derived in terms of the set $\{k_1, k_2, k_3\}$. This is more easily done if the non-closing condition is established, instead. The set $\{k_1, k_2, k_3\}$ will define a set of link lengths that does not constitute a quadrilateron if the radical of eq (9a) is negative for all real values of ψ , as already remarked. A rigorous analysis of this situation, that is not presented here due to space limitations, leads to the following set of inequalities:

$$(k_{2} - k_{1}k_{3})^{2} \ge k_{3}^{4} , (k_{1}^{2} - k_{2}^{2} + k_{3}^{2} - 1)^{2} \ge 4(k_{2} - k_{1}k_{3})^{2} , k_{1}^{2} - k_{2}^{2} - 1 \ge k_{3}^{2}$$
(10a)
$$(k_{2} - k_{1}k_{3})^{2} \le k_{3}^{4} , (k_{2} - k_{1}k_{3})^{2} \le k_{3}^{2}(k_{1}^{2} - k_{2}^{2} - 1)$$
(10b)

The following has then been proved: If the set $\{k_1, k_2, k_3\}$, as given by eqs (1), verifies either set of inequalities, (10a) or (10b), then the link lengths $a_i(i = 1, ..., 4)$ do not define a linkage.

It can be readily verified that if the second relation (10b) holds, then a_3 , as given by eq (5), turns to be imaginary. The non-existence of a linkage can thus be due to either of two possibilities: i) One of its links turns to have an imaginary length (if all three k's in (5) are real, only a_3 can have an imaginary value); ii) Even though all four lengths are real, they do not meet the closure condition.

Now, if the k's are obtained as the solution of eqs (2) with n = 3, then the said k's will yield a closing real quadrilateron, i.e.neither of relations (10) will hold. Although this is not formally proved here, this fact can be realized from the continuity of both sides of Freudenstein's equation, the real-valuedness of $\{(\psi_i, \psi_i)\}_1^3$ and the (assumed) nonsingularity of matrix A appearing in (4).

The application of formulae(9) to obtain the real values ϕ_1 , ϕ_2 corresponding to one single value of ψ should be made taking into account possible cancellations due to round-off errors. This can be taken care of if formula (9) is rewritten in a form that is more suitable for numerical stability, as indicated in [14]. This topic, however, is not further discussed here for it falls without the scope of the paper.

The two real values obtained from eq (9) are distinct, except at deadpoint positions of the input link. Those distinct values correspond to the two conjugate configurations of the linkage. Next, conditions for the full rotatability of the input link are obtained. To this end, the discriminant of eq (9a) is expanded and then zeroed. This leads to

$$\cos^2\psi + 2b\,\cos\psi + c = 0 \tag{11a}$$

with

$$b = \frac{k_2 - k_1 k_3}{k_3^2}, c = \frac{k_1^2 - k_2^2 - 1}{k_3^2}$$
(11b)

The real roots $\cos \psi_1$, $\cos \psi_2$, of eq (11a) then yield the two extremal values of ψ . These two roots are always real if a_3 is, but nothing prevents them from having absolute values larger than unity. The roots of eq (11a) are, in fact,

$$\cdots \quad \cos\psi_{1,2} = b \pm (b^2 - c)^{1/2}$$
 (12a)

which, in terms of the link lengths, produce

$$\cos\psi_{1,2} = \frac{a_1^2 + a_2^2 - (a_3 + a_4)^2}{2a_1^2a_2}$$
(12b)

from which it is clear that both roots are real if a_3 is. Now, if the input link is to be of the crank type, then these two roots should yield complex values of ϕ_1 and ϕ_2 , which is only possible if both roots (12a) have absolute values larger than unity. This is exactly the necessary and sufficient condition for an input link to be of the crank type, provided the k's were obtained from a well-posed synthesis problem. This condition is, then

$$|-b \pm (b^2 - c)|^{1/2} |> 1$$
 (13)

which, if squared, still holds, for both sides of (13) are positive. Thus, it is equivalent to the following two inequalities:

$$2b^2 - c - 1 > 2b (b^2 - c)^{1/2}, 2b^2 - c - 1 > -2b (b^2 - c)^{1/2}$$
 (14)

which are in turn equivalent to the single one given next:

$$2b^2 - c - 1 > 2|b|(b^2 - c)^{1/2}$$
 (15)

which requires that its left-hand side be positive, i e

$$2b^2 - c - 1 > 0$$
 (16)

Now, given inequality (16), relation (15) still holds if its both sides are squared. This leads to

$$(c+1)^2 - 4b^2 > 0$$
 (17)

If definitions (11b) are recalled and substituted into relations (16) and (17), the next two inequalities are obtained, in the space of k_1 , k_2 , k_3 :

$$2 (k_2 - k_1 k_3)^2 - k_3^2 (k_1^2 - k_2^2 + k_3^2 - 1) > 0$$
 (18 a)

$$[(k_1 - k_3)^2 - (k_2 - 1)^2] [(k_1 + k_3)^2 - (k_2 + 1)^2] > 0$$
 (18 b)

Summarizing, then, one has proved that: The necessary and sufficient conditions for the synthesis problem leading to cqs (2) to produce an input crank is that both inequalities (18 a and b) hold.

The full rotatability conditions for the output crank can be obtained analogously. It is far simpler, however, in view of the symmetry of definitions (11b) with respect to a_2 and a_4 in Freudenstein's equation, to exchange the roles of variables k_2 and $-k_3$ in inequalities (18 a and b). This yields

$$2(-k_3 + k_1k_2)^2 - k_2^2(k_1^2 + k_2^2 - k_3^2 - 1) > 0$$
 (19 a)

$$[(k_1 + k_2)^2 - (-k_3 - 1)^2] [(k_1 - k_2)^2 - (-k_3 + 1)^2] > 0$$
 (19 b)

One then has: The necessary and sufficient conditions for the synthesis problem leading to eqs (2) to produce an output erank is that both inequalities (19 a and b) hold.

Conditions for the existence of rockers are next derived. One possible way of establishing them is saying that "an input link is of the rocker type if inequalities 18 a and b) do not hold simultaneously, whereas an output link is of the rocker type of inequalities (19 a and b) do not hold simultaneously". Since the violation of the said inequalities presents various alternatives, it does not guarantee the existence of a rocker link. An alternative approach, specifying the extremal values of the variable of interest, either ψ or ϕ , is presented next.

If an input rocker is required to have mobility only within the range $\psi_1 \leq \psi \leq \psi_2$, then the discriminant of eq (9a) is zeroed at these values, i.e.

$$\cos^2 \psi_i + 2 b \cos \psi_i + c = 0, \quad i = 1.2$$
 (20)

The two given roots of eq (20) satisfy

$$b = -\frac{1}{2} (\cos \psi_1 + \cos \psi_2), c = \cos \psi_1 \cos \psi_2$$
 (21)

In terms of $k_1^{}$, $k_2^{}$ and $k_3^{}$, eqs (21) lead to

$$2(k_2 - k_1 k_3) + k_3^2 (\cos \psi_1 + \cos \psi_2) = 0$$
 (22a)

and

$$k_1^2 - k_2^2 - 1 - k_3^2 \cos \psi_1 \cos \psi_2 = 0$$
 (22b)

Eqs (22) constitute a nonlinear algebraic system in two equations and three unknowns, i.e. it is underdetermined. These are subject to the second

order constraints guaranteeing that ψ_1 be a minimum and ψ_2 be a maximum, which are

$$\psi''(\phi_1) \ge 0, \ \psi''(\phi_2) \le 0$$
 (23)

where, ϕ_i , defined as $\phi(\psi_i)$ can be obtained by substitution in eq (7). At a stationary value of ψ , ψ' vanishes, $\psi''(\phi)$ reducing to

$$\psi''(\phi) = \frac{(A^2 + B^2)^2 A}{2 G^2}$$
(24a)

with A and B given as in definitions (8), and

$$G = [B^{2} (1 - k_{3}) - A^{2} (1 + k_{3})] \sin \psi + 2 AB \cos \psi \qquad (24b)$$

Relations (23) then can be expressed as

$$A(\psi_1) \ge 0, A(\psi_2) \le 0$$
 (25a)

or, from definition (8a), in terms of k_1 , k_2 and k_3 , as

$$k_1 - k_2 + (1 - k_3) \cos \psi_1 \ge 0$$
 (25b)

and

$$k_1 - k_2 + (1 - k_3) \cos \psi_2 \le 0$$
 (25c)

Relations (22a and b) and (25b & c) alone do not allow the computation of k. These should be incorporated into an optimization problem, e.g. one minimizing a norm of the structural error or maximizing a norm of the mechanical advantage within the range of motion.

Now, if an output rocker is required, whose motion be defined in the interval $\phi_1 \leq \phi' \leq \phi_2$, this can be accomplished paralleling the foregoing procedure. It is far simpler, however, to derive the corresponding

relations by exchanging the roles of k_2 and $-k_3$ in relations (23a & b) and (25b & c). This produces

$$2(k_1k_2 - k_3) + k_2^2 (\cos\phi_1 + \cos\phi_2) = 0$$
 (26a)

$$k_1^2 - k_3^2 - 1 - k_2^2 \cos\phi_1 \cos\phi_2 = 0$$
 (26b)

and

$$k_1 + k_3 + (1 + k_2) \cos \phi_1 \ge 0$$
 (27a)

$$k_1 + k_3 + (1 + k_2) \cos \phi_2 \le 0$$
 (27b)

Mobility conditions for the coupler are next derived. An analysis similar to the one leading to Freudenstein's equation yields

$$m_1 + m_2 \cos \theta + m_3 \cos \psi = \cos (\psi - \theta)$$
 (28)

with

• •

$$m_1 = \frac{a_4^2 - a_1^2 - a_2^2 - a_3^2}{2 a_2 a_3} , \quad m_2 = \frac{a_1}{a_2} , \quad m_3 = \frac{a_1}{a_3}$$
(29)

Substitution of identities (6) for angle Ψ in the latter equation yields

$$J \tan^2(\psi/2) - 2K \tan(\psi/2) + L = 0$$
 (30)

with

•

$$J = m_1 + m_2 \cos\theta - (m_3 - \cos\theta)$$
(31a)

K = sin0 (31b)

$$t = m_1 + m_2 \cos \theta + m_3 - \cos \theta$$
 (31c)

One then has, for a given value of 0, from eq (30),

$$\Psi_{1,2} = 2 \tan^{-1} \left[\frac{K \pm (K^2 - JL)^{1/2}}{J} \right]$$
 (32)

Eq (32) thus yields two different values of ψ , corresponding to the conjugate configurations of the linkage, except at extremal positions of the coupler, where the radical vanishes. If the coupler is to have full rotatability, the radical should not vanish for real values of θ . The conditions under which this happens are derived paralleling the procedure leading to relations (18a & b), which produces the following set of inequalities:

$$1 + m_1^2 - m_2^2 + m_3^2 > 0 \tag{32a}$$

$$(1 - m_1^2 - m_2^2 + m_3^2)^2 - 4m_1^2 > 0$$
 (32b)

One then has proved: The well-posed synthesis problem producing m_1 , m_2 and m_3 from eq (28) yields a coupler link possessing full rotatability if, and only if, relations (32a f b) hold.

Mobility conditions for the coupler, considering its motion with respect to the output link, are derived analogously. These are obtained from the equation

$$n_1 + n_2 \cos \theta - n_3 \cos \phi = \cos(\phi - \theta) \tag{33}$$

with

$$n_1 = \frac{a_1^2 - a_2^2 + a_3^2 + a_4^2}{2a_3a_4}, \qquad n_2 = \frac{a_1}{a_4}, \qquad n_3 = \frac{a_1}{a_3}$$
(34)

Eq (33) leads to $M \tan^2(\psi/2) = 2N \tan(\psi/2) + P = 0$ (34)

where

where

. .

$$M = n_1 + n_2 \cos \theta + n_3 + \cos \theta$$
 (35a)

$$P = n_1 + n_2 \cos\theta - (n_3 + \cos\theta)$$
 (35c)

The roots of eq (34) are, thus,

$$\psi_{1,2} = 2\tan^{-1} \left\{ \frac{N \pm (N^2 \pm MP)^{1/2}}{M} \right\}$$
 (36)

The conditions sought are then derived from the zeroing of the radical of eq (36). This leads to the following set of inequalities:

$$1 + n_1^2 - n_2^2 + n_3^2 > 0$$
 (37a)

$$(1 - n_1^2 - n_2^2 + n_3^2)^2 - 4n_1^2 > 0$$
 (37b)

That is: The coupler link of a linkage whose lengths are derived from eq (33) has full rotatability with respect to the output link if, and only if, its lengths satisfy relations $(37a \ 6 \ b)$.

What is meant in the last paragraph under full rotatability of the coupler with respect to the output link is that $0'(\phi)$ does not vanish, for any real value of ϕ . This does not mean, however, that this is equivalent to full rotatability of the coupler link. In fact, if the output link does not possess full rotatability, then $\phi'(\phi)$ vanishes for two distinct values of ψ . Hence, even if $0'(\phi)$ does not vanish, $0'(\psi) = 0'(\phi)\phi'(\psi)$ does, the coupler link thus lacking full rotatability.

On the other side, conditions for oscillating couplers can be obtained paralleling the procedure followed for the input and the output links. The obtention of such conditions is straightforward from the foregoing analysis, for which reason the subject is not further discussed here. Finally, changepoint mechanisms are characterized within this context as those for which the left-hand side of at least one of inequalities (18a) or (18b) (or, equivalently (19a) or (19b)) vanishes.

The inequalities derived in this paper are now used as constraints of an optimization problem of linkage synthesis. This is solved using Newton-Raphson's method, which requires not only first, but also second derivatives of both the objective function and the constraints. The procedure is illustrated with one fully solved example.

Example

Synthesize a RRRR plane linkage, as the one shown in Fig 1, to produce the input-output relation appearing in Table 1. This linkage should approximate the synthesis equations with the least possible r.m.s. error, while its input link should be a crank.

This problem was solved in [15] without considereing the crank-type restriction. The least-square error linkage thus obtained turned to be of the rocker-rocker type.

The synthesis problem at hand is formulated as follows: The synthesis equations are of the type of eq (3), except that A is now a 5 x 3-matrix, whereas vector b is 5-dimensional. The ith row of matrix A and the ith component of vector b are, respectively.

$$A_{i} = [1, \cos\phi_{i}, -\cos\psi_{i}], b_{i} = \cos(\psi_{i} - \phi_{i})$$
 (38)

This problem thus leads to an overdetermined system of linear equations which, in general, has no exact solution. In this case, a vector k is sought that minimizes a norm of the error e, defined as

e = Ak - b (39)

which is clearly a 5-dimensional vector. If the Euclidean norm is to be minimized, without imposing any further constraint, then the minimizing value of k can be expressed explicitly in terms of the Noore-Penrose generalized inverse [16], which can be computed very efficiently using Householder reflections as already shown in [17, 18]. Should a further constraint be imposed on the synthesis problem, then the above-mentioned generalized inverse is not applicable any more. A possible way of solving this problem is via a penalty function [19], which is next introduced. The problem is now formulated as: "Minimize, under k, the objective function z given as

 $z = \frac{1}{2} (Ak - b)^{T} (Ak - b)$ (40)

i.e. half the square of the Euclidean norm of the error, subject to inequalities (18a & b), whose holding is necessary and sufficient for a crank-type input link".

Solution:

Let f_1 and f_2 represent the left-hand sides of inequalities (18a & b), respectively. Now assume a feasible linkage is given, i.e. one defined by a particular vector k^0 that satisfies both given inequalities, which in general is not optimal. This linkage can be improved by minimizing the new objetive function

$$\psi(k; r_1 = z(k) + r_1(\frac{1}{f_1} + \frac{1}{f_2})$$
 (41)

subject to no further constraints.

The second term of the right-hand side of eq (41) is referred to as the penalty term. This is the product of the positive weighting factor r_1 times the sum in parenthesis, referred to as the penalty function. Criteria for selecting a suitable value of r_1 are given in [19, pp 156-196], but the simplest one is to choose it so as to render the penalty term a given fraction of $z(k^0)$. Next, the value of k minimizing ψ , k^1 , is used as a "guess" value of k to minimize a new objetive function with a new penalizing factor, r_2 , a fraction of r_1 . The procedure is repeated a few number of times, say p, which produces an equal number of pairs $\{(k^i, r_i)\}_{1}^{p}$ with $r_i > r_{i+1} > 0$; these can then be fitted to a suitable function, as shown next, the solution to the original constrained problem being obtained by extrapolation, with $r \to 0$.

Since both function z and the penalty term are infinitely many times differentiable, each unconstrained optimization problem meant to minimize $\psi_i = \psi(k; r_i)$, for i = 1, ..., p can be solved using a gradient method or even the Newton-Raphson method [20, pp 249-251]. In any instance, the roots of the gradient of ψ_i with respect to k, are to be computed. The said gradient is given as

$$\nabla \psi_i = A^T (Ak - b) - r_i \left(\frac{\nabla f_1}{f_1^2} + \frac{\nabla f_2}{f_2^2} \right); \quad i = 1, ..., p$$
 (42)

with

ζ,

$$\nabla f_{1} = 2 \begin{bmatrix} (k_{1}k_{3} - 2k_{2})k_{3} \\ 2(k_{2} - k_{1}k_{3}) + k_{2}k_{3}^{2} \\ (k_{1}^{2} + k_{2}^{2} - 2k_{3}^{2} + 1)k_{3} - 2k_{1}k_{2} \end{bmatrix}$$
(43).

$$\nabla f_{2} = 2 \left(\begin{bmatrix} k_{1} - k_{3} \\ 1 - k_{2} \\ k_{3} - k_{1} \end{bmatrix} \begin{bmatrix} (k_{1} + k_{3})^{2} - (k_{2} + 1)^{2} \end{bmatrix} + \begin{bmatrix} k_{1} + k_{3} \\ -k_{2} - k_{1} \end{bmatrix} \left[(k_{1} - k_{3})^{2} - (k_{2} - 1)^{2} \end{bmatrix} \right)$$

$$+ \begin{bmatrix} k_{1} + k_{3} \\ -k_{2} - 1 \\ k_{1} + k_{3} \end{bmatrix} \left[(k_{1} - k_{3})^{2} - (k_{2} - 1)^{2} \end{bmatrix} \right)$$

$$(44)$$

The application of Newton-Raphson's method to the computation of the roots of $\nabla \psi_i$ requires computing the Jacobian matrix J of $\nabla \psi_i$, i.e. $\nabla^2 \psi_i$, with respect to k. This is readily computed as

$$\nabla^{2} \psi_{i} = A^{T}A - r_{i} \left[\frac{f_{1}^{2} \nabla^{2} f_{1} - 2f_{1} \nabla f_{1} (\nabla f_{1})^{T}}{f_{1}^{4}} + \frac{f_{2}^{2} \nabla^{2} f_{2} - 2f_{2} \nabla f_{2} (\nabla f_{2})^{T}}{f_{2}^{4}} \right] = A^{T}A - r_{i} \left[\frac{\nabla^{2} f_{1}}{f_{1}^{2}} + \frac{\nabla^{2} f_{2}}{f_{2}^{2}} - 2(\frac{\nabla f_{1} (\nabla f_{1})^{T}}{f_{1}^{3}} + \frac{\nabla f_{2} (\nabla f_{2})^{T}}{f_{2}^{3}}) \right]$$
(45)

with

$$\nabla^{2} f_{1} = 2 \begin{bmatrix} k_{3}^{2} & -2 k_{3} & 2 (k_{1}k_{3} - k_{2}) \\ 2 + k_{3}^{2} & 2 (k_{2}k_{3} - k_{1}) \\ sym & (k_{1}^{2} + k_{2}^{2} - 6 k_{3}^{2} + 1) \end{bmatrix}$$
(46)

19.

$$\nabla^{2} f_{2} = 4 \begin{bmatrix} 3k_{1}^{2} - k_{2}^{2} - k_{3}^{2} - 1 & 2(k_{3} - k_{1}k_{2}) & 2(k_{2} - k_{1}k_{3}) \\ & -k_{1}^{2} + 3k_{2}^{2} - k_{3}^{2} - 1 & 2(k_{1} - k_{2}k_{3}) \\ sym & -k_{1}^{2} - k_{2}^{2} + k_{3}^{2} - 1 \end{bmatrix}$$
(47)

The Newton-Raphson method with damping, implemented with subroutine NRDAMP [18, pp 39-50] was used to solve the foregoing problem. The results obtained are shown in Table 2, where k^0 was chosen randomly so as to produce a linkage of the input-crank type. The successive values of r_i (i=1,2,3) employed were 0.1, 0.01 and 0.001.

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	ψ =		140°	130°	: . 110°	100°	90°
•	¢ •	· · ·	80°	74°	64°.	58°	50°

	0.205592		k ² =	0.206277	1	i	0.195852
k ¹ =	0.577607			0.793621	•	k ³ -	0.843946
	0.195602			0.220257			0.208259

The values of Table 2 were interpolated to the curve

$$k(r) = \alpha + \beta r^{1/2} + \gamma r$$

which produced the following:

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 $\alpha = \begin{bmatrix} 0, 189301 \\ 0, 918406 \\ 0, 199494 \end{bmatrix}, \qquad \beta = \begin{bmatrix} 0, 224439 \\ -2, 496524 \\ 0, 309335 \end{bmatrix}, \qquad \gamma = \begin{bmatrix} -0, 546832 \\ 4, 486709 \\ -1, 017054 \end{bmatrix}$

Thus the optimizing value k* was obtained as

 $k^* = k(0) = \alpha$

1

which produced the linkage

 $a_1 = 1$, $a_2 = 1.088843$, $a_3 = 5.024554$, $a_4 = 5.012682$ for which the least-square error is $e = [0.001602, 0.011487, -0.034524, -0.032521, 0.013597]^T$; ||e|| = 0.050685

Conclusions

Mobility conditions for RRRR planar linkages have been derived, that are equivalent to Grashof's mobility criteria. The conditions presented here differ from the usual ones in that they are established as relations on analytic functions of variables that are nonlinear combinations of the link lengths, rather than on unsmooth (because of the absolute-value function appearing there) functions of the link lengths. The incorporation of the conditions derived here as inequality constraints of optimization problems allows their solution via gradient-dependent methods, as shown with an example of constrained least-square approximate synthesis. This problem was solved using the Newton-Raphson method, for both the approximation error and the constraints are readily differentiable infinitely many times. Hence the computation of second partial derivatives, as required by the Newton-Raphson method, is quickly executed. Finally, the quadratic-convergence property of the said method, close to a solution, was made apparent by the quick convergence of each of the three nonlinear systems of equations that were solved in the example. In fact, each solution was obtained after at most three iterations.

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I.

DIVISION DE EDUCACION CONTINUA FACULTAD DE INGENIERIA U.N.A.M.

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DISEÑO CINEMATICO DE MAQUINARIA

OPTIMAL SYNTHESIS OF LINKAGES USING HOUSEHOLDER REFLECTIONS

JORGE ANGELES

JUNIO 1984

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OPTIMAL SYNTHESIS OF LINKAGES USING HOUSEHOLDER REFLECTIONS

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ABSTRACT

The uncostrained overdetermined problem of Minematic linkage synthesis is solved in an efficleat way using Householder reflections. The problem formulation leads to a system of either linear or conlinear equations in more equations then unknowns. The linear problem is solved directly by application of a finite number of successive reflections to the space of unknowns with the purpose of taking the system of equations into upper triangular form, which allows for the computation of the unknowns by back substitution. The nonlinear problem is solved via the Newton-Raphson method which computes, at each iteration, the correction to the vector of unknowns as the least-square solution to an overdotermined linear system in exactly the same way as described before for linear problems. Introduction of the said method produces accurate results in relatively short processing times, as shown in the examples presented.

ZUSAMMENTASSUNG

Das uneingeschränkte und überbestimmte Problem der kinematischen Getrichesynthese wird effisient gelost wit Hilfe der Householder-Spiegelungen. Die Problemstellung laitet zu einem Systam Von entweder linearen oder nichtlinearen Gleichungen mit mehr Gleichungen als Unbekannten, Das lineare Problem wird direkt golöst mittels Anvendung einer finiten Zahl aufeinanderfolgenden Spiegelungen zum Raum der Unbekannten mit dem Ziel des Übertragens des Cleichungssystems zu einer höheren dreiczkigen Form, welche die Rechaung der Unbekannten durch Rückersetsung erlaubt. Des nicht-lineare Problem wird mitteles der Newton-Raphson-Methode gelöst, die zu jeder lteration die Besserungen der Unbekannton sus der wenigeren Quadraten-Lösung zu einem Überbestimmten linearan Gleichungssystem arrochnet, auf der gleichen Weise wie bei der Methode für lineare Systeme schon beschrieben wurde. Die Einführung dieser Methode führt zu deutlichen Erfolgen in relativ hurzen Prozessierzeiten, wie mittels der eingeschloßenen Brispielen gezeigt wird.

NOMENCLACURE

A; upper-case underlined character, an exe matrix. <u>-</u>1, Ŷ the inverse of A, when A is square and nonsin-

- 1.11ar
- Lie. Cranspose of A
- lower-case underlined latin character, an m-(incusional vector
- 4 1 the absolute value of a, when a is real; the

modulus of a, when a is complex.

- all:the Eucliean norm of vector a, i.e. the square root of the sum of the squares of its components
- detA: the determinant of the square matrix A f(x): an m-dimensional vector function of the m-dimen-
- sional vector argument X

f'(x):the Jacobian man matrix of f with respect to x

PROBLEM FORMULATION

The equations srising in the reals of kinematic synthesis of linkages constitute either linear or nonlinear algebraic¹ systems (1,2), whose unknowns are the geometric parameters (langths and angles) of the linkage. If these paramenters are arranged within the n-dimensional vector x, the said equations are of the form

1.25

where A and b are a known man matrix and an m-dimensional known vector, respectively, when the system is linear. If it is nonlinear, then the synthesis equations are of the form f(x)=0 (2)

f being an m-dimensional vector containing a set of m scalar functions f. (x) whose arguments are the unknown parameters of the linkage. When the number of specifiad conditions to be net by the linkage matches that of the unknowns, matrix A in (1) is square and vector f is of dimension n. In most technical problems, however ar, the number of prescribed conditions surpasses that of geometric paramenters available, the linkage synthesis problem thus leading to an overdetermined system of equations. This class of systems in general does not admit an exact solution, but it is possible to find a vector x that randers the quadratic error a minimum. Thus, the least-squares problem can be stated

"Pind the value of x that minimizes the Euclidean norm of either Ax-b, or that of f(x), depending on whether the system Is linear of nonlinear".

The linear overdetermined system (1) admits a unique solution x that renders | [Ax-b] & minimum, provided A is of full rank, i.e., if rankA-n. This value is given is (3)

$$\underline{\mathbf{x}}_{1} = (\underline{\mathbf{A}}^{\mathrm{T}} \underline{\mathbf{A}})^{-1} \underline{\mathbf{A}}^{\mathrm{T}} \underline{\mathbf{b}}$$
(3)

where (A A) A is called a "Moore-Penrose Jeneralized

Zsea the nomenclature for the definition of this term,

algebraic as opposed to differential or integral equations.

inverse of A". An extensive treatment of the linear least-squral problem is found in $(\underline{4})$.

The nonlinear problem may admit multiple local minima; these can be found by application of the Newton-Raphson method (5), which at each iteration, computes the currection vector Δx_i as the least-square solution to the overdetermined linear system

$$f^{1}(x_{1})\Delta x_{1} = -f(x_{2})$$
 (4)

This is a system like that appearing in eq.(1). Thus, its least-square solution is

$$\Delta \mathbf{x}_{\mathbf{k}}^{*} - \left[\mathbf{f}^{*}_{\mathbf{k}} (\mathbf{x}_{\mathbf{k}})^{T} \mathbf{f}^{*}_{\mathbf{k}} (\mathbf{x}_{\mathbf{k}}) \right]^{-1} \mathbf{f}^{*}_{\mathbf{k}} (\mathbf{x}_{\mathbf{k}})^{T} \mathbf{f}^{*}_{\mathbf{k}} (\mathbf{x}_{\mathbf{k}})$$
(5)

The new value of the unknown vector is then

$$\sum_{k+1} \sum_{k} \frac{4\Delta x}{\Delta k}$$

The procedure is stopped when the fuclidean norm of the correction vector is sufficiently small within the imposed accuracy, i.e. when

 ε boing a "amall" real positive number. The problem thut, whether linear or nonlinear, reduces to compute the minimizing value x given by eq. (3). An efficiert way of computing this value, outlined next, doe, not require to invert any matrix. The computation is upper by application of Householder reflections.

HOUS EHOLDER REFLECTIONS

An extensive account of this topic can be found in the specialized literature $(\underline{4},\underline{6})$. For this teason, this theory is not treated here. A Householder telfection is a linear, improper orthogonal and symmetric transformation, i.e., if R is its max matrix representation, then

220-2	())
Qx, =d	(10)

where U is an upper triangular nummatrix and Q is the (m-n)un zero matrix, r and d being n-and (m-n)dimensional vectors, with $d \neq 0$. Thus, eq. (9) is determined and can readily be solved by back substitiation, its solution x being the least-square solution to the overdetermined system. Eq. (10) is inconsistent and ||d|| represents the Euclidean norm of the error in the approximation. Since the original system (1) is transformed into (9),(10) via a succession of orthogonal transformations, the error in the transformed coordinates, d, has the same Euclidean norm as that in the original coordinates. Hence, ||d||is the error associated with the original system.

APPLICATIONS TO KINEMATIC LINKAGE SYNTHESIS

Although in many practical applications the problems of linkage synthesis involve inequality constraints, still a considerably large class of synthesis problems are unconstrained. Moreover, afficient optimization techniques exist that handle inequality constraints by introducing suitable penalty functions (2), thus turning the problem an unconstrained one, for these reasons, the study of unconstrained optimig ation problems is of substantial technical interest. Applications to linkage synthesis problems are next illustrated with two examples.

Example 1. Synthesis of an RSSR function generator

The layout of an RSSR linkage, shown in Fig 1, indicates the different geometric parameters of this linkaget a_1, a_2, a_3 , are the lengths of the output-, coupler-and input links, respectively; a_4 is the distance between the axes of the input and the output links; a_4 is the angle between the aforementioned axes, positive about ED; a_4 and a_5 are distances of points C and O along the axes of the cutput-and the input links, respectively. All over, the sign convention of Denavit and Hartenberg (1,pp. 344-345) is observed. The input angle is ψ and the output angle is ϕ . For matching six pairs of input-output values (ψ_1, ϕ_2) with this linkage, Denavit and Hartenberg (1,pp. 335-362) established the following relation

$$k_{1}\cos\phi_{j}+k_{2}\sin\phi_{j}+k_{3}\cos\phi_{j}+k_{4}\sin\phi_{j}+k_{5}(\sin\phi_{j}\cos\phi_{j}-\cos\phi_{4}\sin\phi_{j}\cos\phi_{j})+k_{6}=$$

$$=\cos\phi_{1}\cos\phi_{2}\cos\phi_{1}+\cos\phi_{2}\sin\phi_{3}\sin\phi_{3}$$
(11)

where

$$k_{1} = \frac{a_{4} + a_{4} \pi i n \alpha_{4} \pi i n \alpha_{4}}{a_{3}} + k_{2} = \frac{a_{4} \pi i n \alpha_{4} - a_{4} \pi i n \alpha_{4}}{a_{3}}$$

$$k_{3} = \frac{a_{4}}{a_{1} \cos \phi_{0}} + k_{4} = \frac{a_{1} \pi i n \alpha_{4}}{a_{1} \cos \phi_{0}} + k_{5} = t \operatorname{an} \phi_{0}$$

$$k_{4} = \frac{a_{1}^{2} - a_{2}^{2} + a_{3}^{2} + a_{4}^{2} + a_{4}^{2} + 2a_{1}^{2} + 2a_{1}^{$$

o, being assigned.

In the latter definitions, ϕ_0 measures the location of the zero of the output dial from the dotted line passing through C, parallel to line ED, as shown in Fig 1.



Fig 1 An RSSR linkage

For six precision-point synthesis, eqs. (11) yield a system of six linear equations in six unknowns which, when nonsingular, produces unique values k_1, k_2, \dots, k_n . With these values known, the linkage parameters are computed from eqs. (12) for a given value of a_1 . If more than six precision points are required, however, the system becomes overdetermined, in which case in efficient method to obtain its leastsquare solution is via Howseholder reflections.

In (6), Sub and Mecklenburg molve the overdetermined uncomprained problem of this linkage with 19 prescribed input-output values. For comparison purposes, the solution developed in this example makes use of the same prescribed values. These are shown in Table 1

The method employed in (8) is that of Powell's (9), which does not require the computation of derivatives and teads quickly to convergence for quadratic functions of the independent variables. At this point, two remarks are in order: First, the defivatives of the synthesis equations are easily computed from either Devanit and Hartenberg's formulation, eqs. (11), or from 5uh and Redcliffe's formulation (2), the first one being advantageous because of producing a linear system of equations. Second: The objective function of Suh and Hecklenburg's (8) is quadratic in the synthesis which, in turn, are quadratic in the independent variables; thus, their objective function is quartic in the independent variables, for which reason the quick convergence properties of Powell's method are not fully utilized. Furthermore, squaring the synthesis functions may introduce spurious local mining, as is apparent form the fact that three optimal eplutions are reported in (8).

TABLE 1. Specified input-output pairs for the synthemis of the RSSR function generating linkage.

	¢(degraces)	¢(degraek)
1	0.0	0.0
2	5.0	2,4
3	10.0	5,1
4	15.0	8.2
5	20.0	11.5
6	25.0	15.2
7	30.0	19.1
8	35.0	23,3
9	40.0	27.7
10	45.0	32.3
1	50.0	37.2
12	55.0	42.3
13	60.0	47.5
4	65.0	\$3.0
15	70.0	58.7
6	75.0	64.6
7	80,0	70,9
8	65.0	78.0
.9	90.0	90.0

One advantage of using Householder reflections is that no explicit equating is required, and the unique solution is obtained <u>directly</u> by the application of n(=6) reflections. Another advantage is that, since less computations are required, as compared to Powell's method, the round-off error is lowered. The approximation error obtained using each method is shown in Table 2.

The root mean square errors were exampledly the same: thet obtained by Powell's method was 0.00185269; whereas the one obtained by Householder teflections, 0.00182254; However, the differences in the resulting linkage parameters were more notorious. These are

Solution by	Solution by		
Fowell's method	Howseholder's method		
n - 1.253803	4 = 0.911269 .		
7.759566	4 = 2.620568		
7.435003	4 = 0.603577		

2.262110	■,=-1,L86240
•41.375270	e ¹ 2.417556

In this problem, $\mathbf{a}_{\underline{A}}$ was set equal to L, whereas $\mathbf{a}_{\underline{A}}$, equal to 90?

TABLE 2. Approximation error in overdetermined RSSR linkage synthesis

AFROXIMATION EXERCE USING APROXIMATION FRADE USING POWELL'S METHOD HOUSEHOLDER'S METHOD

	(degrees)	(degrees)
1	0.00000000	0.01420369
2	00120000	0.00100785
3	03060000	03515343
4	0.02330000	0.0159082/
5	-,02680000	03451603
6	0.03260000	0.02596072
7	0.01600000	0.01079286
8	0.03830000	0.03440434
9	0.01520000	0.01196148
10	03790000	04106664
11	00640000	00968497
2	0.02270000	0.01930349
13	~.04220000	04497542
4	00020000	00153544
15	0.03350000	0.03434200
t6	0.01600000	0.01324614
17	0.00020000	0.00139267
8	01250000	01970407
19	0.02980000	0.00412793

Example 2. Synchesis of the RR plane dyad for tigidbody guidance

A rigid body (shaded rectangle) appears in Fig 2, in "reference" configuration C_0 and in a different configuration C. Each configuration is defined by the position of a point, R, and angle, 0. In that figure, 0 represents the origin of the complex plane, and the Prove represent complex numbers associated with the location of the labelled points. The purpose of this class of synthesis problem is to locate point A whose reference and successive positions, A, A, (j=1, ...,n) lie on a circumference centered at B, lor which reason, A, and B are called, respectively, "circular" and "centfel" points, within the Burmester Theory (10). Thus, AB, can constitute a rigid link to guide the rigid body. This is an SR plane dyad.



Fig 2 An RR plane dyad to guide a rigid body through a successive configurations

. The constancy of the length of line BA throughout its a configurations leads to

$$|\mathbf{e}^{10}_{\mathbf{j}}(\mathbf{s}_{0}-\mathbf{r}_{0})+\mathbf{r}_{1}-\mathbf{b}|^{2}-|\mathbf{s}_{0}-\mathbf{b}|^{2},\mathbf{j}-\mathbf{1},\ldots,\mathbf{n}$$
(13)

where $\theta^* \Xi \theta_1 - \theta_2$. Eqs. (13) constitute the synthesis equations for this problem, z_1 and b being the unknowns. It is well known (2, p. 146) that this problem allows to conduct a rigid body through five specified configurations. Some technical problems, however, may require to guide the body through more than tive configurations, as shown in Table 3. Different synthese's were obtained for these, starting from the first 6 coofigurations, then adding the pext ones, one at each time, until the 16 configurations were included

TABLE .	э.	Success	va	¢anf	Égurat:	ions	of	a.	rigid	body
---------	----	---------	----	------	---------	------	----	----	-------	------

j.	x, (ca)	y _j (cm)	θ _i (degrees)
¢	7.880	-0.260	313.720
1	8.490	-7.290	332.330
2	7.680	2.820	349.930
3	6.300	4.210	353.180
4	4.58G	4.950	359.870
5	2.740	5.010	355.840
6	1.010	4.410	356.300
1	0.259	3.880	3.900
6	-0.400	-3.090	3.670
9	G. 250	-3.760	3.690
10	1.000	-4.290	4.150
11	2.730	-4.890	5.120
12	4.360	-4.830	6.810
13	6.280	-4.090	10.000
14	7.660	-2.700	13.000
15	8.440	-0.610	18.000
16	7.790	-2.690	46.270

The procedure converged for all given inital guesses, produced by means of a random number generating subprogram, in less than 50 iterations (usually around 20)Kontrary to the determined case (5 prescribed configurations), for which two different mean ingful solutions exist, for the cases tried here the procedure converged always to the same single solution, except for 6 and 17 configurations, which produced two different solutions. The error in the approx imation was normalized, to yield a dimensionisse number, in the following way: Let

$$f_{j} = |\mathbf{a}^{10}_{j}(\mathbf{a}_{0} - \mathbf{v}_{0}) + \mathbf{r}_{j} - |\mathbf{a}_{0} - \mathbf{b}|^{2}, \mathbf{j} = 1, \dots, \mathbf{a}$$
(14)

If the synthesis were exact, then all f, would be negligibly small. In approximate synthesed, however, these functions attain finite values. The kinematic meaning of these values is that they represent the difference between the length of the RR dyad in its initial configuration, and that in its jth fonfiguration, i.e. A, B-A, B, if the synthesized linkage were to satisfy the prescribed conditions exactly. The dimensionless error in the approximation, a., associated with the jth configuration, is then $e_j = |f_j|/|a_0 - b|^2, j = 1, \dots, m$ (15)

where a, and b are those obtained from the leastsquare solution to the nonlinear system of equations. Notice that the errors thus defined are quadratic. To obtain a representative value of the overall error, the everage of the square roots of the m errors defined in (14) should be taken, i.e.

Some of the results obtained are shown next,

TABLE 4. Overdetermined synthesis of the LR dyad for rigid-body guidance. For 6 configurations.

First solution:	Second solution:		
-0.961467-12.826960	<pre>#_=7.690390+12.700030</pre>		
1,643590-17.997190	b ^v =0.748466-i0.609952		
Error = 17.02%	Error - 13.49%	•	
For 17 configurations,	•		
First solution;	Second solution:		1
-5.123750+i2.254620	#_=1.443950-16.704520		•
-0.549476-17.03377	b ⁰ =6.370810-19.315060		
Error = 38.74%	Error = 60.77%		

CONCLUSIONS

Householder reflections appear to be far more efficient in solving linear problems arising within the field of unconstrained optimal synthesis of linkages. As to poplinear problems, the extension is straightforward. Regarding constrained problems, these could be handled using this method by introducing suitable slack variables and penalty functions. As to processor times, the first example consummed 11.8 sac, whereas the time reported (8) using Powell's method is 2.2 min, the method introduced here thus appearing to be more economical. With regard to the synthesis for rigid-body guidance, it is necessary to investigate whether for overdetermined problems, in general two different solutions can be expected, thus enabling the designer to syntabsize RRER plane linkages for overdetermined rigid-body guidance problems.

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DISERO CINEMATICO DE MAQUINARIA

· UNA GENERALIZACION DEL TEORENA: DE ARONHOLD - KENNEDY

_ DR. JORGE ANGELES

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JUNIO, 1984.

Una generalización del teorema de Aronhold-Kennedy

Jorge Angeles Alvarez*

INTRODUCCION

En el presente trabajo se hace una generalización del teorema de Aronhold-Kennedy, comúnmente conocido en los textos elementales de mecanismos como teorema de los tres centros.

Este teorema establece que cuando tres cuerpos están en movimiento, tienen tres centros instantáneos de movimiento relativo, los cuales se encuentran alineados. Carece de precisión, pues en el caso más general de movimiento de un cuerpo rígido no es posible hablar de centros instantáneos de movimiento relativo, quedando reducido este término al caso particular de movimiento plano, que no es el único que se presenta en problemas de ingeniería; baste citar el caso del movimiento de una junta universal o el de un tren de engranes cónicos, hipoidales o corona sinfín.

Surge entonces la necesidad de hacer una generalización de este teorema que comprenda al movimiento tridimensional, por lo que en este trabajo se establece la condición necesaria y suficiente para que exista movimiento de rotación pura entre los elementos de un mecanismo, y al lievar el teorama a tres dimensiones, se habla, en términos más generales, de ejes instantáneos de rotación.

Como en general, al diseñar un mecanismo, se requiere que este sea lo más eficiente posible (hay casos en que no importa tanto la eficiencia como la ventaja mecánica), es necesario evitar el deslizamiento de las superficies de contacto, el que es nulo cuando hay movimiento de rotación pura.

Existe otra generalización del teorema, que garantiza la existencia de ejes de deslizamiento mínimo y que es útil en el diseño de mecanismos en los cuales es imposible evitar el deslizamiento, como sucede en los acoplamientos hipoldales y en los corona sintín. Esta generalización se puede consultar en libros avanzados de mecanismos (raf 1),

En su forma més general, el teorema de Aronhold Kennedy establece que existen ejes de velocidad relativa mínima (llamados ejes de tornillo instantáneo) y que dichos ejes, para tres cuerpos régidos en movimiento, son tres y tienen la propiedad de ser normales a un eje común.

La originalidad del presente trabajo estriba en que se llege a la forma general del teorema utilizando vectores carteslanos, en lugar de métodos matriciales.

TEOREMA

Un cuerpo C, que gira con respecto a un cuerpo B, que a su vez está animado de un movimiento de rotación con respecto a un tercer cuerpo A, tiene un eje instantáneo de rotación con respecto a A si, y solo si, los ejes instantáneos de C con respecto a B, y de B con respecto a A son concurrentes; entonces, el eje instantáneo de C con respecto a A, también es concurrente con los dos anteriores.

Para su demostración, se dividirá el teorema en dos partes: en la primera se supondrá que existen los tres ejes, y se demostrará que son concurrentes; en la segunda, se consideraránque los dos primeros ejes son concurrentes, y, como consecuencia de ello, se demostrará que existe el tercer eje, y que es concurrente con los dos anteriores.

Obsérvese que como lo que interese es el movimiento relativo de los tres cuerpos, puede considerarse que uno de ellos tiene una velocidad arbitraria. Si por ejemplo, la velocidad angular det cuerpo A con respecto a un sistema de ejes inercial o nevitoniano es nuta, la velocidad angular de cualquier cuerpo referida a A es absoluta.

Según lo anterior, los ejes instantáneos de 8 y de C con respecto a A son conjuntos de puntos de velocidad absoluta nuía.

Sean: x-y-z, un marco newtoniano al que está tijo A; $\overline{w}_{B/A}$ $\overline{w}_{C/B}$ y $\overline{w}_{C/A}$ las velocidades angulares de B con respecto a A, C con respecto a B, y C con respecto a A, respectivamente. En la fig 1, se muestran los tres ejes instantáneos.

El punto Q está alojado sobre el eje instantáneo de C con respecto a A; por lo tanto, puede considerarse como punto de C o de A indistintamente, y tiene, en consecuencia, velocidad absoluta nula. El punto P está alojado sobre el eje instantáneo de C con respecto a B, por lo que puede considerarse alojado en B o en C indistintamente.

Los vectores de posición de P y de Q son, respectivamente, $\vec{r_p}$ y $\vec{r_q}$

Llamando \overline{r}_{12} al vector de posición de la intersección de los ejes 8 con respecto a A y C con respecto a A, \overline{r}_{13} al vector de posición de la intersección de los ejes 8 con respecto a A y C con respecto, a B, y \overline{r}_{23} a la intersección de los ajes C con respecto a A y C con respecto a B, la primera parte del teorema quedará demostrada al cumplirse las siguientes igualdades:



En efecto:

$$\overline{v}_q = \overline{v}_p + \overline{v}_{q/p}$$

Pero:

$$\overline{v}_{p} = \overline{w}_{B/A} \overline{r}_{p}$$

Y:

$$\overline{v}_{q/p} = \overline{w}_{C/B} \left(\overline{r}_{q} - \overline{r}_{p} \right)$$
(1)

Per lo tanto

$$\overline{v}_q = \overline{w}_{B/A} \overline{r}_p + \overline{w}_{C/B} \overline{r}_q - \overline{w}_{C/B} \overline{r}_p$$

Pero como Q es un punto de A, su velocidad absoluta es nula, por lo que

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 $\overline{w}_{B/A} \overline{r}_{a} + \overline{w}_{C/B} \overline{r}_{a} - \overline{w}_{C/B} \overline{r}_{a} = 0$

$$\widetilde{\mathbf{w}}_{\mathbf{C}/\mathbf{B}}\,\widetilde{\mathbf{r}}_{\mathbf{p}} = \widetilde{\mathbf{w}}_{\mathbf{B}/\mathbf{A}}\,\widetilde{\mathbf{r}}_{\mathbf{p}} + \,\widetilde{\mathbf{w}}_{\mathbf{C}/\mathbf{B}}\,\widetilde{\mathbf{r}}_{\mathbf{q}}$$

Multiplicando ambos miembros por $\overline{w}_{B/A}$, se tiene

$$\overline{w}_{B/A} \ \overline{w}_{C/B} \ \overline{r}_p = \overline{w}_{B/A} \ \overline{w}_{B/A} \ \overline{r}_p + \overline{w}_{B/A} \ \overline{w}_{C/B} \ \overline{r}_q$$
(2)

El segundo término del miembro a la derecha de la última expresión es nulo, por lo que

$$\overline{w}_{B/A} \ \overline{w}_{C/B} \ \overline{r}_{p} = \overline{w}_{B/A} \ \overline{w}_{C/A} \ \overline{r}_{q}, \tag{3}$$

Pero, de la ec 1:
$$\overline{v}_{q/p} = \overline{w}_{C/B} \overline{r}_{q/p}$$

por lo que $\overline{v}_q = \overline{w}_{B/A} \overline{r}_p + \overline{w}_{C/B} \overline{r}_{q/p} = 0$
Es decir $\overline{w}_{B/A} \overline{r}_p = -\overline{w}_{C/B} \overline{r}_{q/p}$

Multiplicando la expresión anterior por $\overline{w}_{C/A}$, y tomando en cuenta que $\overline{w}_{C/A} = \overline{w}_{C/B} + \overline{w}_{B/A}$:

$$(\vec{w}_{C/B} + \vec{w}_{B/A}) \vec{w}_{B/A} \vec{r}_p = - \vec{w}_{C/A} \vec{w}_{C/B} \vec{r}_{q/p}$$

Deserrollando el primer término, y eliminando el término que se anula

$$\overline{v}_{C/B} \overline{w}_{B/A} \overline{r}_{\rho} = -\overline{w}_{C/A} \overline{w}_{C/B} \overline{r}_{g/\rho}$$

Electuando una rotación no cíclica con los factores del segundo miembro de esta expresión para eliminar el signo negativo, se tiene

$$\overline{w}_{C/B} \ \overline{w}_{B/A} \ \overline{r}_{p} = \overline{w}_{C/D} \ \overline{w}_{C/A} \ \overline{r}_{q/p}$$
(4)

Pero

$$\overline{w}_{C/H} \, \overline{w}_{C/A} \, \overline{\tau}_{q/p} = \overline{w}_{C/B} \, \overline{w}_{C/A} \, \overline{\tau}_{q} - \overline{w}_{C/B} \, \overline{w}_{C/A} \, \overline{\tau}_{p}$$

Sustituyendo \overline{w}_{CFB} por su valor en función de las otras velocidades angulares, en el primer tármino del miembro a la derecha de la última expresión, se tiena

$$\widetilde{w}_{C/B} \, \widetilde{w}_{C/A} \, \widetilde{\tau}_{q/p} = (\widetilde{w}_{C/A} - \widetilde{w}_{B/A}) \, \widetilde{w}_{C/A} \, \widetilde{\tau}_{q} - \frac{1}{2} \widetilde{w}_{C/B} \, \widetilde{w}_{B/A} \, \widetilde{\tau}_{p}$$

Desarrollando el paréntesis, y eliminando el término que se anula, se tiene

$$\overline{w}_{C/B} \ \overline{w}_{C/A} \ \overline{r}_{q/p} = - \ \overline{w}_{B/A} \ \overline{w}_{C/A} \ \overline{r}_{q} - \overline{w}_{C/B} \ \overline{w}_{B/A} \ \overline{r}_{p}$$
(5)

Efectuando una rotación no cíctica con los factores del segundo término del miembro a la darecha de la ac 5, a fin de cambiar su signo,

Según la ec 3 el miembro de la derecha de esta última expresión es cero, por lo que

$$\overline{w}_{C/B} \ \overline{w}_{C/A} \ \overline{\tau}_{g/P} = 0 \tag{6}$$

De la ec 4 resulta

$$\overline{w}_{C/B} \ \overline{w}_{B/A} \ \overline{r}_p = 0 \tag{7}$$

De la ec 2,

$$\overline{w}_{B/A}/\overline{w}_{C/B}/\overline{r}_q=0$$

Sustituyendo en esta última expresión $\overline{w}_{C/B}$ por su valor $\overline{w}_{C/A} = \overline{w}_{B/A}$, se tiene

$$\overline{w}_{B/A} \ \overline{w}_{C/A} \ \overline{r}_q = \overline{w}_{B/A} \ \overline{w}_{B/A} \ \overline{r}_q = 0$$

Pero el segundo término de esta expresión se anula, por lo que

$$\overline{w}_{B/A} \, \overline{w}_{C/A} \, \overline{\gamma}_{A} = 0 \tag{8}$$

De las ecs 6, 7 y 8, se deduce que los ejes se intersecan dos a dos. Falta ahora demostrar que se intersecan en un punto común. Para esto, basta verificar que el vector de posición es el mismo para las tres intersecciones; es decir, $\vec{r}_{12} = \vec{r}_{13} = \vec{r}_{23}$.

En la notación anterior, se sigue la misma nomenclatura que se expuso al enunciar el teorema,

Previamente se determinará la expresión que define a cada uno de los vectores de posición anteriores, para lo cual debe recurrírse a la fig 2.

Apàrecen dos ejes que se intersecan en el punto I. Estos ejes están determinados por los puntos :L y M, dados a su vez por sus vectores de posición \overline{r}_L y \overline{r}_M respectivamente, y por los vectores unitarios que dan su dirección/ \overline{L} y \overline{m} . El problema es daterminar el vector \overline{r}_L .

$$\overline{r}_{L} = \overline{r}_{L} - \overline{r}_{L}$$

 \overline{r}_{L} es dato; $\overline{r}_{L1} = -$

El segmento L1 se determina por el teorema de los senos, a partir de la fig

$$\frac{LI}{\text{sen B}} = \frac{LM}{\text{sen A}} ; \text{ de donde LI} = \frac{\text{sen B}}{\text{sen A}} LM$$

LIĨ

La longitud del segmento LM es dato, pues está dada por los vectores de posición de L y de M, que también son datos; los senos de A y de 8 se determinan de la siguiente manera:

$$sen B = \frac{\left|\vec{r}_{L,M} \ \vec{m}\right|}{\left|\vec{r}_{L,M}\right|}; sen A = \left|\vec{T} \ \vec{m}\right|$$

$$Por lo tanto LI = \left|\vec{r}_{L,M}\right| \frac{\left|\vec{r}_{L,M} \ \vec{m}\right|}{\left|\vec{r}_{L,M}\right| \left|\vec{T} \ \vec{m}\right|} = \frac{\left|\vec{r}_{L,M} \ \vec{m}\right|}{\left|\vec{T} \ \vec{m}\right|}$$

$$\vec{r}_{LI} = -\frac{\left|\vec{T}_{L,M} \ \vec{m}\right|}{\left|\vec{T} \ \vec{m}\right|} T$$

$$De donde \vec{r}_{I} = \vec{r}_{L} + \frac{\left|\vec{T}_{L,M} \ \vec{m}\right|}{\left|\vec{T} \ \vec{m}\right|} T$$

Esta expresión servirá para determinar el vector de posición de cada intersección.

En la fig 3 aparece nuevamente la disposición de los tres ejes instantáneos a que se hace mención, mostrándose las intersecciones de cada par de ejes y los vectores de posición de cada una de estas.

Los vectores unitarios $\overline{T}_1,\overline{T}_2,\gamma\overline{T}_3$ están dados por

$$\frac{\overline{w}_{B/A}}{\left|\overline{w}_{B/A}\right|} = \frac{\overline{w}_{C/A}}{\left|\overline{w}_{C/A}\right|} = y - \frac{\overline{w}_{C/B}}{\left|\overline{w}_{C/B}\right|} = respectivamente.$$

De la expresión que da r₁, se tiene:

$$\overline{r}_{12} = -\frac{\left|\overline{r}_{q} \overline{r}_{2}\right|}{\left|\overline{1}_{l} \overline{r}_{2}\right|} T_{1} = -\frac{\left|\overline{r}_{q} \overline{w}_{C/A}\right|}{\left|\overline{w}_{C/A}\right| \left|\frac{\overline{w}_{B/A} \overline{w}_{C/A}\right|}{\left|\overline{w}_{B/A} \overline{w}_{C/A}\right|} \frac{\overline{w}_{B/A}}{\left|\overline{w}_{B/A} \overline{w}_{C/A}\right|}$$





$$\overline{r}_{12} = -\frac{\left|\overline{r}_{q} - \frac{\overline{w}c/A}{\left|\overline{w}_{B/A} - \overline{w}c/A\right|}\right|}{\left|\overline{w}_{B/A} - \overline{w}c/A\right|} \quad \overline{w}_{B/A}$$
(9)

$$\widehat{\mathbf{r}}_{1,3} = -\frac{\left|\overline{\mathbf{r}}_{\mathbf{p}} | \mathbf{T}_{3}\right|}{\left|\widehat{\mathbf{t}}_{1} | -\overline{\mathbf{T}}_{3}\right|} | \overline{\mathbf{t}}_{1} = -\frac{\left|\overline{\mathbf{r}}_{\mathbf{p}} | \overline{\mathbf{w}}_{C/B}\right|}{\left|\overline{\mathbf{w}}_{B/A} | \overline{\mathbf{w}}_{C/B}|} \frac{\overline{\mathbf{w}}_{B/A}}{\left|\overline{\mathbf{w}}_{B/A}\right|}$$

Simplificando

$$\overline{\tau}_{12} = -\frac{\left|\overline{\tau}_{p} - \overline{w}_{C/B}\right|}{\left|\overline{w}_{B/A} - \overline{w}_{C/B}\right|} \qquad (10)$$

$$\overline{\tau}_{2,3} = \overline{\tau}_q - \frac{\left|\overline{\tau}_{qp} T_{3}\right|}{\left|\overline{T_2} - \overline{\tau}_{3}\right|} T_2 = \overline{\tau}_q - \frac{\left|\overline{\tau}_{qp} \overline{w}_{C/B}\right|}{\left|\overline{w}_{C/B}\right| \left|\frac{\left|\overline{w}_{C/A} - \overline{w}_{C/B}\right|}{\left|\overline{w}_{C/A} - \overline{w}_{C/B}\right|} - \frac{\overline{w}_{C/A}}{\left|\overline{w}_{C/A}\right|}$$

de donde

$$\overline{r}_{23} = \overline{r}_{q} - \frac{\left|\overline{r}_{qD} \ \overline{w}_{C/B}\right|}{\left|\overline{w}_{C/A} \ \overline{w}_{C/B}\right|} \quad \overline{w}_{C/A}$$
(11)

Primero se demostrará que $\overline{r}_{12} = \overline{r}_{13}$

En efecto
$$\overline{r}_q \ \overline{w}_{C/A} = \overline{w}_{C/A} \ \overline{r}_q \ , \ \gamma$$

 $\overline{r_p} \ \overline{w}_{C/B} = \overline{w}_{C/B} \ \overline{r}_p$

En estas expresiones, por comodidad se ha invertido el sentido de los factores de los productos vectoriales, inversión que resulta de un cambio de signo al producto; pero como se manejan módulos, el cambio de signo no afecta a éstos y la manipulación es justificaga.

Entonces

$$\begin{split} \overline{w}_{C/B} \, \overline{r}_p &= (\overline{w}_{C/A} - \overline{w}_{B/A}) \ (\overline{r}_q + \overline{r}_{qp}) = \\ &= \overline{w}_{C/A} \, \overline{r}_q + \overline{w}_{C/A} \, \overline{r}_{qp} - \overline{w}_{D/A} \, \overline{s}_q - \overline{w}_{B/A} \, \overline{r}_{qp} = \\ &= \overline{w}_{C/A} \, \overline{r}_q + \overline{v}_p - \overline{v}_q - \overline{w}_{B/A} \, \overline{r}_{qp} \end{split}$$

Pero $\overline{v}_{ij} = 0$, como ya se vio en un principio. De donde

$$\vec{w}_{C/D} \, \vec{r}_p = \vec{w}_{C/A} \, \vec{r}_q - \vec{v}_p - \vec{w}_{B/A} \, (\vec{r}_p - \vec{r}_q)$$
$$= \vec{w}_{C/A} \, \vec{r}_q - \vec{v}_p - \vec{v}_q$$

Observese que $\overline{w}_{B/A}$ $\overline{r}_{p} = \overline{v}_{p}$, y que $\overline{w}_{B/A}$ $\overline{r}_{q} = \overline{v}_{q} = 0$

Simplificando, se obtiene

 $\overline{W}_{C/B} \overline{T}_{p} = \overline{W}_{C/A} \overline{T}_{q}$, y al tomar módulos

$$\overline{w}_{C/B} \overline{r}_{p} = \overline{w}_{C/A} \overline{r}_{q}$$
(12)

Por otra parte $\overline{w}_{B/A}$ $\overline{w}_{C/B} = \overline{w}_{B/A}$ $(\overline{w}_{C/A} - \overline{w}_{B/A})$

de donde
$$\overline{w}_{B/A} \ \overline{w}_{C/B} = \overline{w}_{B/A} \ \overline{w}_{C/A}$$
 (13)

Dividiendo (12) entre (13), resulta

$$\frac{\overline{r}_{q}}{\overline{w}_{B/A}} \frac{\overline{w}_{C/A}}{\overline{w}_{C/A}} = \frac{\overline{r}_{p}}{\overline{w}_{B/A}} \frac{\overline{w}_{C/B}}{\overline{w}_{C/B}} \text{ con lo que se demuestra que}$$

Llamando I a esta intersección común, se tiene que i está en los ejes 1 y 2, al mismo tiempo que en los ejes 1 y 3, por lo que, i está en los ejes 2 y 3; por lo tanto

$$\overline{r}_{23}=\overline{r}_{13}=\overline{r}_{13}$$

De lo anterior se concluye que si existan los tres ejes, estos concurren en un punto común, con lo cuel se ha demostrado la primera parte del teorema.

Faita demostrar la segunda parte, en la que se supone que dos de los ejes se intersecan en un punto, y que como consecuencia de elto existe el tercer eje, damostrándose además, que este es concurrente con los dos anteriores.

Supóngase que los ejes de 8 con respecto a A, y de C con respecto a B, se intersecan en un punto que, por comodidad, se considerará como el origen de coordenadas. Sea x-y-2 un marco newtoniano, con origen en la intersección de los dos ejes. Debe demostrarse que exíste un conjunto de puntos con velocidad absoluta nuta que están en el cuerpo C y que en consecuencia, constituyen el eje instantáneo de C con respecto a A.

Sea P un punto cualquiera de C. Su velocidad absolute con respecto a la velocidad de un punto de B, como Q por ejemplo, será

$$\overline{\mathbf{v}}_{p} = \overline{\mathbf{v}}_{q} + \overline{\mathbf{v}}_{p/q}$$
$$\overline{\mathbf{v}}_{q} = \overline{\mathbf{w}}_{B/A} \cdot \overline{\mathbf{r}}_{q}, \mathbf{Y}$$
$$\overline{\mathbf{v}}_{p/q} = \overline{\mathbf{w}}_{C/B} \cdot \overline{\mathbf{r}}_{qp}$$

Entonces $\overline{v}_{p} = \overline{w}_{B/A} \overline{\tau}_{q} - \overline{w}_{C/B} \overline{\tau}_{qp}$

$$\begin{split} \overline{v}_{p} &= \overline{w}_{B/A} \, \overline{r}_{q} + \overline{w}_{C/B} \, \overline{r}_{qp} = \overline{w}_{B/A} \, \overline{r}_{q} + \overline{w}_{C/B} \, (\overline{r}_{p} - \overline{r}_{q}) \\ \\ \overline{v}_{p} &= \overline{w}_{B/A} \, \overline{r}_{q} + \overline{w}_{C/B} \, \overline{r}_{p} - \overline{w}_{C/B} \, \overline{r}_{q} \end{split}$$

Pero de la fig 4, $T_q = m \overline{w}_{C/B}$, en que m es un escelar cualquiera que hace que el módulo de T_q sea m veces el de $\overline{w}_{C/B}$.



Entonces $\overline{w}_{C/B}$ $\overline{r}_q = \overline{w}_{C/B}$ m $\overline{w}_{C/B}$, que es evidentemente cerro. Por tanto

$$\overline{v}_{p} = \overline{w}_{B/A} \overline{r}_{q} + \overline{w}_{C/B} \overline{r}_{p}$$

Para que \overline{v}_p sea cero como condición para que sea eje instantáneo de rotación de C con respecto a A, se requiere que

 $\overline{w}_{\rm B/A} \, \overline{r}_{\rm a} + \overline{w}_{\rm C/B} \, \overline{r}_{\rm b} = 0$

Es decir, se requiere que

$$\overline{w}_{a/A} = \overline{w}_{C/B} + \overline{w}_{C/B} \overline{\tau}_{a} = 0$$

Esta última expresión puede ponerse en la forma

$$\overline{W}_{C/B} \overline{T}_p - m \overline{W}_{C/B} \overline{W}_{B/A} = 0$$

Sacando como factor común a W_{C/B}:

$$\overline{w}_{C/B} \left(\overline{r}_{p} - m \, \overline{w}_{B/A} \right) = 0$$

igualdad que se cumple si $\overline{r}_p - m \overline{w}_{B/A} = 0$, o sea, si $\overline{r}_p = m \overline{w}_{B/A}$, lo cual quiere decir que P, punto de C, está alojado sobre el eje instantáneo de rotación de B con respecto a A. Esto sería un caso trivial, pues entonces C y A serían el mismo cuerpo. También se cumple la igualdad anterior si $\overline{r}_p - m w_{B/A} = n w_{C/B}$, en que n es un escalar. Se tendría entonces que

$$\overline{T}_{\mathbf{r}} = \mathbf{m} \, \overline{\mathbf{w}}_{\mathbf{U}/\mathbf{A}} + \mathbf{n} \, \overline{\mathbf{w}}_{\mathbf{C}/\mathbf{B}}.$$

De esta manera, todos los puntos como P, en que sus vectores de posición \overline{r}_p son línealmente dependientes con $\overline{w}_{B/A} = \gamma - \overline{w}_{C/B}$, cumplen con la condición de tener velocidad absoluta nula. Es decir, constituyen el eje instantáneo de rotación de C con respecto a A. Además se observa que cuando m y n son simultáneamente nulas, el lugar geométrico pasa por la intersección de los dos primeros ejes, que es lo que se quería demostrar.

COROLARIO. Los tres ejes son coplanares y guardan la relación: $\overline{w}_{C/A} = \overline{w}_{C/B} + \overline{w}_{B/A}$, y como se demostró que los, ejes concurren en un mismo punto, se desprende que son coplanares.

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DISEÑO CINEMATICO DE MAQUINARIA

AUTOMATIC COMPUTATION OF THE SREW PARAMETERS OF RIGID-BODY HOTIONS PART. I: FINITELY - SEPARATED POSITIONS

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JORGE ANGELES

JUNIO, 1984

AUTOMATIC COMPUTATION OF THE SCREW PARAMETERS OF RIGID-BODY MOTIONS.

PART. I: FINITELY - SEPARATED POSITIONS

Jorge Angeles¹

Abstract

A novel approach, based on invariants, is introduced, that leads to efficient algorithms for computing the screw parameters of rigidbody motions. Both finitely and infinitesimally-separated positions are treated. The computer implementation of the algorithm allows the real-time computation of the parameters defining the position and orientation of a rigid body.

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Introduction

Robotics applications, calling for efficient algorithms for the determination of the position and orientation of a rigid body from a reduced set of measurements, have motivated current research on this topic [1-4] . Although this subject is well known from a theore<u>t</u> ical standpoint [5,pp 1-25, 6, 7 & 8, pp 35-62], the need of means for the efficient real-time computation of the parameters defining a rigid-body motion, commonly referred to as the screw parameters, has called for a revisitation of the underlying theoretical basis. In fact, as pointed out in [2, pp85-118]generally acepted formulae can fail to apply under special, though rather frequent, circumstances.

Presented in [2] are algorithms that take into account all possible particular cases. These algorithms, however, are rather lengtby, and lack symmetry, in the sense of considering one particular point as a body-fixed reference. The latter item is disadvantageous in applications, as pointed out in [3]. The approach introduced in [3] solves the problem of lack of symmetry, but introduces spurious singularities; it is, additionally, limited to the infinitesimally-separated-positions case. Introduced in the present paper is an algorithm based upon invariant concepts that allow a fast and reliable computation of the accrew parameters, for finitely-separated positions. Infinitesimally-separated positions are discussed in an accompanying paper. Given the reduced number of operations involved, this algorithm can be applied to the real-time computation of the said parameters, as required in robotics applications. The procedure, however, is based upon an exact

knowledge of the coordinates of three noncollinear points in two distinct configurations of the rigid body which they belog to. In practice, such coordinates are known only up to random measurement errors. Thus, filtering off of these errors requires either taking measurements of over three points of the rigid body or perform redundant computations, as outlined in the paper. The principles presented here, nevertheless, can be applied even if the computations are based upon measurements of over three points. This subject, however, is not discussed here, but only proposed for further research.

Description of the algorithm

The motion associated with two finitely-separated positions of a rigid body is fully described by the following [2,pp 85-119] : a) the axis of the screw, given by the three coordinates of one of its points (preferably the one lying closest to the origin) and three direction cosines, b) the sliding of the screw along its axis, and c) the angle of rotation about the axis of the screw, supplied with sign, given a positive direction defined on the axis. The set of scalar screw parameters of the rigid-body motion is, thus, the following: the three components of a vector \mathbf{r}_0 , locating point \mathbf{R}_0 of the screw axis *l*, that lies the closest to the origin; the three components of a unit vector \mathbf{e} , paralel to *l* and defining the positive direction along *l*; two scalars, u and θ , representing the sliding along and the rotation about *l*. This gives 8 scalar components, which are subject to the following two scalar constraints:

$$e^{T}e = 1$$
 (1a)
 $r_{g}^{T}e = 0$ (1b)

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the superscript $\binom{T}{}$ standing for transposition. Thus, the number of independent screw parameters is six. The computation is based upon that of the orthogonal matrix defining the rotation involved. The latter is based, in turn, upon the computation of the principal directions of the second-moment tensor of the three points, about their centroid. This is equivalent to the moment of inertia of a rigid system composed of three unit-mass particles. Let p_i be the position vector of the \pounds^{th} point, and c, that of their centroid. Hence, c is given by

$$c = \frac{1}{3} + \frac{3}{1} p_{i}$$
 (2)

whereas its second-moment tensor with respect to its centroid, by (9,pp 393-397) 3 2

$$= \sum_{j} \left(\rho_{ij}^{2} - \rho_{ij} \rho_{ij}^{T} \right)$$
 (3a)

with

$$p_{\mathcal{L}} \equiv p_{\mathcal{L}} - c_{\mathcal{L}}$$
 (3b)

The second-rank tensor I is invariant, symmetric and positive definite. The last two properties are obvious from definition (3a). Invariance, on its behalf, means that, under a change of coordinates the three proper values of I do not change, its proper vectors both in the original configuration, $\{e_1, e_2, e_3\}$, and those in the new one, $\{f_1, f_2, f_3\}$, being related by

$$f_{i} = Q e_{i}$$
 (4)

where Q is the matrix associated with the rotation involved. The foregoing is illustrated in Fig. 1. Horeover, if A denotes the 3 x 3 matrix containing the components of I in the original configuration, whereas



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Fig 1 Proper vectors of the second-moment tensor of a three-poin rigid system

B, those in the new one, then

$B = Q A Q^T$

Clearly, all matrices appearing in (5) should be expressed in the same coordinate frame. Now, if the three given points are noncollinear, it is a simple matter to verify that tensor I, and hence either matrix A or B, is nonsingular. The aforementioned matrix Q is, of course, prop er orthogonal.

The algorithm described next requires the computation of the three proper values and vectors of tensor I. Although this computation leads, in general, to the solution of a nonlinear problem which has to be solved iteratively, for the problem at hand a direct solution is possible, as shown next. Furthermore, once the foregoing eigenvalue problem has been solved, tensor I is expressed with respect to its proper vectors, indentified in what follow with the two orthonormal triplets $\{e_1, e_2, e_3\}$ and $\{f_1, f_2, f_3\}$, of eq (4). Given these two triplets, the computation of Q is a simple matter. The procedure to compute the three proper values and vectors of Q directly(as opposed to iteratively) is, thus, fundamental to this algorithm, for which reason the said procedure is described first.

Let all coordinates be given with respect to a hegerence frame, hence forth referred to as X, Y, Z. Moreover, let A and B be the matrices representing I with respect to the reference frame in the original and the new configuration of the rigid body, respectively. Since tensor I is associated with a plane body, the one defined by the three given points, two of its proper vectors lie in the plane of the body, the remaining one being perpendicular to this plane, as shown in elementary mathematics and mechanics texts. Furthermore, the proper value associated with the third vector, I_3 , equals the sum of the proper values, I_1 and I_2 , asso cjated with the first two proper vectors. Thus,

(6)

6)

(5)

Next the first two invariants of I, tr (I) and tr (I^2), also called the two first moments of I [20, p 67] are expressed both in terms of the proper values of I, and in terms of matrix A, and then equated, which leads to the following system of equations for I₁ and I₂:

$$2(1_1 + 1_2) = tr A$$
 (7a)

$$I_1^2 + I_2^2 + (I_1 + I_2)^2 = \operatorname{tr} A^2$$
 (7b)

with a similar set of equations for matrix B which, if no roundoff nor measurement errors were present, would be identical. If these are present, then both sets can be used to filter the said errors. Since (7a & b) represent a system of two independent equations for two unknowns, I_1 and I_2 , they suffice to compute these:

Elimination of I_2 between eqs (7a & b) produces the following quadratic equation for I_1 :

$$4I_{1}^{2} - 2(trA)I_{1} + (tr^{2}A - 2trA^{2}) = 0$$
 (8)

whence

$$I_{1} = \frac{\text{trA} + \sqrt{\partial \text{tr} A^{2} - 3 \text{tr}^{2} A}}{4}$$

Assuming that the three proper values of I are ordered such that

$$\mathbf{I}_{1} \leq \mathbf{I}_{2} \leq \mathbf{I}_{3} \tag{9}$$

then

$$I_{1} = \frac{trA - \sqrt{8} tr A^{2} - 3 tr^{2} A}{4}$$
(10a)

From eq (7a) it is readily realized that the second root of eq (8) equals I_{η} , i.e.

 $(\overline{\mathbf{r}})$

$$I_2 = \frac{\text{tr } A + 8 \text{ tr } A^2 - 3 \text{ tr}^2 A}{4}$$
 (10b)

and hence

$$I_{3} = \frac{1}{2} tr A$$
 (10c)

In order to compute the proper vectors of I, two possibilities are considered, namely either $I_1 < I_2$ or

I1- I2 (noturally, to machine precision).

In the first case, the null space of $A - I_{i}$, for i = 1, 2, 3, 1 equals the 3 x 3 identity matrix, is of dimension 1. Thus, any vector f the said space can be normalized to produce e_{i} , the problem thus reducing to the determination of this space. This is most efficiently done with the aid of Householder reflections [11, pp 111 - 118], which reduce matrix A to a row echelon form. That is to say, if H is the product of the three involved reflections, then

$$H (A - I_{i} I) = \begin{bmatrix} a I \\ . I \\ .$$

with

$$a_{1} = [a_{11}, a_{12}, a_{13}]^{T}$$
 (11b)
 $a_{2} = [0, a_{21}, a_{22}]^{T^{-1}}$ (11c)

$$a_3 = [0, 0, 0]^T$$
 (11d)

Hence,

$$\mathcal{L} = \frac{a_1 \times a_2}{\left| \left| a_1 \times a_2 \right| \right|}$$
(11e)

In practice the foregoing computations need be executed only for two proper vectors, the remaining one being computed simply as the cross product of the two previously computed ones.

Now, if $I_1 = I_2$, this means that the three-particle rigid system has a cylindrically symmetric inertia tensor. Since each particle has been assumed of unit mass, this can only happen if the three particles are located at the vertices of an equilateral triangle. Hence, any vector lying in the plane defined by the three points is a proper vector of I. This means that e_1 and e_2 can be chosen arbitrarily within that plane, though mutually orthogonal. In order to uniquely define these vectors, e_1 can be chosen, for instance, as

 $e_{j} \frac{p_{j} - c}{|| p_{j} - c||}$ (12a)

Since e_3 is uniquely defined perpendicular to the plane of the given points, e_7 can be readily computed as

 $\mathbf{e}_{2} = \mathbf{e}_{3} \times \mathbf{e}_{1} \tag{12b}$

Notice that symmetry is not destroyed by the fact of defining arbitrarily vector e_1 as appearing in eq (12a), for the inertia tensor itself is cylindrically symmetric, as said previously.

The set $\{f_{\lambda}\}_{1}^{3}$ should be computed correspondingly, i.e. for the case $\lambda_{1} < \lambda_{2} < \lambda_{3}$, f_{λ} should correspond to λ_{λ} ; for the second case, $\lambda_{1} = \lambda_{2}$, if e_{1} , and e_{2} are computed as given by eqs. (12 a &b), then

$$f_{j} = \frac{p_{j}^{2} - c'}{\{j p_{j}^{2} - c'\}}$$
(13a)

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and

fg= fg x fj

 p_{λ}^{+} (i= 1,2,3) and c' being the position vectors of the given points and their centroid, respectively, in the new configuration.

The entries of matrix Q, q_{ij} (*i*, *j* = 1,2,3), representing the rotation in the $\{e_j\}_{j=1}^{2}$ basis, are now computed by simply recalling the definition of the matrix representation of a linear transformation [12, p 65]. Thus, if both $\{e_{i}\}_{i=1}^{3}$ and $\{f_{i}\}_{i=1}^{3}$ are given in referenceframe coordinates, then

$$q_{ij} = e_{i}^{T} f_{j} \qquad (14a)$$

The rotation expressed in reference-frame coordinates, referred to as matrix Q_R, is obtained as

$$Q_{\rm R} = EQE^{\rm T}$$
 (14b)

with

 $\mathbf{E} = \begin{bmatrix} \mathbf{e}_1 & \mathbf{e}_2 & \mathbf{e}_3 \end{bmatrix}$ (14c)

Parameters e, in reference-frame coordinates, and 0 can now be obtained as

 $e \sin \theta = vect (Q_R), \cos \theta = \frac{1}{2} tr (Q) -1$ (15a)

Now, denoting by q_{i} the *i*-th component of vect (Q), the foregoing invariants are given by

> $q_{j} = \frac{1}{2} \epsilon_{ijk} q_{kj} tr (Q) = q_{jj}$ (15b)

and so

$$e \sin \theta = E c$$
 (15c)

The index convention has been used in eq (15b), ϵ_{ijk} being the alternating tensor commonly used in tensor analysis. The remaining screw parameters are now computed: Line L is defined here by the position vector \mathbf{r}_0 of one of its points, R_0 , whose distance to the origin (of the reference frame previously introduced) is a minimum, and by vector \mathbf{e} , giving its direction [2, pp 85-119]. Vector \mathbf{r}_0 can be computed, in turn, using either a closed-form expression or the minimum-norm solution to an underdétermined linear algebraic system. In the first case, the formula is [2,p 91] :

$$r_0 = \frac{1}{2} \cot \frac{\theta}{2} \exp (p_1^* - p_1^*) - \frac{1}{2} \exp (p_1^* + p_1^*)$$
 (16)

with similar expressions for vectors P_2 , P_2' , P_3 and P_3' . The three resulting formulae are, of course, redundant and, if no measurement nor roundoff errors were present, all three would yield one and the same vector r_0 . Since such errors are always present, the said three formulae do in fact produce slightly different results. The involved error can be filtered by taking the mean of the three computed values. Alternatively, vector r_0 can be computed as the solution of a linear algebraic system. Indeed, any point R, of position vector r, lying on L satisfies the following equation (2, p 89):

 $(Q - 1)^{T} (Q - 1) r = (Q - 1)^{T} (Qp_{i} - p_{i})$ (17) which is valid for i = 1, 2, 3. None of the three equations (17) can be solved for r, however, for matrix Q - 1, and hence $(Q-1)^{T}(Q - 1)$, is singular. In fact, were this matrix nonsingular, then eq(17) would define not a set of points R of L, but one single point. Eq (17) can now be

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expressed in a more compact form as

 $A \equiv (Q-1)^T (Q-1), b = (Q-1)^T (Q p_{\chi} - p_{\chi})$ (18b)

The point of L whose distance to the origin is a minimum is now ...selved via the following minimization problem

$$\mathbf{A}^{t}\mathbf{r} = \mathbf{b}^{1} \tag{20}$$

with A' being a full-rank 2 x 3 matrix. Matrix A' can be readily obtained from (18a) if it is taken into a row echelon form, as discussed before, via Householder reflections. This would imply a corresponding transformation of vector b of eq (18a), for which reason, the righthand side of eq (20) is changed to b'.

(20). The solution to this problem, r_o, is given in closed form by the matrix A', namely as

$$T_{0} = (A^{\dagger})^{\dagger} b^{\dagger}$$
 (21a)

. with

 $(A^{\prime})^{\dagger}$ $(A^{\prime})^{T} [A^{\prime} (A^{\prime})^{T}]^{-1}$ (21b) Computing $(A^{\prime})^{\dagger}$ explicitly, as given by eq (21b), however, is not recommended, for the matrix in brackets is usually ill-conditioned;

(12)

in fact, its condition number is the square of that of matrix A', which is a number larger than 1 [13, p 223]. A safe means of computing r_0 from eq (20) is now outlined, as recommended by Lawson and Hanson [14, pp 74-76]. Let H be the product of the two Householder reflections reducing (A')^T to row echelon form. Since matrix H is orthogonal, H^TH = 1, and hence, eq (20) can be rewritten as

Let

 $H(A')^{T} = T, Hr = y$

(23a)

(22)

with T defined as

$$\mathbf{T} = \begin{bmatrix} \mathbf{U} \\ -\frac{\mathbf{U}}{\mathbf{0}} \end{bmatrix}_{\frac{1}{2}}^{+}, \quad \mathbf{y} = \begin{bmatrix} \mathbf{y} \\ -\frac{\mathbf{J}}{\mathbf{y}} \end{bmatrix}_{\frac{1}{2}}^{+}$$
(23b)

the 2 x 2 matrix U being upper triangular. Eq (22) can thus be expressed $\frac{1}{100}$ as

$$\begin{bmatrix} \mathbf{u}^{\mathrm{T}} & \mathbf{0} \end{bmatrix} \begin{bmatrix} \mathbf{y}_{1} \\ -\mathbf{1} \\ \mathbf{y}_{2} \end{bmatrix} = \mathbf{b}^{\mathrm{T}}$$
(24a)

or, alternatively, as

$$u^{T}y_{1} + 0y_{2} = b^{1}$$
 (24b)

Eq (24b) can be readily solved for y_j by forward substitution, for U^T is lower triangular, y_2 being left undefined so far. However, since y_2 is being multiplied times a zero matrix, it can be given any real value. Since the minimum norm solution to eq (22) is being sought, then

the proper choice of y, is

$$y_{j} = 0$$
 (25)

Thus, the minimum-norm solution to eq (24b), yo, is given as

$$\mathbf{y}_{0} = \begin{bmatrix} \mathbf{U}^{-T} - \mathbf{b}_{1}^{T} \\ \mathbf{0} \end{bmatrix}$$
(26)

Vector r, is now computed simply as

$$\mathbf{r}_{o} = \mathbf{H}^{\mathrm{T}} \mathbf{y}_{q} \tag{27}$$

Again, each value of Z produces one corresponding system of equations (17), or, alternatively (18a). All three systems can be solved for r_0 , which would produce three slightly different values for this unknown, out of which the mean value would be the one containing the minimum error.

It is pointed out here that the foregoing computations, simple as they are, are sought to be executed with the highest efficiency, for they are aimed at applications demanding real-time computations, while keeping the roundoff error low enough. Finally, the single parameter that is to be computed is u, the sliding of the arew. This is readily computed as the projection on L of the displacement undergone by any of the three given points, i. e. as

$$u = e^{T} \left(p_{\vec{\lambda}}^{1} - p_{\vec{\lambda}} \right)$$
 28)

which is valid, again, for all three points. Hence, the likeliest u can be taken as the mean value of those three, thereby completing the computation of the parameters sought.

In the foregoing discussion the only source of singularities is tensor I. In fact, if the three given points are collinear, then the

(H)

orientation of the body about the line defined by the three points is undefined. Practical applications, rather frequent, call for the determination of the position and orientation of a line of a body, and not of the body itself. Such applications arise, for instance, in the positioning of axially-symmetric workpieces or tools. The problem of the determination of the screw parameters for such a situation and others, referred to as incompletely-specified motions, has been studied in [15], whereas the motion of a rigid line in two infinitesimally-separated pom itions only, has been treated in [16].

Apart from the singular cases mentioned, ill- conditioned problems should be considered, as well. These arise whenever the ratio I_2/I_1 becomes very large, as compared to unity. Such a ratio becomes large when the legs of the triangle defined by the three different points have very different lengths. Hence, the more such a triangle "approaches" an equilateral one, the better conditioned is the problem. Ill- conditioned problems may lead to cancellations in the computation of I_1 , as given by eq (10a), for which reason the computation of I_1 and I_2 should be executed in the following order, according to [17,pp20-23]: I_2 is first computed with eq (10b); next I, is computed as

$$I_{1} = \frac{tr^{2}A - 2tr A^{2}}{I_{0}} \qquad (29)$$

An example is next included, that illustrates the foregoing proce dure. (5

Example

Determine the screw parameters of the motion of a rigid body whose original and final configurations, referred to as C and C', respectively, are given by the coordinates of three of, its points, P_1 , P_2 and P_3 in C, and P_1 , P_2 and P_3 in C¹. The position vectors of these points are given as: $p_{f} = [1, 0, 0]^{T}, \qquad p_{f}^{T} = [2, 0, -1]^{T}$ P_2^{T} [1, 1, 0] T, P_2^{T} = [2, 0, 0] T $p'_3 = [3, -1, 0]^T$ $p_{3} = [2, 1, -1]^{T}$ Hence. $c = \frac{1}{3} [4, 2, 1]^{T}, c' = \frac{1}{3} [7, -1, -1]^{T}$ c and c' being the position vectors of the centroids in and '', respectively. Vectors p_i , p'_i , defined as $p_i = p_i - c$, and $p'_i = p'_i - c'_i$. ∠ = 1, 2, 3, are thus $p_{1} = \frac{1}{3} [-1, -2, 1^{T}] p_{1}' = \frac{1}{2} [-1, 1, 2]^{T}$ $p_{g} = \frac{1}{3} [-1, 1, 1, T] p_{g} = \frac{1}{3} [-1, 1, 1]$ $P_3 = \frac{1}{3} [2, 1, -2^T] P_3' = \frac{1}{3} [2, -2, 1]^T$ Matrices A and B are thus computed as $A = \frac{1}{3} \begin{bmatrix} 4 & -1 & 2 \\ -1 & 4 & 1 \\ 2 & 1 & 4 \end{bmatrix} , B = \frac{1}{3} \begin{bmatrix} 4 & 2 & -1 \\ 2 & 4 & 1 \\ -1 & 1 & 4 \end{bmatrix}$ · Hence, trA = 4, $tr A^2 = 20/3$

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Eqs (10a - c) thus yield

$$I_1 = 1 - \frac{\sqrt{3}}{3}, \quad I_2 = 1 + \frac{\sqrt{3}}{3}, \quad I_3 = 2$$

Now e₃, the unit vector spanning the null space of A- I₃1, is determined. This matrix is $A-I_3 = \frac{1}{3} \begin{bmatrix} -2 & -1 & 2 \\ -1 & -2 & 1 \\ 2 & 1 & -2 \end{bmatrix}$

which can be taken to the following row echelon form via Householder reflections: $\begin{bmatrix} 3 & 2 & -3 \end{bmatrix}$

$$H (A - I_3 I) = \frac{1}{3} \begin{bmatrix} 0 \sqrt{2} & 0 \\ 0 & 0 & 0 \end{bmatrix}$$

Thus, letting

$$a_1 = [3, 2, -3^T] a_2 = [0, \sqrt{2}, 0]^T$$

in eq (11c), one readily obtains

$$e_2 = [0.7071, 0, 0.7071]^T$$

Similarly,

and

$$e_1 = e_2 \times e_3 = [-0.6280, -0.4597, 0.6280]^T$$

Similar results are obtained for matrix B as follows: $f_1 = \begin{bmatrix} -0.6280, 0.6280, -0.4597 \end{bmatrix}^T$

Hence, the rotation matrix Q is obtained from formula (14) as

$$Q = \begin{bmatrix} -0.1833 & -0.6124 & 0.7691 \\ -0.6124 & 0.6833 & 0.3980 \\ -0.3980 & 0.5000 \end{bmatrix}$$

and
$$\begin{bmatrix} 0 & 0 & -1 \\ -1 & 0 & 0 \\ 0 & 1 & 0 \end{bmatrix}$$

Thus,
$$e \sin \theta = \frac{1}{2} [1, -1, -1]^{T}, \cos \theta = -\frac{1}{2}$$

i.e.
$$e = \frac{\sqrt{3}}{3} [-1, 1, 1]^{T}, \theta = 4\pi/3$$

...Eq. (17) is now written for $i = 1$, which yields eq (18a) with
$$A = \begin{bmatrix} 2 & 1 & 1 \\ 1 & 2 & -1 \\ 1 & -1 & 2 \end{bmatrix}, b = \begin{bmatrix} 3 \\ 2 \\ 1 \\ 1 \end{bmatrix}$$

Bouncholder reflections applied to A and b yield
$$A^{T} = \begin{bmatrix} -\sqrt{6} & -\sqrt{6}/2 & -\sqrt{6}/2 \\ 0 & -3\sqrt{2}/2 & 3\sqrt{2}/2 \end{bmatrix}, b^{T} = \begin{bmatrix} -3\sqrt{5}/2 \\ -\sqrt{2}/2 \\ -\sqrt{2}/2 \\ \end{bmatrix}$$

Hence, the minimum-norm solution of eq (24a) is
$$y_{0} = (-\frac{\sqrt{6}}{2}, -\frac{\sqrt{2}}{6}, 0)$$

trank on,
$$y_{0} = [1, \frac{2}{3}, \frac{1}{3}]^{T}$$

Finally,

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$$u = e^{T} (p_{1}^{1} - p_{1}) = \frac{\sqrt{3}}{3} [-1, 1, 1] \begin{pmatrix} 1 \\ 0 \\ 1 \\ \vdots \end{pmatrix} = -\frac{2\sqrt{3}}{3}$$

thereby completing the solution

Conclusions

An algorithm was presented, that allows the efficient computation of the screw parameters of the motion undergone by a rigid body between two finitely-separated positions. The efficiency of the algorithm refers to its low number of operations, as well as to its avoidance of large roundoff errors and of spurious singularities. Means of avoiding ill conditioning were discussed. Ways of filtering roundoff and/or meas-"Tement errors were outlined, but not treated in detail, since this fall outside the scope of the paper. The computation of the screw parameters for two infinitesimally-separated positions is discussed in an accompanying paper [18].

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AUTOMATIC COMPUTATION OF THE SCREW PARAMETER OF RIGID-BODY MOTIONS

PART. 11: INFINITESIMALLY - SEPARATED POSITIONS

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JUNIO, 1984.

AUTOMATIC COMPUTATION OF THE SCREW PARAMETERS OF RIGID-BODY MOTIONS. PART II: INFINITESIMALLY-SEPARATED POSITIONS. Jorge Angeles¹

Abstract

The approach introduced in an accompanying paper, simed at the computation of the screw parameters of a rigid-body motion defined by two finitelyseparated positions, is now applied to that defined by two infinitesimallyseparated positions. Given the economy of computation of this algorithm, it should allow the real-time computation of the screw parameters under study. The algorithm assumes perfect knowledge of the position and the velocity of three noncollinear points of the body.

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Introduction

A novel approach, simed at the computation of the screw parameters of a rigid-body motion defined by two finitely-separated positions, is introduced in an accompanying paper [1]. Many an application, either in robotics or in mechanism design, however, require the computation of the said parameters when the motion is defined by two not finitely-, but infinitesimallyseparated positions. Moreover, since infinitesimally-separated positions give rise to linear problems, whereas finitely-separated positions, to nonlinear ones, the latter are solved frequently by first solving the former, and then performing a time integration,

The algorithm introduced here is based upon that presented in [2], but modifies it in the sense of eliminating the spurious singularity contained therein. Moreover, the computations are simplified and the procedure extended to the computation of all the independent parameters of the screw motion under study.

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The motion defined by two infinitesimally-separated positions of a rigid body is fully described by the following [3, pp 119-148]: a) the axis of the instantaneous screw, given by the position vector of one of its points (preferably the one lying the closest to the origin) and three direction cosines, b) the sliding of the screw along its axis, and c) the rate of rotation about the axis, supplied with sign, given a positive direction defined on the axis. The set of scalar screw parameters of the rigid body motion is thus, the following: the three components of a vector r_0 , locating point R_0 of the screw axis 1, whose distance to the origin is a minimum: the three components of a vector e, parallel to 1 and defining the positive direction along L; two scalars, u and 0, representing the sliding along and the rate of rotation about L. This gives 8 scalar components, which are subject to:

 $e^{T}e = 1$ $r_{e}^{T}e = 0$

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To describe the present algorithm, let p_{i} (i = 1, 2, 3) be the position vectors of three noncollinear pints P_{i} (i = 1, 2, 3) of a rigid body, their velocities being denoted by p_{i} (i = 1, 2, 3). Now the angular velocity ω is computed as follows:

Let c be the centroid of the three points, v being its velocity. These two vectors are clearly given by

$$c = \frac{1}{3} \frac{3}{1} p_{i}, \quad v = \frac{1}{3} \frac{3}{1} \frac{p_{i}}{p_{i}}$$
 (1)

The velocity of any point of the body can be expressed in terms of that of one point A, a, whose position vector is represented by a, and the angular velocity ω of the body. If the angular velocity matrix Ω is used instead of vector ω , then the foregoing relation, written for each point P and C, is

 $\dot{p} = \dot{a} + \Omega(p_i - a), \ \dot{\iota} = 1, 2, 3$ (2a) $v = \ddot{a} + \Omega(c - a)$ (2b)

The relation between ω and Ω is

 $\omega = \text{vect}(\Omega)$

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 $\omega_{i} = \frac{1}{2} \varepsilon_{ijk} \Omega_{kj}$ (3b)

where the standard index notation is being used, ϵ_{ijk} being the alternating tensor. Subtraction of eq (2b) from eq (2a) yields

$$\dot{\mathbf{p}}_{i} = \mathbf{v} = \Omega(\mathbf{p}_{i} = \mathbf{c}), \ i = 1, 2, 3$$
 (4)

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(3a)

Following the approach introduced in [2], matrices P and P are next defined as

$$\mathbf{P} = \left[\mathbf{p}_{1} - \mathbf{v} \right] \quad \mathbf{p}_{2} - \mathbf{v} \left[\mathbf{p}_{3} - \mathbf{v} \right] \quad \mathbf{p}_{3} - \mathbf{c} \quad \mathbf{p}_{3} -$$

All three relations (4) can thus be expressed as

F = ΩP (6)

which is an equation not depending upon the location of the origin. Thus, the algorithm is not affected if two of the three given points turn to be collinear with the origin The algorithm, moreover, does not depend upon the regularity of matrix P; indeed, this matrix, as given by definitions (5), is identically singular, but this is no drawback, as shown next.

Taking the vector of both sides of eq (6) produces $\begin{bmatrix} 4 \\ \end{bmatrix}$:

$$\frac{1}{2} (1 \operatorname{tr} P - P) \omega = \operatorname{vect}(P)$$
(7)

which can be solved for w provided matrix M = 1 tr P - P is invertible. If this is the case, then

 $\omega = 2 (1 \text{ tr } P - P)^{-1} \text{ vect } (\dot{P})$ (8)

matrix in becoming singular only if P becomes a rank-one matrix which, in turn, implies that the three given points are collinear. This fact is supported by

Theorem 1. The trace of matrix P, as defined in eq (5), is identically differ-

Proof :

With no loss of generality, axes X-Y-Z are assumed to be orientated so that Z is perpendicular to the plane defined by the three given points. Furthermore,

5

.../...

the origin is placed at C, whereas X is orientated parallel to line P_3P_1 , so that $x_1 > 0$. Hence, $y_1 = y_3 < 0$, $y_2 = -2y_1 > 0$, the corresponding layout appearing in Fig 1





Matrix P is this given by

$$P = \begin{bmatrix} x_1 & x_2 & -(x_1 + x_2) \\ y_1 & -2y_1 & y_1 \\ 0 & 0 & 0 \end{bmatrix}$$

Henc:,

$$tr P = x_1 - 2y_1$$

which, under the foregoing assumptions, is the sum of two positive real numbers, and mence, never vanishes, q. e. d.

Theorem 2. Matrix 1 trP - P, with P defined as in eq(5), is singular if, and only if the three given points P_i (i = 1, 2, 3) are collinear. Proof:

Only necessity will be proved, for sufficiency can be proved by inverting the nicessity proof given. To this end, let π_i and μ_i be the proper values of P

and M = 1 tr P = P, respectively. Given the definition of M, these values \dots are related by

$$\mu_{i} = tr P - \pi_{i}, i = 1, 2, 3$$

However, the three columns of matrix P represent the components of three coplanar vectors, hence P is singular. Thus, at least one proper value of P vanishes. Let $\pi_3 = 0$. Hence,

$$tr P = \pi_1 + \pi_2$$

and

$$\mu_{j} = \operatorname{tr} P - \pi_{j}, \ \mu_{2} = \operatorname{tr} P - \pi_{2}, \ \mu_{3} = \operatorname{tr} P$$
(10b)

Therefore, the determinant of M, det(M), is given by

 $det(M) = \mu_1 \, \mu_2 \, \mu_3 = \pi_1 \, \pi_2 \, tr \, P \, . \tag{11}$

where relations (10a & b) have been taken into account. From eq(11) and Theorem 1, it is clear that det(M) vanishes only if at least one of \mathbf{v}_1 and \mathbf{v}_2 does. However, both of them cannot vanish simultaneously, for this would imply $P_1 = P_2 = P_3$. The vanishing of, say: \mathbf{v}_1 , implies that matrix P has two vanishing proper values, i. e. that it is a rank-one matrix. This implies, in turn, that the three column vectors of P are parallel to one single vector, i. e that the three points are collinear, thereby completing the proof.

Now, if the three points are given along the edges of a right-angled tribedron, so that their distances to the spex are identical, and vector c is defined now not as the centroid, but as the position vector of the spex, then a proper normalization renders matrix P orthogonal. In this case, matrix 1 tr P - P can be inverted explicitly as [4]:

(9)

(10a)

$$(1 \text{ tr } P - P)^{-1} = \frac{1}{\text{tr}^2} (P \text{ tr } P + P^T)$$
(12)

In this case, matrix 1 tr P - P turns to be singular if, and only if, P represents a rotation of the tribedron, with respect to the given reference frame, of an angle of rotation $\theta = \pi/2$, π or $3\pi/2$. This singularity can be readily removed by introducing a redefinition of the Cartesian axes attached to the reference frame, as discussed in [4].

In general, however, and particularly if c is defined as in eq (1), matrix P is not orthogonal. The computation of ω can be executed safely if matrix 1 tr P - P is not only nonsingular, but also well conditioned. Ill conditioning of this matrix arises only when the three given points, though not collinear, are close to it. A measure of the closeness to collinearity can be given, thus, by the inverse of the condition number [5] of this matrix. Most linear-equation solvers supply the user with a good estimate of the said number.

The remaining parameters are now computed. Given ω , it is a simple matter to construct Ω . Symbolically they are related by

 $\Omega = 1 \times \omega \tag{13}$

where x denotes the standard cross product, in dyadic notation.

Now, vector r_0 defining the position of point R_0 on the instantaneous screw axis, whose distance to the origin is a minimum, can be computed in two alternate ways, one resorting to an explicit formula, the second one, via a minimization problem. Both approaches are now discussed.

The formula giving r_o is [3, p 129]:

$$\mathbf{r}_{o} = \mathbf{p}_{i} + \frac{1}{\omega^{2}} \quad \omega \ge \mathbf{p}_{i} - (\omega + \mathbf{p}_{i})\omega \qquad (14)$$

which is valid for i = 1, 2, 3. If no roundoff nor measurement errors were present, formula (14), as applied to all three points, would yield one and the same value for r_0 . In practice, this is not the case; thus the said errors can be filtered if the formula is applied to the three given points, then defining r_0 as the mean value of the three distinct values thus obtained.

Alternatively, r_0 can be computed via the following optimization problem [3, pp 126-129]:

$$\operatorname{Min} \quad \frac{1}{2} r^{\mathrm{T}} r \qquad (15a)$$

subject to

А́Г = Ъ.

 $\Omega[p_{i} + \Omega(r - p_{j})] = 0$ (15b)

Matrix Ω being 3 x 3 and skew-symmetric, is singular; hence r cannot be solved for from eq (15b). In fact, if r could be solved for from that equation, then the said equation would produce one single value of r, not a set, that defining the screw axis. Eq (15b) contains exactly 2 linearly independent equations, for rank (Ω) = 2. These can be readily extracted from eq (15b) if Householder reflections are applied to both sides of eq (15b), as discussed in [1]. This would yield:

(16a)

.../...

with

$$A = \begin{bmatrix} a_{11} & a_{12} & a_{13} \\ 0 & a_{22} & a_{23} \\ 0 & 0 & 0 \end{bmatrix} \qquad b = \begin{bmatrix} b_1 \\ b_2 \\ 0 \end{bmatrix}$$

Eq (16a) represents then an underdetermined linear algebraic system possessing infinitely many solutions, its minimum- (Euclidean) norm solution being given by the generalized inverse of the upper $\ell \ge 3$ submatrix of A, referred to as A_u. Analogously, let b_u be the upper ℓ -dimensional subvector of b. Thus,

$$\mathbf{r}_{o} = \mathbf{A}_{u}^{+} \mathbf{b}_{u}$$
(17a)
with
$$\mathbf{A}_{u}^{-} \equiv \mathbf{A}_{u}^{+} (\mathbf{A}_{u}^{-} \mathbf{A}_{u}^{T})^{-1}$$
(17b)

Given the frequent ill conditioning of matrix $A_{\mathcal{U}}^{\mathbf{T}}$, it is not recommended to invert this matrix explicitly. In fact, r_0 can be more efficiently computed resorting again to Householder reflections, as discussed in [1]. This would produce r_0 as follows: let H be the product of Householder reflections rendering $A_{\mathcal{U}}^{\mathbf{T}}$ upper triangular. Then

$$H A_{\mu}^{T} = T = \begin{bmatrix} U \\ --- \\ 0 \end{bmatrix} \begin{bmatrix} 2 \\ 1 \end{bmatrix}$$
(18a)

U being upper triangular. Vector yo is computed from

$$\mathbf{y}_{0} = \begin{bmatrix} -\frac{\mathbf{y}_{\mu}}{2} \\ \mathbf{0} \\ \mathbf{1} \end{bmatrix}^{2} \mathbf{y}_{\mu}^{T} \mathbf{y}_{\mu} = \mathbf{b}_{\mu}$$
(18b)

from which the 2-dimensional vector y_u can be readily computed, since u^T is lower triangular. Then

$$\mathbf{r}_{o} = \mathbf{H}^{\mathrm{T}} \mathbf{y}_{o}$$
 (18c)

Clearly, each value of p_{i} and p_{i} (i = 1, 2, 3) would produce one pair (A, b), to form then three different minimization problems (15a & b). Again, if no roundoff nor measurement errors were present, then all three problems would produce one and the same value r_0 . Since such errors are always present, the likeliest value of r_0 can be chosen as the mean value of the three values thus obtained.

Having determined ω_{i} , e and θ can be readily computed from

 $\mu = e \theta$ (19) with e defined as a unit vector. Thus, only parameter u need be computed. This is done simply by projecting any of \dot{p}_{i} (i = 1, 2, 3) on e, thus obtaining

 $u = e^{T}p_{i}$ (20) which would again produce three slightly different values, out of which the likeliest would be their mean value.

Example

Given the three position vectors $\mathbf{p}_{\hat{\mathcal{L}}}$ of points $\mathbf{P}_{\hat{\mathcal{L}}}$, as well as their corresponding velocities, $\mathbf{p}_{\hat{\mathcal{L}}}$, for $\hat{\mathcal{L}} = 1, 2, 3$, all vectors referred to the same coordinate frame X-Y-Z, and shown pext, determine the parameters defining the instantaneous screw of the corresponding motion.

$$\mathbf{p}_{1} = \begin{bmatrix} 1 & 1 & 7 \end{bmatrix}^{T}, \quad \dot{\mathbf{p}}_{1} = \begin{bmatrix} 7 & -5 & 1 \end{bmatrix}^{T}$$
$$\mathbf{p}_{2} = \begin{bmatrix} 4 & 7 & 1 \end{bmatrix}^{T}, \quad \dot{\mathbf{p}}_{2} = \begin{bmatrix} -5 & 4 & 4 \end{bmatrix}^{T}$$
$$\mathbf{p}_{3} = \begin{bmatrix} 7 & 10 & 10 \end{bmatrix}^{T}, \quad \dot{\mathbf{p}}_{3} = \begin{bmatrix} 1 & -2 & 4 \end{bmatrix}^{T}$$

The vectors given above are taken from Example 1 of [2] for comparison purposes. They define in fact the positions and the velocities of three points of a rigid body. Indeed, the following compatibility condition holds for the said vectors:

$$(\dot{p}_{j} - \dot{p}_{j})^{T}(p_{j} - p_{j}) = 0, i = 1, 2, 3; j = 1, 2, 3; j \neq i$$

Now c and v are computed as given in eq (1):
 $c = [4, 6, 6]^{T}, v = [1, -1, 3]^{T}$

Matrices P and P, defined in eq (5), are given next:

•	- 3	0	3			6	- 6	0	Ì
P -	- 5	1	4	,	P -	- 4	5	- 1	
	1	- 5	4			- 2	: 1	1	

.../...

Substitution of the foregoing matrices and vectors in eq (8) produces $\omega = \begin{bmatrix} 1, & 1, & 1 \end{bmatrix}^{T}$

thus obtaining the same value of ω as that reported in [2]. The remaining parameters are now obtained, To this end, matrix Ω is now computed. From eq (13),

Now eq (15b) is expressed, for 4 = 1, as

	- 2	11	ī	[×]	$\begin{bmatrix} 0 \end{bmatrix}$
1	1	- 2	1	y -	0
	L 1	1	- 2	z .	[0]

x, y and z being the components of vector r of that equation, as yet to be determined. This is done as the solution to problem (15a & b), which is readily obtained by application of Householder reflections to both sides of the last equation, to render it into row echelon form. Letting H and A be the product of the two reflections required for the aforementioned transformation, and the matrix of the system involved, one has

HA =
$$\begin{bmatrix} 2 \alpha & -\alpha & -\alpha \\ 0 & \beta & -\beta \\ 0 & 0 & 0 \end{bmatrix}$$
, $\alpha \equiv \frac{\sqrt{6}}{2}$, $\beta = \frac{3\sqrt{2}}{2}$

Hence, eq (15b) represents the null space of matrix HA given above. Any

13.

vector r of that space has clearly the following components:

 $r = ||r|| [1, 1, 1]^T$

|| r || representing its magnitude. Hence the minimum-norm vector r_0 , solving problem (15a & b) is simply

 $r_{o} = [0, 0, 0]^{T}$

for the axis of the instantaneous screw, given by vector r, as found above, passes through the origin.

Parameters θ and e are readily obtained from vector ω , as $\dot{\theta} = || \omega || = \sqrt{3}$, $e = \frac{\omega}{\dot{\theta}} = \frac{\sqrt{3}}{3} [1, 1, 1]^{T}$

Finally, parameter u is found simply as the projection of any of p on e. For instance, for i = 1,

 $\mathbf{u} = \mathbf{e}^{\mathrm{T}}\mathbf{p}_{j} = 0$

and hence the motion is a pure rotation about the origin. The following remarks are now in order:

a) This problem turned to be very simple to solve, given that it reduced, in the last stage of computing r_0 , to finding the null space of a rank-2 2 x 3 matrix. In the general case, it would have given rise to an underdecermined linear system of 2 equations in 3 unknowns, whose minimum-norm solution would have been determined as outlined in [1].

b) The solution reported in [2] did not make evident that the motion of this example is a pure rotation about the origin.

A method was presented that allows the automatic computation of the screw parameters of a motion defined by two infinitesimally-separated positions. This method is simpler than previously reported ones, making use of a smaller number of operations. The latter feature can allow the real-time computation of the said parameters, which is essentially necessary in robotics applications. The method does not depend on the location of the origin of the coordinate reference frame; but its orientation can produce a "spurious singularity, that can be readily removed. Furthermore, it gives no preference to any of the three given points. It fails only if the three given points are collinear, but in this case no method can provide the parameters sought, for they are undetermined. A possible source of numerical instability is contained only in matrix 1 tr P - P, this matrix becoming ill conditioned as the three given points approach collinearity. Hence, this matrix is better conditioned, the more triangle $P_1P_2P_3$ approaches an equilateral triangle.

Acknowledgements

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DISENO CINEMATICO DE MAQUINARIA

SUBPROGRAMAS PARA EL ANALISIS DE CADENAS CINEMATICAS CON 7 ESLABONES ACOPLADOS MEDIANTE PARES DE ROTACION

DR. JORGE ANGELES ALVAREZ

JUNIO, 1984.

SUBPRICIANAS PARA PL THALIEIS DE CARENAS CURINATICAS CON 7 ECLAPONES ACOPLAS : ENDERRIE ESTES DE FORACION

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Abstract

Computer subgrograms for the analysis of powered ink revolute-coupled kinematic chains, both enough and open are presented. The subgrograms are based upon an algorithm that beings the displacement kinematic equalions that were obtained from invariance concepts. Newton-Supheon's method was used to solve the displacement constructs. These subgrograms can be used in real time, thus allowing the control of rebst manipulators.

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Lenneor

Se presentan subprogramas para el mélisis de cadenas cinculticas abiertas o corretas, compacentes par 7 estabates acoplados mediante pares de retación. Los unbragamas están bata dos en un algoritmo que resuelve ins eruaciones titemáticas que se obtavieros utilizados concep tos de invortencia. Se emples el actodo de deston-legimon para obtener na coloción de las eccaciones consultante enaminas. Estos nub programas porten foncionar un timpo real para el control de manipuladore.

introducción

El problema del calculo de los 6 ángeios dosconocídos en una cadena cinemático computata por 7 estabones acoplados mediante pure de ratación se conoce como problem cinemilico inver so y ha atroleko la atención de muchos inventiga dores, sobre todo en los últimos 10 años [1-4] pri su aplicabilidad en rebots manipuladores, En lus trabajos existentes se sigue Lúsicamente la misma tendencia, esto es', reducir las ecuaciones cinemánicos involucradas a un polinemio en una incógnita, cato ha dado lugar a que, pa ra obtener solucionan, se tengan que ballar raí ces de determinantes de matrices equ clementos polinómicos. Pur ejemplo, en 💈 se trabaja con una matriz de 12x12, cuyos elementos son de 4ª grado en la tangente de la mitad del ángelo de entrada, en Casto que en [4] se establace una matriz de 16x16 con elementos de 2ª grado en el miumo promento, Hosturiomento, Albua [5] obtiene una mutriz de 16x16 con elementos du 2º grado en la tangente de la mitae del ángu le de salida, lo cual consume demasiaria ticajo para obtiner una coloción aceptable, para esta fonaulación produce un excesivo ocror de redon dee. The Angulo's intermedica, 07, ..., 06, se colcalam ballando las raíces de polingaios de grado b* o memori. Sin embargo, na el proceso de eliminación de los ángulos internation de in troduten taleas aspurlas, por la qua debe veri ficarse cada solución ballada [6].

El algoritmo que agai de emplea se obticue aplicando concepcos de invariancia cost permitenun cálcolo auy oficiento de las occivadas norm lugradas, lo cual, hasta la fecto, no ba abús ajdizado jou ningún tuvest igador. - Oto da lu gar a und inghabita suferde jubu mit jerterfer fig portigada [7] r con esto ne puest le nor, con un número may limitado de operaciones (ny ite ración y con un error de reslondes terrelacite bajo, el análisis de cañenas electórico, alter Las a corrados, composition de contre existences articulado, medicate porte de returión, de ca dona abierto, los datos del prelivat son les Vo rishies can definen of sovimizato del fugado terminal (07) on table que, en codena corrada, son las que definer el servizionte del calabér wotor o "do entrada".

pescripción del almonitas

Las couteignes part el málicis de desplaça miento de una cadenta citerática ao resto grado de Histori pictor da morte de las condiciones de cerradula de dasplavanicato y de estación que se dissection de sensible. De laberda con el mátedo y la notación da benavit y malechery. [5, 8], los n' estalones de vici calita chi di Lica se mejeran ordenno-bette, de la n y al i* estaván so fija el sistema coordenado Si, Yi, Zi. Asī, [Qi, 143] i tertemeta uta matriz ortogonal, refering a low even Aj, Y., 71 que gita éstos a una posición concidente. con los correspondientes $X_{i+1}, Y_{i+1}, B_{i+1}$ por su parte, $\begin{bmatrix} a_i & i+1 \\ a_i & i+1 \end{bmatrix}_1$ es el vertes que une los origenes $\begin{bmatrix} a_i & i+1 \\ a_i & i+1 \end{bmatrix}_1$ do los entre unte riores, dirigido del granero al regenció y rete rido a ojes fijos al ut erlatšu. Joi, Lou condiciones de corradura son (ver Fig.1):

$$\begin{bmatrix} \mathbf{g}_{1,2} \end{bmatrix}_{1} \begin{bmatrix} \mathbf{g}_{2,3} \end{bmatrix}_{2} \cdots \begin{bmatrix} \mathbf{g}_{n,2} \end{bmatrix}_{0} = \begin{bmatrix} \mathbf{j} \\ \mathbf{j} \end{bmatrix}_{1}$$
(1)

para rotación y

$$\begin{bmatrix} a_{1,2} \\ a_{2,3} \end{bmatrix}_{1} + \begin{bmatrix} a_{2,3} \\ a_{2,3} \end{bmatrix}_{1} + \cdots + \begin{bmatrix} a_{2,7} \\ a_{3,7} \end{bmatrix}_{1} + \begin{bmatrix} a_{1,1} \\ a_{2,3} \end{bmatrix}_{1} + \cdots + \begin{bmatrix} a_{2,7} \\ a_{3,7} \end{bmatrix}_{1} + \begin{bmatrix} a_{1,1} \\ a_{2,7} \end{bmatrix}_{1} + \begin{bmatrix} a_{2,3} \\ a_{3,7} \end{bmatrix}_{1} + \cdots + \begin{bmatrix} a_{2,7} \\ a_{3,7} \end{bmatrix}_{1} + \begin{bmatrix} a_{2,7} \\ a_{3$$

Las consciones (1) y (2) constituyen on sintens de 12 constitues electrons, de las curles pôlo 6 seu independentes. So superes recursej des $[\widetilde{O}]_1$, que represente la retación (ella gi tar los ejes X₁, Y₁, Z₁ a cua avienterión idéntica a la de los ejes X₂, Y₂, X₃, así co mo $[\mathbf{z}]_1$, que es el vector de proteción del parto p del GC (origen del 7ª sustera condena del. Aplicando in invariancia del vector aci. [9] as obtienes vect (Q) = 0 set ϕ

- (H

Signdo c of vector característico real asocia de con el valor característico real +1 y ϕ , el ámplio de retación, "Combinado (1) y (3) su tiene:

vect
$$([g_1, 2], [g_2, 3]_2, \dots, [g_6, 7]_3)^{m} \in \text{nen } (4)$$

Las consciones (2) y (4) constituyes on sigtoma algebraico no lineal de sexto orden de la format

$$f(\overline{0}) = \begin{bmatrix} \overline{1}_1(\overline{0}) \\ f_2(\overline{0}) \end{bmatrix} = \overline{0}$$
 (5a)

donates

 $\underbrace{f_{2}(0)}_{1} = \begin{bmatrix} a_{1,2} \\ a_{1,2} \end{bmatrix}_{1}^{+} \begin{bmatrix} a_{2,3} \\ a_{2,3} \end{bmatrix}_{1}^{+} \cdots + \begin{bmatrix} a_{6,7} \\ a_{6,7} \end{bmatrix}_{1}^{-} \begin{bmatrix} x \\ x \end{bmatrix}_{1}^{+0} \qquad (5c)$

La ecuación (5b) aparece multiplicada par 2 con el fin de rvitar divisiones posteriores en tre 2 que pueden incrementar el número de obra ciones necesarias:

Su puede obtener la solución do (Sa) aplican do el mítodo de Newton-Japhion que, con un valor inicial próximo a una solución, converge cuadrá ticamento, requiricado así pero ticapo para obte ner una solución [10].

El esquena iterativo de Newton-Raphana es el signicitar

$$\underbrace{\underline{\theta}^{k+1}}_{i=k} \mapsto \underbrace{\underline{\theta}^{k}}_{i=k} \wedge \underbrace{\underline{\theta}^{k}}_{i=k} \qquad (6n)$$

donde &0 . us la selución de

 $\tilde{\mathfrak{I}}(\tilde{\mathfrak{h}}_{\mathbf{k}}) ? \tilde{\mathfrak{h}}_{\mathbf{k}} = - \tilde{\mathfrak{I}}(\tilde{\mathfrak{h}}_{\mathbf{k}})$

y $J(\theta^{2})$ on la matriz Jacobiana, evaluada en $\theta = f^{4}$, del sistema (5a), que se calcula combi

$$\tilde{\Sigma}(\tilde{\theta}) = \begin{bmatrix} \tilde{\sigma} \tilde{\Gamma}^{2}(\tilde{\theta}) \sqrt{2\tilde{\theta}} \\ \tilde{\sigma} \tilde{\Gamma}^{1}(\tilde{\theta}) \sqrt{2\tilde{\theta}} \end{bmatrix}$$
(3)

donde

$$\frac{\partial f_1(\underline{\theta})}{\partial \theta_1} = 2 \text{ vect } \left(\begin{bmatrix} \underline{\theta}_1, \underline{2} \end{bmatrix}_1 \begin{bmatrix} \underline{\theta}_2, \underline{3} \end{bmatrix}_2 \cdots \begin{bmatrix} \frac{\partial \underline{\theta}_{1, i+1}}{\partial \theta_1} \end{bmatrix}_i \\ \begin{bmatrix} \underline{\theta}_{i+1, i+2} \end{bmatrix}_{i+1} \cdots \begin{bmatrix} \underline{\theta}_{6, 7} \end{bmatrix}_6 \right)$$
(8a)

Pura ol cilculo de $\partial f_2 / \partial 0$, definase $x_1(0) = \begin{bmatrix} a_1, 2 \\ b_1, 2 \end{bmatrix} ; \begin{bmatrix} a_2, 1 \\ b_2, 1 \end{bmatrix} ; 2^{a_2} ... ; \begin{bmatrix} a_2, 1 \\ c_1, 2 \end{bmatrix} ; \begin{bmatrix} a_2, 1 \\ c_2, 1 \end{bmatrix} ; 2^{a_2} ... ; \begin{bmatrix} a_2, 1 \\ c_2, 1 \end{bmatrix} ; 2^{a_2} ... ; \begin{bmatrix} a_2, 1 \\ c_3, 2 \end{bmatrix} ; 2^{a_3} ... ; \begin{bmatrix} a_3, 2 \\ c_3, 2 \end{bmatrix} ; 2^{a_3} ... ; \begin{bmatrix} a_3, 2 \\ c_3, 2 \end{bmatrix} ; 2^{a_3} ... ; 2^{a_3} .$

que puede calcularse mediante el algoritmo de Normer para evaluación da polimentos [10] come:

$$\underline{\mathbf{x}}_{\mathbf{k}} = \begin{bmatrix} \mathbf{e}_{\mathbf{k}_{1},\mathbf{k}+1} \end{bmatrix}_{k}^{k} \begin{bmatrix} \mathbf{e}_{\mathbf{k}_{1},\mathbf{k}+1} \end{bmatrix}_{k}^{k} \begin{bmatrix} \mathbf{e}_{\mathbf{k}_{2},\mathbf{k}+1} \end{bmatrix}_{k}^{k} \end{bmatrix}_{k}^{k} \begin{bmatrix} \mathbf{e}_{\mathbf{k}_{2},\mathbf{k}+1} \end{bmatrix}_{k}^{k} \begin{bmatrix} \mathbf{e}_{\mathbf{k}_{2},\mathbf{k}+$$

45.1

$$\frac{2n^2}{2n^2} = \frac{2n^2}{2n^2}$$
(9.1)

La equación (8.) puede sumplificarse 🕖 en la forma

$$\frac{\partial v_1}{\partial t_1} = (1 \text{tr} \ \mathbf{p} - \mathbf{p}) \left[\mathbf{p}_1, \mathbf{z} \right]_1 \cdots \left[\mathbf{p}_{n-1}, \mathbf{1} \right]_{n-1} = (2)$$

siendo e, el vector anitario paralelo al eje de recaelon del 1º par einentícico de rocación, y

$$= \left[\underline{\rho}_{1,2} \right]_1 \left[\underline{\rho}_{2,3} \right]_2 \cdots \left[\underline{\rho}_{2,3} \right]_0 \tag{10}$$

La velocidad se calcula a partir de la suguien te relación:

$$\operatorname{Voct}\left(\underline{p}, \underline{p}^{T}\right) = \omega \tag{11}$$

dondo y os la velocidad anjular especticula del OT, referida al pistona l, en tamio que

$$\frac{\partial z_1(\underline{\theta})}{\partial \theta} \stackrel{\text{d}}{=} \underbrace{\nabla} \tag{12}$$

siendo y la velocidad del punto p del Ofreferida al sistema L. Las acuationes (LL) y (L2), después de elganas súmplitionetores [i], quedan como:

$$J(\theta) = d \tag{13.4}$$

donde
$$d = \begin{bmatrix} \psi \\ y \\ y \end{bmatrix} y \psi = (1 \text{ tr } P - P) \psi$$
 (13b)

La acolezación se calcula a partir de las ecua ciones (11) y (12) como $[\underline{U}]$:

3(9) 0 - 9

9=

$$\begin{bmatrix} (1 \text{ tr } p-p) (\dot{\omega} - \frac{1}{2} \dot{\omega}^{T} - \frac{\eta^{2} \dot{\omega}}{-\eta} \dot{\psi}) \\ \underline{a} - \frac{1}{2} \dot{\psi}^{T} - \frac{\eta^{2} \dot{\omega}}{\eta} \dot{\psi} \end{bmatrix}$$
(14b)

en las cuales \hat{w} es la aceleración angular del OF y a es la aceleración del ponto p del OT, ambas prescritas y referido, al sistema considenado 1.

De acuerdo con el algoritmo exposito, los sub lorgianais realizan la solución numérica de las ceuaciones (5A) (13a) y (14a), y producen los valores de 0, 6 y 0 correspondientes a los pres critos de 0 lo equivalentemente, de c y 0, r, y, ω , a y ω . A continuación se describen las variables que intervienen en el cólculo, así como las sobratibas y sus argumentos. Según la Fly 1, y la notación de Lemavit y Hartenberg, Se tiene, con c () \equiv cos () y s () \equiv sen ():

$$\begin{bmatrix} \mathbf{a}_{1,i+1} \end{bmatrix}_{i} = \begin{bmatrix} \mathbf{a}_{i} \mathbf{c} \mathbf{0}_{i}, & \mathbf{a}_{i} \mathbf{s} \mathbf{0}_{i}, & \mathbf{b}_{i} \end{bmatrix}^{T}$$
(15)

donde \mathbf{x}_1 of la distancia entre los ejes \mathbf{z}_1 y \mathbf{z}_{1+1} , luego siempre positivo, y b₁ es la coordenada de la intersección de \mathbf{x}_{1+1} con \mathbf{z}_1 en el sistema \mathbf{x}_1 , Y₁, Z₁. Así, la matriz $\begin{bmatrix} \mathbf{Q}_{1,1+1} \end{bmatrix}_1$ se puede expresar como

$$\begin{bmatrix} \alpha_{i,i+1} \end{bmatrix}_{i} = \begin{bmatrix} c\alpha_{i} & -a\theta_{i} & c\alpha_{i} & a\theta_{i} & a\alpha_{i} \\ c\theta_{i} & c\theta_{i} & c\alpha_{i} & -c\theta_{i} & a\alpha_{i} \\ 0 & a\alpha_{i} & c\alpha_{i} \end{bmatrix},$$
(16)

donde α_1 es el Jugalo entre π_1 y π_{1+1} , medi éx en la dirección positiva de π_{1+1} .

Con lo anterior se pueden definir los arreglos THETA y P. Los primeros 32 elementos de P son dutos y deben almacenarse como se indica a continuación:

```
1+1,...,6
```

THETA(1) :	contiene los ángulos asocia
	dos a cada par de la cadena
	cinemática, oca (15) y (16)
P(J):	almacena a, de la ec (15)
P(1+6);	almacena b_1^2 de la ec (15)
· P(2+12);	almacena Gi y posturiormente,
	$\cos \omega_1$ de la uc (15)
P(1,13):	almacena sea u _i do la oc (15)
P(25);	almacena 🖗 de la uc (3)

```
3=1,2,3
```

P(J+25): almacona g de la ec (3) P(J+20): almacena las coordenadas del OT reforidas al sistema l (lado derocho de la ec (2)) P(J2): almacena 27

D: P(33) hasta P(32) so definen en la t<u>a</u> bla 1.

En la Pig 2 se muestra el diagrama jerárquí co del programa.

Las subrutinas que se emplean son;

NRDAMP(THETA, P, UF, 1P, P, TOLX, TOLF, DAMP, N, ITER, MAX, MAX)

Obtiene las raíces de un sistema algebraico no linuil da orden N, aplicando el método de Newton-Raphson con amortiguamiento (04DADD*1) 8. El amortiguamiento tiene por objeto acele rar la convergencia. En esta subrutina, P re presenta el miembro izquierdo de la se (5a), que es de dimensión 6 y se calcula en FRN: DF es la matriz Jacobiana de la ec (7) y de dimen sión 6x6 y se calcula en DSPX. La solución de (6b) se obtiene por medio de las subrutinas DECAMP y SOLVE que se describen posteriormente. TOLX en la tolerancia impuesta en la aproxima ción a la solución, en tanto que TOLF es la to lorancia que se acepta en la función l'EMAX es el número múximo de iteraciones permitidos, ITER en el número de la itoración ojocutada y KMAX, ul número máximo de amortiguamientos que

se permiten per iteración: la es un arregto en tero de dimensión 6, se obtiene en DEXONE y es el indico pivotal en la descumposición LO.

FURITHERA, F. P.N)

Forma las equaciones de desplazamiento y de rotación que, al anularse, proporcionan el v_d lor de los ángolos 6, para una posición dada, ec (5a). Para este tran se requieren las subr<u>u</u> timas PROD y VECX. Estas requieren a sa vez el arregio TeS, definido como:

I#1,...,6

$$TCS(I) = THETA(I)$$
$$TCS(I+G) = COS(THETA(I))$$
$$TCS(I+I2) = SIS(THETA(I))$$

con lo cual se conservan los valores de Thura, un cada (teración.

PROD(TCS, P)
Efectual los productos :

$$\begin{bmatrix} 0 \\ 1, 2 \end{bmatrix} 1$$

 $\begin{bmatrix} 0 \\ 1, 2 \end{bmatrix} 1$
 $\begin{bmatrix} 0 \\ 2, 3 \end{bmatrix} 2$
 $\begin{bmatrix} 0 \\ 1, 2 \end{bmatrix} 1$
 $\begin{bmatrix} 0 \\ 2, 3 \end{bmatrix} 2$
 $\begin{bmatrix} 0 \\$

y los almacena de acuerdo al arresto costinita en la Tabla 1. Para lograrlo evalúa (10) con los valores de $0_i \neq u_1$ para i(1=2,...,6) y ejecuta el producto matricial con el último arregio almacenado. El receltado es quardado en memoria con la ubicación mostrada en la tabla 1.

VECX(TCS, P)

Calculu el vector x_1 , ecs (6b y 5c), así como su derivada con respecto a 0, ec (8d) y los alunc<u>e</u> na en el arreglo P, como se nuestra en la Tubla 1. Esto se debe a que existen conjuntos de el<u>e</u> racionos comunes para el câlculo de x_1 y su d<u>u</u> rivadu.

DEDX(THETA, DF, P, N)

Calcula la matriz Jacobiana definida en la ec. (7) y la almacena en (7)(6,6).

DECOMP (N, N, DE, COND, IP, WORK)

Ecompone una matriz real (DF) en el producto DJ y estima su condición (CXD) [11], con lo que se conoce la amplificación del error de rodendeo. Para efectuar esta estimación se emplea el arregio WORK. IP en el vector pivote, que almacena información molere los intercambios de rempiones necesarios para evitar un exercivo error de redondeo.

SOLVE (R, N, DF, P, IP)

Resulve el sistema <u>d</u>(p+f) expleando la matriz DF factorizada en DFCOMP por sustitución $t \leq \frac{1}{2}$ greniva, considerando los intercampins de re<u>n</u> glones efectuados en DECAMP.

VISLAC(N, DP, 1P, P, V, A)

CAlcula $D \neq Q$ cuando se tiena convergencia en NRDAP, intorna al programa principal y llega a Ceto cua la matriz Jacobiana doscompleta en al producto LD. En caso de que casualmente, o a propósito, se de a NRDAMP valoren solución a las oca (Sa), esta subrutina no utiliza DECOMP, por lo que no se tiene la matrix Jacobiana dog compuesta: en esta caso, llama a DFDX y a LEOSU para entrar a VELAC. A osta subrutina se le su ministran los valores de May y en el arreglo V, \tilde{w} y n en el arreglo A y regrona los resultados de \tilde{y} y \tilde{y} en V y A respectivamente, o sea, obtie no la colución de las ecuaciones (13a) y (14a).

Ejemplo

Los subprogramas se aplicaron para obtener el análisis de una malla cerrada de 7 estabones que se muestra en la Pig 3, cuyos parámetros Gon:

 $a_1 = 1$ (u de long)
 i = 1, ..., 6

 $a_7 = 3$ (u de long)
 i = 1, ..., 7

 $b_1 = 0,$ i = 1, ..., 7

 $a_1 = a_3 = a_4 = a_5 = a_7 = 0^*$ $a_2 = a_5 = a_7 = 0^*$

Los resultados se muestran en la fig 4 en la que se graficó sólo θ_2 ve θ_1 , ya que por las simultías de este mocanismo, $\theta_2 \cdot \theta_5 \cdot \theta_3 \cdot \theta_6$ y $\theta_4 \cdot 2\theta_1$ [7]. El error que só obtuvo al comparar θ_2 con la solución analítica obtenida en [7] se muestra en la Fig 5: on las Figs 6 y 7 se moestra los errores de θ_2 y θ_2 , cospectivamente, compara dos con la solución analítica.

Conclusiones y extensiones

ios resultados hasta abora obtenidos demues tran la potencialidad del algoritmo. En la so lución del ejemplo se requirieron aproximadamen te 17 segundos, para la rama superior o sea, pa ra $120^{-4} \partial_1^{-5} 240^{\circ}$ y $0^{-5} \partial_2^{-5} 120^{\circ}$. Este timpo pue de disminuirse al empluar un lenguaje de máquina.

La subrutina VELAC no calcula <u>0</u> ni <u>0</u> cuando la matriz Jucobiana es singular, lo cual sucedo cuando su (ienen configuraciones or remas de las cadenas cinemíticas, por lo que se continúa la investigación en este sentido.

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Fig 1 Parámetros que dufinen la arquitectura de una cadena cuemática



Fig 2 Diagrama jerárquico del programa



Plg 3 Cadena cinemática corrada 7R

- Eje que apunta hacia afuera e i plano de la fogura.
- X Eje que apunta bacia adentro del plano de la figura



Fig 5 Error on el cálculo de $\theta_2 \, \cos \theta_2 \ge 0$







Pig 4 Respecto cinenálica del necalismo 74 de la Pig J.

Tabla 1

P	P_10_	¹ '39	¥.,2	P.45	² 48	P_51	Г ₅₄	T 57
34	⁹ 37	10	ذه ''	P46	P 44	۴ ₅₇	F	Г 18
P15	°38	P41	۲ 44	1° 47	1° ₅₀	1 ¹⁷ 51	P-6	1
<u> </u>			£122			P1P2P3		

^р 60	р 63	P 66	Р ₆ э	P72	175	Т.	² е1	i'H4
P ₆₁	F 64	67	۶ 70	P 73	1'76	774	P ₈₇	P _{HS}
P 62	P 65	ά ^τ .	۲,	1. 74	۳ ₇₇	1' 80	103	1.
<u> </u>	^Q ₂ Q ₃ Q ₄		$\overline{v_1}\overline{v_2}$	1. P. 185	·	$\left[\frac{c_1c_2c_3c_4c_5c_5}{c_1c_5c_5}\right]$		

P ₈₇	P.90	P. 13	F 90	¹ '93	10
P UR	⁷ 91	[£] 94	97	100 B	
Р. Н2	1 [.] 92	P.95	ETEI	101	10-1
×1	×2	X.I	3.	X	x

P 105	P 108	1 111	114	P1.7	1:0
P106	"109	P 112	115	P ₁₁₀	1.1
P 107	¹ 110	² 11)	110	°119	122
³⁶ 2	1 ¹ 2	$\frac{1}{2}$	' [〔] 2	f	P' 2
90 1	32.	-10 	СТ.		ມ_

 $(Q_i, Q_{i+1})_i$ so indications Q_i



DISENO CINEMATICO DE MAQUINARIA

MODELIZACION DE CADENAS CINEMATICAS CON PARES DE ROTACION Y PRISHATICOS

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MODELIZACION DE CADENAS CINEMATICAS CON PARES DE ROTACION Y PRISMATICOS

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Abst<u>ract</u>

A model of kinematic chains with prismatic and revolute pairs is presented. The model is based upon an algorithm using invariance concepts and solves the arising kinematic equations via the method of Newton-Raphson. The aplication is illustrated with a fully solved example, namely the analysis of a spatial RSRC linkage

Resumen

Se presenta un modelo de cadenas cinemáticas con pares de rotación y prismáticos. El modelo está basado en un algoritmo que emplea conceptos de invariancia y resuelve las ecuaciones cinemáticas resultantes mediante el método de Newton-Raphson. La aplicación se ilustra con un ejemplo que consiste en el análisis de un mecanismo espacia) RSRC.

Introducción

El problema cinemático inverso consiste en determinar los ángulos o los desplazamientos en una cadena cinemática compuesta por n eslabones acoplados

mediante pares inferiores, suponiendo conocidas las historias de posición. velocidad y aceleración de uno de los eslabones. Esta cadena puede ser abier ta (manipuladores) o cerrada (mecanismos). Dada la importancia que tienen actualmente los manipuladores industriales y la complejidad del problema inverso, numerosos investigadores han propuesto diferentes soluciones, principalmente en forma cerrada, para casos particulares simples, empleando matrices de 4x4 que contienen rotación y desplazamiento simultáneamente [1,2,3], o bien reduciendo las ecuaciones cinemáticas involucradas a un polinomio en una incógnita. Este aparece expresado como un determinante de una matriz de 12x12 con elementos de 4º grado en la tangente de la mitad del ángulo de entrada [4], o bien de una de 16x16 con elementos de 2º grado en el mismo argumento [5,6]. Alizade y Duffy [7] proponen la relación entre los datos y las incógnitas a través de un conjunto de 30 ecuaciones, algunas de ellas dependientes de la topología, las cuales no son de fácil aplicación para el cálculo de velocidad y aceleración. Whitney [8] orienta su trabajo hacia la solución del problema inverso aplicando cambios diferenciales en la posición, con lo que se obtiene el jacobiano que relaciona los cambios en posición y orientación dados, como variables dependientes, con las variables de los pares considerados independientes, calculándose fácilmente sólo la velocidad. Un mét<u>o</u> do para calcular la aceleración en cada par.conocidas la posición y la velocidad en cada uno de éstos, así como el momento o la fuerza de entrada, ha sido pro puesto por Walker y Orin [9]. En este trabajo se presenta un modelo general aplicable en tiempo real a manipuladores con 7 eslabones con arquitectura arbitraria, conociendo las historias de posición, velocidad y aceleración del órgano terminal (OT) del manipulador. El modelo es una generalización del expuesto con anterioridad [10,11], el cual no tiene la flexibilidad de sustituir un par prismático por uno de rotación. El presente modelo es aplicable

a cadenas con seis pares de rotación o bien a una combinación de cinco pares de rotación y un prismático, el cual puede estar ubicado en una posición arbitraria con respecto a los de rotación. Esta última topología se encue<u>n</u> tra en una gran cantidad de manipuladores, tales como el brazo stanford (USA) [1], el VWR30 (RFA), el ACMA Cribier H-80 (Francia) y el Komatsu RCA70 (Japón) [12].

Descripción del algoritmo generalizado

De acuerdo con el método y la notación de Denavit y Hartenberg [13], que se resume en la Fig l, los n eslabones de una cadena cinemática se numeran ordenadamente de l a n y al i² eslabón se fija el sistema coordenado X_i , Y_j , Z_i como sigue:

Z se dirige a lo largo del eje del par i , que une los eslabones i e i+1, si es de rotación, o alineado con la traslación del par i cuando éste es prismático.

X_i se define sobre la perpendicular común a Z_{i-1} y Z_i, dirigida de Z_{i-1} a Z_i. Y_i completa el sistema coordenado dextrógiro del iº eslabón. La posición relativa de dos eslabones adyacentes está completamente definida por los siguientes parámetros:

a_i, distancia entre los ejes Z_i y Z_{i+1} α_i , ángulo entre Z_i y Z_{i+1}, medido en la dirección positiva de X_{i+1}.

 b_i , coordenada entre los ejes $X_i \neq X_{i+1}$, constante cuando el par es de rotación. θ_i , ángulo entre los ejes $X_i \in X_{i+1}$ medido en la dirección positiva de Z_i , constante cuando el par es prismático.

Sean

$$\begin{bmatrix} 0_{i,i+1} \end{bmatrix}_{i}^{i} = \begin{bmatrix} c\theta_{i} - s\theta_{i}ca_{i} & s\theta_{i}sa_{i} \\ s\theta_{i} & c\theta_{i}ca_{i} & c\theta_{i}sa_{i} \\ 0 & sa_{i} & ca_{i} \end{bmatrix}$$
(1)

$$\begin{bmatrix}a_{i,i+1}\end{bmatrix}_{i} = \begin{bmatrix}a_{i}c\theta_{i}, a_{i}s\theta_{i}, b_{i}\end{bmatrix}^{T}$$
(2)

donda c() $\equiv \cos()$ y s() $\equiv sen()$ $[Q_{i,i+1}]_i$ y $[a_{i,i+1}]_i$, abreviadas como Q_i y a_i , respectivamente, representan una matriz ortogonal referida a los ejes X_i , Y_i , Z_i , que gira éstos a una posición coincidente con los correspondientes X_{i+1} , Y_{i+1} , Z_{i+1} , y un vector que une los orígenes O_i y O_{i+1} , dirigido del primero al segundo y referido a los ejes del iº eslabón, respectivamente.

Las condiciones de cerradura en orientación y desplazamiento son:

$$Q_1 Q_2 \cdots Q_5 Q_6 = Q \tag{3}$$

$$[a_1]_1 + [a_2]_1 + \dots + [a_6]_1 = r$$
(4)

que constituyen un sistema de 12 ecuaciones escalares en 6 incógnitas, siendo independientes sólo 6 de ellas. Las cantidades conocidas son Q y r, que representan el giro necesario para sobreponer los ejes X_1 , Y_1 , Z_1 con X_7 , Y_7 , Z_7 y el vector que une el origen del sistema inercial con un punto R del OT. Al obtener el vector axial [14] de Q se tiene:

vect
$$(Q) = e \operatorname{sen} \phi$$
 (5)

donde e es el vector característico real de Q asociado con el valor característico +1 y ϕ , el ángulo de rotación. Al combinar (3) y (5) se tiene:

$$\operatorname{vect} (Q_1 \ Q_2 \dots \ Q_6) = e \, \operatorname{sen} \phi \tag{6}$$

Las ecuaciones (4) y (6) constituyen un sistema algebraico no lineal de sexto orden de la forma:

$$f(\theta) = \begin{bmatrix} f_r(\theta) \\ f_t(\theta) \end{bmatrix} = 0$$
(7)

donde $\theta = [\theta_1; \theta_2, \dots, \theta_6]^T$ representa cinco valores diferentes asociados con pares de rotación y uno que puede estar asociado con un par prismático, o bien, con otro par de rotación; además:

$$f_{r}(\theta) = 2 \operatorname{vect}([Q_{1,2}]_{1} [Q_{2,3}]_{2} \cdots [Q_{6,7}]_{6}) - 2 \operatorname{sen} \phi = 0$$
(8a)

$$f_{t}(\theta) = [a_{1,2}]_{1} + [a_{2,3}]_{1} + \dots + [a_{6,7}]_{1} - [r]_{1} = 0$$
(8b)

La solución de (7) se obtiene aplicando el método de Newton-Raphson, que converge cuadráticamente con valores próximos a la solución (15]. El esquema iterativo de Newton-Raphson es el siguiente:

$$\theta^{k+1} = \theta^k + \Delta \theta^k \tag{9a}$$

donde ∆0^k es la solución del sistema

t

$$J(e^{k}) \Delta e^{\tilde{k}} = -f(e^{k})$$
(9b)

y $J(\theta^k)$ es la matriz jacobiana evaluada en $\theta = \theta^k$ del sistema (9b), que se

calcula a partir de:

$$J(\theta) = \begin{bmatrix} \partial f_r(\theta) / \partial \theta \\ \partial f_t(\theta) / \partial \theta \end{bmatrix}$$
(10)

Cuando θ_1 está asociado a un par de rotación se tiene:

$$\frac{\partial f_r(\theta)}{\partial \theta_i} = (I \text{ tr } P-P) Q_1 Q_2 \dots Q_{i-1} e_i$$
(11a)

siendo

. •

$$P = [Q_{1,2}]_1 [Q_{2,3}]_2 \dots [Q_{6,7}]_6$$

$$\vdots$$

$$e_i: 'vector unitario paralelo al eje z_i$$
(11b)

I : matriz identidad 🕚

$$\frac{\partial f_{\pm}(\theta)}{\partial \theta_{i}} = \frac{\partial x_{1}}{\partial \theta_{i}} = \frac{\partial x_{1}}{\partial \theta_{i}} = Q_{1}Q_{2} \dots Q_{i-1}(e_{i} \times x_{i})$$
(11c)

 x_i se calcula aplicando el algoritmo de Horner [15] para evaluación de polinomios, esto es, como:

$$x_{6} = [a_{6,7}]_{6}$$
(11d)
$$x_{k} = [a_{k,k+1}]_{k} + [Q_{k,k+1}]_{k} x_{k+1} = 5,4,3,2,1$$

Cuando el par es prismático, b_j es una traslación (ver Fig 1) y se emplea:

$$\frac{\partial f_r(\theta)}{\partial \theta_i} = 0 \tag{12a}$$

δ.

$$\frac{\partial f_t(\theta)}{\partial \theta_i} = \theta_1 \theta_2 \dots \theta_{i-1} e_i$$
 (12b)

7.

La velocidad se calcula a partir a las siguientes expresiones:

$$\frac{\partial x_{1}}{\partial \theta} \dot{\theta} = v \tag{13a}$$

$$vect(PP_{i}^{T}) = \omega$$
 (13b)

siendo v la velocidad lineal del punto P del OT en el sistema 1 y ω , la velocidad angular vectorial del OT. Después de algunas manipulaciones, las ecuaciones (13) quedan como [11]:

$$J(\theta) \dot{\theta} = \dot{r}^{\prime} \tag{14}$$

donde

$$\mathbf{r} = \begin{bmatrix} \boldsymbol{\omega} \\ \mathbf{v} \end{bmatrix} \mathbf{y} \quad \boldsymbol{\omega}' = (\mathbf{I} \ \mathbf{tr} \mathbf{P} - \mathbf{P}) \boldsymbol{\omega}$$
(15)

válida sin modificaciones para pares de rotación o prismáticos.

La aceleración se obtiene a partir de:

ŧ

$$\dot{\omega} = \frac{d}{dt} \operatorname{vect} (\dot{P}P^{T}) = \frac{\partial \omega}{\partial \dot{\theta}} \ddot{\theta} + \frac{1}{2} \dot{\theta}^{T} \frac{\partial^{2} \dot{\omega}}{\partial \dot{\theta}} \dot{\theta}$$
 (16a)

$$\ddot{v} = \frac{\partial x}{\partial \theta} \ddot{\theta} + \dot{\theta}^{T} \frac{\partial^{2} x_{1}}{\partial \theta^{2}} \dot{\theta}$$
(16b)

Se requieren ahora las siguientes expresiones cuando se tiene un par de rotación

$$\frac{\partial^2 \dot{\omega}}{\partial \theta_i \partial \theta_j} = [e_i]_1 \times [e_j]_1 \quad i = 1, \dots, 6 \quad (17a)$$

$$\frac{\partial^2 x_1}{\partial \theta_j \partial \theta_j} = \left[e_{,i} \right]_1 \times \frac{\partial x_1}{\partial \theta_j}$$
(17b)

Cuando se tiene un par prismático en la iª posición se tiene

$$\frac{\partial^2 \hat{\omega}}{\partial \theta_j \partial \theta_j} = 0$$
(18a)

$$\frac{\partial^{2} \times_{1}}{\partial \theta_{i} \partial \theta_{j}} \neq [e_{i}]_{1} \times [e_{j}]_{1} \quad i < j$$
(18b)

Despues de algunas simplificaciones [11], se tiene:

2.

$$J(\theta) \ddot{\theta} = \ddot{r}$$
 (19)

donde

$$\overset{"}{r} = \left[\begin{pmatrix} I \ tr \ P-P \end{pmatrix} \begin{bmatrix} \dot{\omega} - \frac{1}{2} \ \dot{\theta}^{T} \ \frac{\partial^{2} \dot{\omega}}{\partial \dot{\theta}^{2}} \ \dot{\theta} \end{bmatrix} \right] (20)$$

$$\overset{"}{r} = \left[\dot{v} - \dot{\theta}^{T} \ \frac{\partial^{2} x_{1}}{\partial \theta^{2}} \ \dot{\theta} \end{bmatrix}$$

El modelo cinemático aquí propuesto está formado por las ecs (8a y b), (14) y (19), que proporcionan los valores de las variables asociadas con los pares cinemáticos de la cadena, así como sus dos primeras derivadas. Las ecuaciones dichas se realizaron en un programa de computadora que es aplicable indistintamente a cadenas cinemáticas provistas de seis pares de rotación o bien de 5 pares de rotación y uno prismático.

Ejemplo

Se desea simular el mecanismo mostrado en la Fig 2, que es del tipo RSRC. El punto O_7 permanece fijo, cambiando de orientación a razón constante de 1 rad/s. En la Fig 3 se muestran los sistemas coordenados involucrados, así como la sustitución del par esférico por tres pares de rotación con los ejes concurrentes en el centro del esférico, en tanto que el par cilindrico.

por un par de rotación y otro prismático. Para este mecanismo particular se obtiene fácilmente una expresión que relaciona el giro 4 con el despl<u>a</u> zamiento \$ [16] mediante:

$$s(t) = a \, sen \phi + \sqrt{b^2 - c^2 - a^2 \cos^2 \phi}$$
 (21)

no siendo así para las demás variables

El error obtenido al comparar el valor calculado de s con el dado por la ec (21), se muestra en la Fig 4. Se emplearon los siguientes valores que aparecen en la Fig 2: a = 1.0, b = 3.0 y c = 2.0 en unidades de longitud.

Conclusiones

Los resultados se obtuvieron en una microcomputadora Apple IIe, siendo 580 el número total de operaciones requeridas para el cálculo de θ , $\dot{\theta}$ y $\ddot{\theta}$ por posición, en base a una iteración en el método de Newton-Raphson, cuando todos los pares son de rotación, reduciéndose este número a 556 el introducir un par prismático. El algoritmo presentado se puede aplicar por lo tanto al control de manipuladores como el brazo stanford [1] o al análisis de mecani<u>s</u> mos de las más variadas topologías, como se mostró con el ejemplo.

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DISENO CINEMATICO DE MAQUINARIA

ANALISIS CINEMATICO DE MANIPULADORES CON ARTICULACIONES REDUNDANTES HACIENDO USO DE MINIMIZACION CUADRATICA

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JUNIO, 1984.

" AMALISIS CINERATICO DE HANIPULADORES CON ARTICULACIONES REDULEANTES HACIENDE USO DE MINIMIZACION CÚADRATICA "

Resumen

Se desarrolla un programa de computadora para la solución de sistemus de ecuaciones no lineales subdeterminados, haciendo uso de minimización cuadrítica. Como aplicaciones particulares se presentá el problema del análisis cinemático de manipuladores con articulaciones redundantes.

El programa que se presenta puede resolver, esimismo, un problema donde la función objetivo, a minimizar, no sea de norma cuadrática.

Introducción

Debido a que en vi estudio de manipuladores se puede presentar el caso de articultaciones redundantes, es decir, que se tengan grados de libertad extra, obteniendose un sistema de ecuaciones subdeterminado es necesario recolver el problema tratando de extremizar alguna función de los parame- v tras del mecanismo, en este cuso, se toma la norma cuadrática de sus ángulos. El programa desarrollado puede aplicarse a cualquier otro tipo de problemas de sistemas de ecuaciones no lineales subdeterminados, extremizando funciones cuadráticas o no cuadráticas.

Desarrollo

El problema consiste en hallar una solución particular de un sistema algobraico no líneal, de la forma:

. (1)

$$f(\underline{x}) = 0$$

donde f, x y 0 son vectores de dimensión m, n y m, respectivamente, con m < n. Dado que el problema es subdeterminado, se dispone, en general, de multiples soluciones, por lo que es necesario proponer una función objetivo z = z(x) a extremizar, siendo z un número real. Por lo tanto, podemos plantear el problema como:

 $\min z = z(x)^{2}$

f(x) = Q

sujeta a

Y

Para establecer un método que resuelva este problema, se recurren a argumentos geométricos. Cada función $f_i(x)$ representa una superficie en un especio Rⁿ, al que pertenece x. El conjunto de ecuaciones representado por (1), es decir:

 $f_{1}(x_{1}, x_{2}, \dots, x_{n}) = 0$ $f_{2}(x_{1}, x_{2}, \dots, x_{n}) = 0$ \dots $f_{m}(x_{1}, x_{2}, \dots, x_{n}) = 0$

define un subconjunto de puntos de Rⁿ que satisfacen simultaneamente l**as** m ecuaciones no lineales. Así, pues, esas ecuaciones definen una curva Cen Rⁿ, de grado de libertad n-m, que representa la intersección de esas m superficies.

Supóngase que se conoce un punto P_0 de Γ en R^n , con vector de posición \underline{x}^0 , que satisface el sistema (1). Se desea ahora determinar un nuevo punto P^1 de Γ para el cual la función objetivo tenga un valor menor que para \underline{x}^0 , esto es, se desea determinar un vector \underline{x}^1 tal que:

 $f(\underline{x}^{1}) = 0$ $z(\underline{x}^{1}) < z(\underline{x}^{0})$

si tel punto no existe, x^o es un mínimo local del problema.

Debido a que Γ es una curva en Rⁿ, no es posible, en general, definir un punto P¹ de Γ que esté muy lejado de P⁰. Se buscará, entonces, un punt: cercano e P⁰, linealizando el problema y buscando un punto u¹, de vector de posición x¹ que este contenido en la tangente de Γ en P⁰, orien-

tado en l. dirección opuesta a $\sqrt{2}$. Como G¹ puede estar tan alejado de P¹ como sea, es necesario imponer una condición adicional, es decir, que en G¹, z alcance su valor mínimo a lo lorgo de la tangente. Para esto definimos:

$$\overline{x}_1 = x_1 + \Delta \overline{x}_1$$
 (2)
$$\Delta \overline{x}_1 = -\Delta \underline{g}$$
 (3)

siendo ξ un vector (no necesariamente unitario) contenido en la tangente de P, siendo perpendicular a'los m vectores $\{\varphi f_i\}$, i = 1, 2, ..., m. Tal vector satisface, entonces, el siguiente sistema de ecuaciones:

$$\begin{bmatrix} (\nabla f_1)^T \\ (\nabla f_2)^T \\ \cdots \\ (\nabla f_m)^T \end{bmatrix} = \begin{bmatrix} 0 & J \\ \\ \end{bmatrix} = 0$$
 (4)

La ecuación (4) establece que ξ se encuentra en el espacio de J y define un plano perpendicular a la tangente a Γ , como se muestra en la fig. 1. Si $J^T \lambda$ es el vector del plano más próximo a $\forall z$, entonces

$$\underline{\mathbf{\xi}} = \nabla \mathbf{z} - \mathbf{J}^{\mathrm{T}} \underline{\mathbf{\lambda}}$$
 (5)

así, λ se obtiene como la solución de mínimos cuadrados de

$$\mathbf{J}^{\mathrm{T}} \mathbf{\lambda} = \mathbf{\nabla} \mathbf{Z}$$
 (6)

que es 🕤

donde

 $\lambda = (JJ^{T})^{-1}J \forall z$

por lo que, sustituyendo (7) en (5)

$$= [1-J^{\mathrm{T}}(JJ^{\mathrm{T}})^{-1}J] \nabla z \qquad (8)$$

(7)

y de (3)

$$\Delta \bar{x}_{1}^{\mu} \propto \left[J (J J^{T})^{-1} J = 1 \right] \forall z$$
(9)

donde \prec se escoge como el valor que minimiza $z(x^{\circ} + b\bar{x}_{1})$, entonces, \prec se determina como la solución al problema:

$$\min_{\mathbf{x}} \mathbf{z}(\mathbf{x}^{O} - \mathbf{A}\mathbf{S}) \tag{10}$$

donde tanto xº como § se conocen.

Э

Para el caso mas general, en el que z sea una función arbitraria de x, « se puede determinar resolviendo un problema no lineal de minimización en una dirección. Una forma eficaz de resolver tal.problema es mediante la subrutina FETN [1].

Si en perticular, la función objetivo es cuadrática, es decir, para algún \overline{x} fijo se tiene:

 $z = \frac{1}{2} \left(x = \overline{x} \right)^{T} W(x - \overline{x})$ (11)

siendo W positiva definida, se tiene que z tendra un mínimo a lo largo de la tangente en un punto donde se anule dz/d4 , es decir, el problema se reduce a la búcqueda del mínimo en una dirección; pero

	$\frac{dz}{dx} = \left(\frac{\partial z}{\partial x^{1}}\right)^{T} \frac{\partial x^{2}}{\partial x} $	12)	
siendo	<u>x¹ - x - x - x⁰ - x<u>5</u> - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x - x </u>	(13)	
por lo que	$\frac{\partial z}{\partial x} = W(x - \bar{x}) = W(x^{\circ} - \bar{x} - 4\xi)$	(14)	
y .	$\frac{\partial x^1}{\partial x} = \frac{\partial}{\partial x} (x^0 - x\xi) = -\xi$	(15)	
por lo tanto	$\frac{dz}{dA} = -\xi^T w(x^0 - \bar{x} - A\xi) = 0$	(16)	
÷	$\mathbf{x} \mathbf{\underline{S}}^{\mathrm{T}} \mathbf{W} \mathbf{\underline{S}} = \mathbf{\underline{S}}^{\mathrm{T}} \mathbf{W} (\mathbf{\underline{X}}^{\mathrm{O}} - \mathbf{\underline{X}})^{-1}$	(17)	٠
pero	$W(\underline{x}^{\circ} = \overline{\underline{x}}) = \overline{\nabla} z$	(18)	
por lo tanto	$\mathbf{K} = \frac{\mathbf{\xi}^{T_{W}} \mathbf{v} \mathbf{z}}{\mathbf{\xi}^{T_{W}} \mathbf{\xi}}$	(19)	
	$\begin{bmatrix} \mathbf{T} \mathbf{W} \mathbf{q} \mathbf{z} = (\mathbf{Y} \mathbf{z})^{\mathbf{T}} \begin{bmatrix} 1 = \mathbf{J}^{\mathbf{T}} (\mathbf{J} \mathbf{J}^{\mathbf{T}})^{-1} \mathbf{J} \end{bmatrix} \mathbf{W} \mathbf{\nabla} \mathbf{z}$	(20)	
	$\boldsymbol{\zeta}^{\mathrm{T}} \mathbb{W} \boldsymbol{\underline{\zeta}} = (\boldsymbol{\nabla} \boldsymbol{z})^{\mathrm{T}} [\boldsymbol{1} - \boldsymbol{J}^{\mathrm{T}} (\boldsymbol{J} \boldsymbol{J}^{\mathrm{T}})^{-1} \boldsymbol{J}] \mathbb{W} (\boldsymbol{\nabla} \boldsymbol{z})^{\mathrm{T}} [\boldsymbol{1}$	$- J^{T}(JJ^{T})^{-1}J$ (21)	، د
Si además, V = 1			
· · · ·	$ \sum_{k=1}^{T} \left[\sum_{n=1}^{T} \left(\nabla_{z} \right)^{T} \left[1 - J^{T} \left(JJ^{T} \right)^{-1} J \right]^{2} \nabla_{z} = $		
	$= (\nabla z)^{T} [1 - J^{T} (JJ^{T})^{-1} J] \nabla z$	(22)	

por lo tanto de 1

Una vez determinada κ , se tiene $\Delta \chi^1$ calculando <u>y</u> de la ecuación (6)

por medio de las subrutinas Hecomp y Holve [2]. Sin embargo se tendrá la configuración de la fig. 2, de donde se observa que, aunque $z(\bar{x}_1) < z(\bar{x}^\circ)$, \bar{P}_1 , cuyo vector de posición es \bar{x}_1 , esta fuera de Γ , esto es,

o sea, \bar{x}_1 no satisface el sistema (1). Es necesario determinar ahora una corrección $\delta \bar{x}_1 = \bar{x}_1$ tal que

$$\underline{x}_1 = \overline{x}_1 + \underline{A} \overline{x}_1$$
 (24)

(23)

si satisfaga a (1). Para resolver este problema existen varias alternativas que son:

Caso 1.- Como $\Delta \bar{x}_1$ está en el complemento del plano normal a Γ , si \bar{P}_1 no esta muy alejado de P_o, se puede pensar que, si se pasa un plano por \bar{P}_1 . paralelo a la normal a Γ , ese plano cortará a Γ , como se muestra en la fig. 3. La simplificación estriba en que $\Delta \bar{x}_1$ se habrá obtenido mediante la solución del sistema

> $f(\vec{x}_{1} + \delta \vec{x}_{1}) = 0$ (25) $\vec{x}_{1} = J^{T} \mu$ (26)

donde

Como se conocen \tilde{x}_1 y J, la incógnita es A, de dimensión m. Así la ecuación (25) se reduce a

토(신) = 일 · (27)

que es un sistema algebraico no lineal de m ecuaciones en m incógnitas, , que se puede resolver por el método de Newton - Raphson para sistemas determinados, esto se hace mediante la subrutina NRDAPP [3] que a su vez requiere de 165 subrutines DECOMP y SOLVE para la solución de sistemas de ecuaciones determinados en base a la descomposicion LU de matrices[4]. Caso 2.- Otra alternativa es seleccionar P₁ como el punto de Γ más próximo a \overline{P}_1^c , esto es, como el punto de tangencia de una esfera centrada en \overline{P}_1 con $\overline{\Gamma}$. En este caso, $\Delta \overline{x}_1$ se puede obtener como la solución al problema (fig. 2):

$$\frac{1}{2} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^{T} \sum_{i=1}^$$

sujeta a $f(\bar{x}_1 + \Delta \bar{x}_1) = 0$ (29)

Incorporando la ecuación (29) a la función objetivo, mediante multiplicadores de Lagrange . Así, se define:

$$\Psi(\Delta \bar{x}_{1}) = \frac{1}{2} \Delta \bar{x}_{1}^{T} \Delta \bar{x}_{1} + \sqrt{f} f(\bar{x}_{1} + \Delta \bar{x}_{1})$$
 (30)

que tiene un extremo cuando

$$\frac{d\Psi}{d\lambda\bar{x}_{1}} = \lambda\bar{x}_{1} + \left(\frac{\partial f}{\partial\bar{x}_{1}}\right)^{T} \dot{Y} = 0$$
(31)

donde

As1

por

$$\mathbf{x} = \bar{\mathbf{x}}_1 + \mathbf{b} \bar{\mathbf{x}}_1 \tag{32}$$

$$\frac{\partial f}{\partial k_{1}^{2}} = \frac{\partial f}{\partial x} = \frac{\partial x}{\partial k_{1}^{2}} = J \cdot 1 = J$$
(33)

Sustituyendo (33) en (31) y despejando, se tiene

$$\Delta \bar{\bar{x}}_1 = J^T \bar{Y}$$
 (34)

lo tanto:
$$f(\ddot{x}_1 + J^T \chi) = f(\chi) = 0$$
 (35)

que es un sistema no lineal de ecuaciones en m incógnitas como en el caso enterior, por lo que también puede resolverse mediante la subrutina NRDANP. Caso 3.- Por último, se puede usar

$$J \Delta x = -f$$
(35)

o bien
$$\Delta x = -J^{T}(JJ^{T})^{-1}f.$$
 (37)

actualizando J. En cualquier caso, se alcanzará un mínimo cuando se satisfagan los criterios de primerio y de segundo orden, obteniendose a continuación. Adjuntando la restricción (1) a la función objetivo z = z(x), obteniéndose

$$\frac{\partial}{\partial t} \left(\frac{x}{2} \right) = 2(\frac{x}{2}) + \frac{1}{2} \frac{T}{f}(\frac{x}{2})$$
(38)

se tiene un punto estacionario de Ψ cuando se anul<u>ă</u> su gradiente, escogiendo <u>X</u> de manera que produzca una <u>X</u> que satisfaga (1). As**i**,

$$\frac{\partial \Psi}{\partial x} = \Psi x + J^{T} \chi = 0$$
(39)

o bien, $J^{T} \lambda = - \forall z$ (40) que son las ecuaciones de normalidad, o sea la condición necesaria de primer orden para la existencia de un mínimo. La interpretación geométricade la ecuación (40) es que, en un punto estacionario, $\forall z$ se encuentra en el plano normal a \Box . La interpretación algebraica de esta ecuación es que, en un punto estacionario, $\forall z$ está en el codominio de J^{T} , un subespacio de R^{n} , de dimensión m < n.

La colución de la ecuación (36) puede obtenerse mediante la subrutina SUBDEL que resuelve un sistema algebraico dáneal subdeterminado, auxiliado por la subrutina HECOMP.

Descripción del programa

El programa desarrollado, SANLSU, puede resolver el problema según los casos 2 y 3 descritos anteriormente. En la fig. 4 se muestra el diagrama de flujo del programa. Los listados del programa y las subrutinas para cada ejemplo resuelto se muestran en el apéndice A.

Una vez obtenidos los resultados del problema, se almacenan en un archivo de datos que es leído por el programa GRAMAN que dibuja el mecanismo en pantalla con opción de graficar en papel empleando la subrutina SCRPRT. El listado del programa GRAMAN también se encuentra en el apéndice A.

La computadora empleada en la solución de estos problemas ès una APPLE IIe de 64K de memoria, con lo que se demuestra que no es necesaria gran cantidad de memoria para resolver este tipo de problemas.

Ejemplos

Ejemplo 1: Considere el siguiente problema en el que m=1, n=2 y f y zestan dadas, respectivamente por:

$$f(x_1, x_2) = x_1^2 + 4x_2^2 - 1 = 0$$

$$z(x_1, x_2) = 1/2(x_1^2 + x_2^2)$$

Determinar los puntos de coordenadas (x_1, x_2) que, satisfaciendo $f(x_1, x_2) = 0$, minimicen z.

Solucion:

La interpretación geométrica del problema se ilustra en la fig. 5, donde le elipse esta representada por la función P y Z representa la mitad del cuadrado de la distancia de un punto al origen. Como se puede observar de la figuro, en los puntos A, B, C y D z alcanza valores estacionarios que representan la solución del problema.

· La matriz jacobiana de f y el gradiente de z son, respectivamente:

$$\mathbf{J} = \begin{bmatrix} 2x_1 & 8x_2 \end{bmatrix} \qquad \forall \mathbf{z} = \begin{bmatrix} x_1 \\ x_2 \end{bmatrix}$$

En los puntos A y C, J y - z toman los siguientes valores:

$$\mathbf{z} = \begin{bmatrix} 2 & 0 \end{bmatrix} \qquad \mathbf{z} = \begin{bmatrix} 1 \\ 0 \end{bmatrix}$$

y en B y D:

 $J = \begin{bmatrix} 0 \\ 4 \end{bmatrix} \qquad \qquad z = \begin{bmatrix} 0 \\ 0.5 \end{bmatrix}$

De estos resultados se verifica la ecuación (6), es decir, la estacionaridad de los puntos A, B, C, y D. Los resultados del problema se muestran : en el apóndice B¹⁷ 57 55

Ejemplo 2: Se requiere posicionar un punto en un plono mediante una cadena cinemática abierta (menipulador articulado) de triple grado de libertod, con cuatro eslabones articulados, mostrada en la figura 6. La longitud de

sus eslabones noviles a_2 , a_3 , a_4 son unitarias. Determinar los ángulos θ_1 , θ_2 y θ_3 entre sus eslabones para que el punto P tenĝa las coordenadas del punto Q(0,1). Más aún, se requiere que el punto P genere una trayectoria prescrita.

Solución:

Debido a que en el problema existen infinidad de soluciones que posicionan el punto P en G, se debe seleccionar la mejor de estas según algún criterio que permita escoger una función z a extremizar.

De la fig. 6 se puede ver que la ecuacion f(x) = 0 es :

$$f_{1}^{(\theta_{1},\theta_{2},\theta_{3})} = a_{1}c(\theta_{1}) + a_{2}c(\theta_{1}+\theta_{2}) + a_{3}c(\theta_{1}+\theta_{2}+\theta_{3}) - x = 0$$

$$f_{2}^{(\theta_{1},\theta_{2},\theta_{3})} = a_{1}s(\theta_{1}) + a_{2}s(\theta_{1}+\theta_{2}) + a_{3}s(\theta_{1}+\theta_{2}+\theta_{3}) - y = 0$$

se puede elegir 2 como :

 $z = 1/2(\theta_1^2 + \theta_2^2 + \theta_3^2)$ donde: c() = cos() y s() = sen()

que púede representar una función de costo, donde se desea minimizar las rotaciones de los eslabones. En base a lo anterior, podemos formar el jacobiano de f y el gradiente de z que son, entonces:

 $J = \begin{bmatrix} -5(0_1) - 5(0_1 + 0_1) - 5(0_1 + 0_1) + 5(0_1 + 0_2 + 0_3) - 5(0_1 + 0_1) - 5(0_1 + 0_2) + 5(0_1 + 0_2) - 5(0_1 + 0_1) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2) - 5(0_1 + 0_2)$

	-	[ө,]
5V	3	0,
. '		[0,]
	-	

Los resultados numéricos, así como las gráficas de los manipuladores se mustran en el apéndice B. En primer lugar se resolvió el problema según el cado 2 antes descrito (fig. 7), la s figuras 8 y 9 representan al manipulador siguiendo una trayectoria recta, habiéndose resuelto el problema, según el caso 3, en la figura 9 sólo se muestran algunas de las configuraciónes que se siguieron para llegar al punto máximo.

Otro problema que se resolvió fue el de generar una trayectoria alea-

toria. En la figura 10 se méloran los resultados según el coso 3 y en la figura 11 se mélora la configuración original y el primer punto de la trayectoria, ya que el caso 2 no conværgió debido a la gran separación entre los puntos de la trayectoria.

Ejemplo 3: Obtener la posición óptima de un manipulador de 6 eslabones articulados, a partir de una posición dada. Solución:

La configuración general del mecanismo se muestra en la fig. 12. Las ecuaciones para el anilisis de desplazamiento de una cadena cinemática pueden obtenerse mediante las condiciones de cerradura de desplazamiento y de rotación. De acuerdo con el método y la notación de Denavit y Hartenberg [5] los n eslabones de una cadena cinemática se numeran ordenadamente, 1 a n y a i⁰ eslabón se fija el sistema coordenado X_1 , Y_1 , Z_1 . Así $[Q_{1,1+1}]_i$ representa una matriz ortogonal, referida a los ejes X_i , Y_i , Z_i que gira éstos a una posición coincidente con los correspondientes X_{i+1} , Y_{i+1} , Z_{i+1} ; por su parte $[a_{1,i+1}]_i$ es el vector que une los origenes O_i y O_{i+1} de los ejes anteriores dirigidos del primero al segundo y referido a ejes fijos al i⁰ eslabón. Así, las condiciones de cerradura son:

 $[a_{1,2}]_1 [a_{2,3}]_2 \dots [a_{5,6}]_1 = [a]_1$

para rotación y

 $[a_{1,2}]_1 + [a_{2,3}]_1 + \dots + [a_{5,6}]_1 = [\underline{r}]_1$

para desplazamiento, donde $[r]_1$ es el vector de posición del punto P del organo terminal. '

Debido a que en nuestro caso no importa la orientación del organo terminal, podemos considerar unicamente las condiciones de cerradura para desplazamiento y, haciendo coincidir el punto B con el A, tendremos un sistema de 3 ecuaciones con cinco incognitas, es decir:

 $f(\theta, \theta, \theta, \theta, \theta) = [a_{1,2}]_1 + \dots + [a_{5,6}]_1 - [r]_1 = 0$

Para el cálculo de 81/86 definamos:

$$\mathfrak{L}_1(\Theta) = \mathfrak{L}_1 + \mathfrak{L}_1\mathfrak{L}_2 + \mathfrak{L}_1\mathfrak{L}_2\mathfrak{L}_3 + \mathfrak{L}_1\mathfrak{L}_2\mathfrak{L}_3\mathfrak{L}_4 + \mathfrak{L}_1\mathfrak{L}_2\mathfrak{L}_3\mathfrak{L}_4\mathfrak{L}_5$$

donde : $a_i = \begin{bmatrix} a_{i,i+1} \end{bmatrix}_i$

 $\underline{\mathbf{q}}_{\mathbf{i}} \cdot [\underline{\mathbf{q}}_{\mathbf{i},\mathbf{i}+1}]_{\mathbf{i}}$

que puede calcularse mediante el algoritmo de Horner para evaluación de polinomios [6] como:

$$\begin{array}{rcl} x_5 &= \frac{a_5}{2} & & \\ & \frac{x_5}{k} &= \frac{a_k}{k} + \frac{Q_k}{k} x_{k+1} & & k = 4, 3, \dots, 1 \\ \end{array}$$
As1: $\begin{array}{rcl} \partial f &= \frac{\partial x_1}{\partial \varphi_1} & & \\ & \partial \varphi_1 & & \partial \varphi_1 \end{array}$

según la fig,13 y la notación de Denavit y Martenberg se tiene

$$\begin{split} \mathbf{\hat{n}}_{i} &= \begin{bmatrix} \mathbf{\hat{a}}_{i} \mathbf{c}(\mathbf{\hat{\theta}}_{i}) \\ \mathbf{a}_{i} \mathbf{s}(\mathbf{\hat{\theta}}_{i}) \\ \mathbf{\hat{b}}_{i} \end{bmatrix} \\ \\ \mathbf{\hat{Q}}_{i} &= \begin{bmatrix} \mathbf{c}(\mathbf{\hat{\theta}}_{i}) & -\mathbf{s}(\mathbf{\hat{\theta}}_{i}) \mathbf{c}(\mathbf{x}_{i}) & \mathbf{s}(\mathbf{\hat{\theta}}_{i}) \mathbf{s}(\mathbf{x}_{i}) \\ \mathbf{s}(\mathbf{\hat{\theta}}_{i}) & \mathbf{c}(\mathbf{\hat{0}}_{i}) \mathbf{c}(\mathbf{x}_{i}) & -\mathbf{c}(\mathbf{\hat{\theta}}_{i}) \mathbf{s}(\mathbf{x}_{i}) \\ \mathbf{0} & \mathbf{s}(\mathbf{x}_{i}) & \mathbf{c}(\mathbf{x}_{i}) \end{bmatrix} \end{split}$$

donde a_i es la distancia entre los ejes $Z_i \neq Z_{i+1}$, b_i es la coordenada de la intersección de X_{i+1} con Z_i en el sistema X_i , Y_i , Z_i , α_i es el ángulo entre $Z_i \neq Z_{i+1}$, medido en la dirección positiva de X_{i+1} .

Para la solución de este problema se utilizaron, además, las subrutinos VECX y PRODQ [7].

Los resultados se muestran en el apéndice B, así como las configuraciones del manipulador en las figuras 14 y 15.

Conclusiones

Como se puede observar, el método más eficiente fué el implementado para el cauo 3, yu que no es necesario que los puntos de la trayectoria esten muy cercanos, como se requiere para el caso 2, además, el tiempo de proceso para el caso 2 fue mayor que para el tercer caso.

Hay que hacer notar que este programa puede ser implementado en cualquier computador ya que no requiere de gran capacidad de memoria.

El programa puede modificarse ficilmente para resolver cualquier problema do este tipo, es decir, solución de sitemas no lineales de ecuaciones subdoterminados.







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FIGURA 3

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FIGURA 6



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PEDGRAM SAMISU

ESTE FROGRAMA CALCULA LA SOLUCION DE UN SISTEMA ALGEBRAICO NO LINEAL SUDDETERMINACO DU LA FORMA ;

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MIN Z=7(X) X

SUJETA AL F (X) #0 DIMENSION 07(2), DELTX(2), X1(2), DELTX1(2), XFIN(2), U(2) REAL MO(1), JAC (1,2), JACT (2,1), J1L (2) COMMON AU (D) J1, N CONHON CSI (2) WRITE(*, *(At) *) * SI DESEA IMPRESORA, TWOLEE : READ(*, (11) *) * 1F (1.EP.1) THEN OPEN(6,FILE# 'FPINTER: ') FUSE DECNIG.FILE= (CONSOLE: 1) **ETIDIE** 18111576,50 WRITE(+,'(AD)')' 1 SI EL PROBLEMA ES DE NORMA QUADEATICA ' READ(*,'(II)') II IF (11.EC.1) ALFA-1.0 WRITE(6.100) FE4D (*.200) H.H WRITE(5,200) M.N. W3ITE(6,300) READ(*.400) (XO(I),1-1,1) WRITE(6,400) (XO(J), In1, N WRITE (6,350) TEL.XH0.0000001 TOLF-0.000001 DOMPHO. 3 NAX+200 FTIAX-109 WRITE(C.COO) TOLY, TOLF, DAME, MAX, KMAX MU(1)=0.2 IONT=FORT+1 CALL JACOBI (JAC) CALL GRAD(02) PO 5 141,N -4.3 JTL (1) =67 (1) CO 5 J⇒1,N JACT (I, J) = JAC (J, I) CALL HECOMP (N, H, M, JACT, U) CALL HOLVE (N.N.K.JACT, U. JTL) DG 3 (#1.4 DO 3 J#1.M JACT((,J)=JAC(J,() CALL HULVED (N.H. JACT, JTL, JTL) DO 4 141.N CSI(I)+67(I)+670(I)

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LELTX (Diserved FABCET (F) X1(1)=X0(1)(D41 T7(1) 6 CALL NEDALF (MILLA), JOCT, TOUX, TOUF, DAKE, ITER, MAX, KMAX) CALL MULVEC (N. P. JAST, MU. DELTXI) 20 7 1~1.8 XFIN(I)=))(J)+DELTX1(E) 9 DO 14 141.00 (F(ABS(XFIN(1)-20(1)).6T.(.0E-6) GD TD 12 CONTINUE 14 60 TO 10 12 DO 13 I-1,N . 13 X0(1)=7F10((1) 60 TO IL WRITE (6, 1/A#) 1) I LA SOLUCIÓN FINAL DEL PROPLEMA ES: 1 10 HRITE 16, 500% ALEA WRITE(6,001) (DEI(1),T=1,N) WRITE(6,505) (MU(T), T+1,M) WRITE(6,T07) OFDUCE),101,4D \$T0P FORNAT(3X, ' PROOFAMA PARA) A SOLUCION DE BISTEMAS DE ECUACIÓ 'NEB'7157.' NO LINEALES SHODETERMINATÚS'7) FORMAT(3Y, ' OAME LAS DIMENSICHES DE F(X) Y X (M Y N)') 50 t. 100 150 FORMAT(11) 200 FORMAT (212) FORMAT(3), NO BE LEEGA A LA BOLUCION DERFUES DE SO VECEET). FORMAT(37, 1448E X01) 200 200 FORMATIC: TOLK, TELF, DAMP, MAX Y KMAX () 350 400 FERMAT (SPICES) 300 EDEMOTIVEX, TALEATVEX, E10.4) FORMAR(37, 101737, 1910, 4) FORMAR(3), 100737, 910, 4) 501 505 507 FOPHAT(32, 1XF(1)/32, 2F10.4) FERMAT (VC/5.5,F5.3,2150) 600 700 FORMAT (2F10.5, f 5.0, 213) END C C SUGROUTINE JACOBI (JAC) C C SUBRUTING FOR CALCULA EL JACOBIANO DE LA FUNCION F(X) ENFLEADA Ç EN EL PROGRAMA BACLOU. C C ENTRADOP: лÇ, C H - DIMENSION LE X Ċ хò VALOR INICIAL Ç SALIDA9: С - JACOBIANO DE F(X) - JVC 1 . . . C REAL JAC(1.2) COMMON XO(2), N.N. PO 1 1*1,M DO 1 3=1,N JAC([,J)≠0.0 1 JAC(1.1)-7.0+YO(1) JAC(1,2/#3.0#X0(2) RETURN END

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Ċ С SUSFOUT ME CRAD (57) c ç - EUBRUTINA DUE CALCULA EL OSADIENTE DE LA FUNCION $\mathcal{I}(\hat{\mathbf{X}})$ EMPLEADA EN EL PROGRAMA SAPLEU. ENTRADAS: N DIMENSION DE Z(X) XO FUNDO INICIAL BALIDASI c GZ - GRADIENTE DE ZIXX c DIMENSION 07(2) COMMON XO(2) .M.N DO 1 (#1.0) 6Z(()=0.0 1 00 2 1+1,14 2 GZ (7) +X0(1) ACTORN -£240 SUBROUTING FUNCE, JACT (P) COMMON X0(2), N, O DIMENSION Y(0),P(2),F(1),X(1) REAL JACT (2,1) DO I INI,H 60 1 U.M. N Y(J)=P(J)+JACT(J,I)=X(I) L F(1) + Y(1) + + 7++ + .) = Y(2) + + 2-1.0 RETURN EMD С ۲C SURROUTINE DEDX (X, DF, JACT, F) CON54CH X0(2).8.01 D1MEN9104 X(1).P(2).DF(1,1) REG. JACT (2,1) DF(1,1)+2.0F(F(1)+JAGT(1,1)+*(1))+JACT(1,1)+ 9.0+(P(2)+JACT(2,1)+X(1))+JACT(2,1) L PETURI END REAL FUNCTION F (ALFA) REAL ALFA COMMON (X0.(2) (1) N COMMON, C51 (2) CALL FUNZ (ALFA, 2) F=Z RETURN END Sé С С SUPPOUTINE FUNZ (ALFA, Z) С C CALCULA LA FUNCION Z (XG-ALFANDSI) C COMMON REPORTS COMMON CEL (2) Z=0.0 DO 1 1-1.N 2=2+(X0(1)-ALF4:C31(1))++2 t Z#0.5×Z RETURN END

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MAGES HARDON, IN MARSHOPPIN, CASH SUSED UNBLIN IN ADEPATH. CODE PROGRAM (CARCISE) C ~ 10 . Ç PRODRAM SANUSO С 1.6 8 ĉ ESTE PROGRAMA CALCURA LA COLUCION DE UN SISTEMA ALGEORATOS NO LINERA, DREDETERMINADO DE LA FORMA I Ç F(X)+0 ¢ DONNER F, X Y O GUN WH YORKER BE DIMERSIBILIS, N Y P REFECCI (WHOLDE C CON HON. - MEDICO A ESTIMAS PROCESSARIO PROPERCY UNA TURCIAN OCULTION ¢ Z=Z(X) A CATREMIZOR, GIENDU Z UM NUMERO ACOU. Ę, EN RESUMENT, TENEROS EL SIGUIENTE PRODUCHA DE PROGRAMMATEMA-С LIC-II ē a (61)1 Z= 2 (31) х C. CUJETA A £ (X) #0 C D31EN3304 02131,01337(3),011(3),E814X1(3),X718(3),U(3), (Y(2) DEDENGED CERD (2) , XYO (2) , MEX (3) , P (2) REAL FORD, 340 (2,3) (4601 (3,2) (470 (5) DOMMON YOLD , H. H. COMPONE (20113) COMPANI N TABLE OPENIS, FROM DE MORDEL DATA1, STATIOS= INSUID PROTECT, CAPTOR DI DECEMI PRODUCTION, TECLES I REORD, 1011 10-1 LE CLEWARD TREE UPERIG FILL- (ERIMITIC) D E1.5;6, OPEN (FL) TO EFICONDOUCH TO ER01F (PR111), (47,56) WRITE (5.450) READ (*, 150) 10050 MAITECS, 1200 NOV60 IF TREASE, OT. 5. ON NOTION LET 10 GO TO A WRITE (6,750) READ (1, 100) NOTINA WRITEG, 1500 BORBA IF (PRICEMULTER, 1), ALF4.=1.0 984(6(6,10)) REODES, 2000 IN, N WRITE (6, 200) H.N. 10LX+1.0E+S T01.F=0.0001 PARP+C. P 1197,+200 KHAX+100 WRITE(6, 350) WRITESS, SOOD TOLY, FOUR, DAIT, MOX, KMAX HRCTE (4,360) REGD:*,4005 (20(1),1*1,14) WRITE(6,400) (SQ(2), I=1,N) WRITE(3,510) (00(1), 5+1,10 hilling (mar. 19 60321-0.10 0000011#0.0 CEL0121=6_0. 1.2 WRTTE(6,450) READCH, WOUL CAY (11, 1-1, 11) REPARTS AND LEFTER LALM

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1	CONTINUE	<u>.</u>			
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Referencias

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DIVISION DE EDUCACION CONTINUA FACULTAD DE INGENIERIA U.N.A.M.

DISENO CINEMATICO DE MAQUINARIA

GUIDANCE OF AXIALLY SYMMETRIC RIGD BODIES USING FIFTH-DEGREE OF FREEDOM REVOLUTE COUPLED MANIPULATORS

DR. JORGE ANGELES ALVAREZ

JUNIO, 1984.

GUIDANCE OF AXIALLY-SYMMETRIC RIGID BODIES USING FIFTH-DEGREE-OF-FREEDOM--REVOLUTE-COUPLED MANIPULATORS.

Jorge Angeles¹

Abstract

A kinematic model is constructed that allows the computation of all joint angles associated with fifth-degree-of-freedom revolute -coupled kinematic chains guiding an axially-symmetric rigid body through a set of prescribed configurations. Noreover, if velocity and acceleration specifications are introduced, the model also provides the first two derivatives of the joint angles. Finally, the applicability of the algorithm presented here, to the analysis of single-degree-of-freedom single-loop-6R closed chains is shown with an example.

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CONDUITE DE CORPS REGIDES AXIALEMENT SYIMETRIQUES AU MOYEN DE , MANIPULATEURS - À CINQ COUPLES ROTOÏDES

Jorge Angeles

Résumé

L'auteur presente un modèle cinématique permettant de calculer tous les angles associés aux cinq couples rotoïdes des chaînes cinématiques à cinq degrés de liberté, destinées à guider des corps rigides axialement symétriques au moyen d'un ensemble de configurations préétablies. En outre, si des specifications de vitesse et d'accélération y sont introduites, le modèle fournit aussi les deux premières dérivées des cinq angles. En fin, l'applicabilité de l'algorithme présenté ici à l'analyse des chaînes fermées à un degré de liberté comportant six couples rotoïdes est montré par moyen d'un exemple.

Nomenclature

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х:	n-dimensional vector over the real field
X:	mxn matrix over the real field
$\mathbf{x}_{i}, \mathbf{y}_{i}, \mathbf{z}_{i};$	Cartesian axes fixed to the ith link of the chain
[x] _i , x _i :	3x1 array containing the components of 3-dimensional vector
	x referred to X_i , Y_i , Z_i
[X] _i , X _i :	3x3 array containing the components of 3x3 matrix X referred
	to X_i , Y_i , Z_i .
tr (X):	the scalar invariant of x matrix X, $X_{11} + \dots + X_{nn}$
к.:	the axis of the ith revolute pair of the chain, coincident
	with Z _i .
e _i :	the angle of rotation of link $i+1$ with respect to link i , its
	sign being defined by the positive direction of Z_{i} .
a: i	the angle between R _ and R $_{\ell^{+}I}$, its sign being defined by the
	positive direction of axis $X_{\mathcal{L} \neq J}$, which is directed in turn
	along the common perpendicular to $R_{\vec{k}}$ and $R_{\vec{k}+\vec{l}}$, from $R_{\vec{k}}$ to
	R ₁₊₁ .
a.: 1	The distance between axes ${\tt R}_{i}$ and ${\tt R}_{i+1}$, hence always positive
٥. •	the coordinate of the intersection of axes $X_{\ell-1}$ and Z_{ℓ} in
	frame X _i , Y _i , Z _i
$[\mathbf{Q}_{i}, \mathbf{i}_{i+1}]_{i}$	Q_{j} : the matrix rotating axes $X_{j}^{}, Y_{j}^{}, Z_{j}^{}$ into an orientation
;	parallel pairwise to X_{i+1} , Y_{i+1} , Z_{i+1} , respectively, referred
	to X_i , Y_i , Z_i coordinates
[a_1,1]1. a	$_{i}$: the vector connecting the origins of X $_{i}$, Y $_{i}$, Z $_{i}$ and
	X_{i+1} , Y_{i+1} , Z_{i+1} , directed from the former to the latter,
r _A , r _B : -	in X, Y, Z, coordinates. the position vector of points A and B, respectively, in the -4
	specified configuration, measured from the origin of X_1 , Y_1 , Z_1 .
v _A , v _B :	the velocity of points A and B, respectively
^а д, а _в :	the acceleration of points A and B, respectively

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Introduction

Industrial robot manipulators in use are frequently required to perform tasks involving the guidance of rigid bodies with axial symmetry, e.g. turned workpieces, painting nozzles, are welding pistols, etc. Most of the time, however, these tasks are realized with sixth-degree-of--freedom manipulators. Since the guidance of axially-symmetric rigid bodies does not involve their orientation about the axis of symmetry, it seems natural to perform these tasks with a fifth-degree-of freedom manipulator. If, on the other hand, a sixth-or greater-degree-of--freedom manipulator is to be used anyway, the redundant degree of freedom can be used to optimize some given performance index. Hence the interest to establish a kinematic model allowing the computation of the joint angles and their time derivatives, associated with multiple-degree-of-freedom manipulators in the presence of incomplete specification of the hand motion.

Following the approach introduced in [1], the set of desired equations is derived regarding the problem as one involving the guidance of a segment of a rigid line, by specifying the displacement, velocity and acceleration of any two points of that segment. Of course, these specified variables should meet the compatibility condition guaranteeing that the distance between these points is preserved throughout any physically possible motion.

The kinematics of rigid segments has been considered in [2] regarding only the velocity and acceleration analyses. A complete account of the kinematics of rigid segments, lines and points is given in [3]. 6 Furthermore, an analysis of SR open kinematic chains to guide axially-symmetric rigid bodies appeared in [4]. The approach followed in this paper differs from these, being aimed at devising a real-time implementable algorithm oriented towards the computer-control of SR robot manipulators, applied to the guidance of axially symmetric rigid bodies.

5,

Displacement analysis

In this Section, reference is made to the fifth-degree-of-freedom kinematic chain depicted in Fig 1. Moreover, the notation and method of Denavit and Hartenberg [5, pp 343-355] is applied throughout. Thus, a set of Cartesian coordinate axes $\{X_{\mathcal{L}}, Y_{\mathcal{L}}, Z_{\mathcal{L}}\}$ is attached to the i th link. According to Denavit and Hartenberg's notation, axis Z_{χ} is placed along the axis of pair R_{χ} , its positive direction denoting the direction in which angle θ_{j} is measured. Axis $X_{\hat{\mathcal{L}}}$ is defined as the common perpendicular to the axes of R_{i-1} and R_i , directed from the former to the latter. Notice that there is absolute freedom to chose X_j as any line in a plane per pendicular to Z. The orthogonal matrix rotating axes labelled λ into an orientation paralel to those labelled *i*+1, referred to axes i, is represented as $[Q_{i, i+1}]_{i}$ or as Q_{i} , for compactness. Finally, the vector connecting the origins 0_j and 0_{j+1} , respectively, directed from the former to the latter, in *i*-coordinates, is represented $[a_{i, i+1}]_{i}$, or as a_{i} , for compactness. According to the as nomenclature, these are²

$$\begin{bmatrix} Q_{i, i+1} \end{bmatrix}_{i} = \begin{bmatrix} c\theta_{i} & -s\theta_{i}c\alpha_{i} & s\theta_{i}s\alpha_{i} \\ s\theta_{i} & c\theta_{i}c\alpha_{i} & -c\theta_{i}s\alpha_{i} \\ 0 & s\alpha_{i} & c\alpha_{i} \end{bmatrix}$$
(1)

$$\begin{bmatrix} a_{i, i+1} \end{bmatrix}_{i} = \begin{bmatrix} a_{i} c \theta_{i}, a_{i} s \theta_{i}, b_{i} \end{bmatrix}^{\mathsf{T}}$$
(2)

Let line AB of Fig 1 be the axis of symmetry of the rigid body, not shown there, that is meant to be guided. Furthermore, let r_A and

7.

[.]Throughout, c() $\equiv \cos($), s() $\equiv \sin($).

 $r_{\rm B}$ be the prescribed values of the position vectors of A and B. respectively, letting x_j and y_j be the synthesized values of the same position vectors. Thus, $r_{\rm A}$ and $r_{\rm B}$ are known constants, whereas $x_j = x_j$ (0), $y_j = y_j$ (0), are functions of 0, where 0 denotes the 5-dimensional vector whose lith component is θ_{i} . Vectors x_j and y_j can be computed recursively as follows:

$$x_4 = a_4$$
 (3a)

$$x_k = a_k + Q_k x_{k+1}, \quad k = 3, 2, 1$$
 (3b)

$$y_5 = a_5$$
 (4a)

$$y_k = a_k + Q_k y_{k+1}, k = 4, 3, 2, 1$$
 (4b)

_ The displacement equations are, then,

$$x_1 = r_A, \quad y_1 = r_B \tag{5}$$

Now, the six-dimensional vector f is defined as

$$\mathbf{f} \equiv \begin{bmatrix} \mathbf{f}_1 \\ \mathbf{f}_2 \end{bmatrix}$$

f f_1 and f_2 being the three-dimensional vectors defined in turn as

$$f_{ij} = x_1 - r_A, \quad f_2 = y_1 - r_B$$
 (6)

Thus, the displacement equations lead to the following sixth-order nonlinear algebraic system in five unknowns:

$$f(\theta) = 0 \tag{7}$$

This system can be solved using Newton-Raphson's method [6, pp 248-253]

as follows: Let θ^{0} be an initial "guess" for the solution. Then generate the sequence $0^{J}, \ldots, 0^{k}, \ldots$ which, if converges, will do so quadratically [6, pp 222-226]. This sequence is generated recursively from Newton-Raphson's iterative scheme:

$$e^{k+1} = e^k + \Delta e^k \tag{8}$$

where $\Delta \theta^k$ is computed from the Taylor expansion of $f(\theta)$ at $\theta^{=}\theta^k$, which leads to

$$J(\mathfrak{o}^k) \Delta \mathfrak{o}^k = -f(\mathfrak{o}^k) \tag{9}$$

Equation (9) is a sixth-order linear algebraic system in five unknowns, its Jacobian matrix J being 6x5. Thus J cannot be inverted properly speaking. However, from the nature of the problem, out of the six scalar equations appearing in (9), only five are independent. In fact, no matter what the value of θ^k is, $f(\theta^k)$, and hence $-f(\theta^k)$, lies in the range of J, if the motion of points A and B is not to violate the rigidity condition. As a matter of example, consider the one-degree-of-freedom 6R linkage³ depicted in Fig 2 for the particular values $\theta_1 = -\theta_3 = \theta_5 = 120^\circ$, $\theta_2 = \theta_4 = 0^\circ$. The motion of this linkage can be analysed as one of an open 5R kinematic chain, as shown in Example 1. For that configuration,

 $J = \begin{bmatrix} 0 & 0 & \sqrt{3} & 0 & 0 \\ -2 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & -1 & 0 \\ 0 & 0 & \sqrt{3} & 0 & 0 \\ -1 & 0 & 0 & 0 & 1 \\ 0 & -1/2 & 0 & -1/2 & 0 \end{bmatrix}$

ч.

⁷³ A wire model of this is for sale commercially under the trade work HEXIFLEX^{TR}.

Now assume that the link connecting 0_6 with 0_1 is removed, thus obtaining an open chain. Moreover, it is desired to compute the angle increments $\Delta 0_2$ ($\ell = 1, ..., 5$) for the following value of Δf :

$$\Delta f = \begin{bmatrix} \Delta x_j \\ \overline{\Delta} y_j \end{bmatrix}; \ \Delta x_j = \begin{bmatrix} 0 \\ 0 \\ \zeta \end{bmatrix}, \ \Delta y_j = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

 Δx_j being specified parallel to the Z_j axis, otherwise arbitrarily.

This motion clearly complies with the rigid-body condition, for $\Delta y_1 - \Delta x_1$, the displacement increment of B with respect to A, is perpendicular to AB, as it should be if the distance between A, and B is to remain constant. Hence Δf lies in the range of J. Eq (9) thus yields, for the assumed values of θ_j ,

$$\Delta \theta_1 = \Delta \theta_2 = \Delta \theta_5 = 0, \quad \Delta \theta_9 = -\Delta \theta_4 = 2.4$$

If, on the other hand, Δf is specified as

$$\Delta \mathbf{f} = \begin{bmatrix} \Delta \mathbf{x}_1 \\ \Delta \mathbf{y}_1 \end{bmatrix}; \quad \Delta \mathbf{x}_1 = \begin{bmatrix} \boldsymbol{\xi} \\ \mathbf{0} \\ \mathbf{0} \end{bmatrix}, \quad \Delta \mathbf{y}_1 = \begin{bmatrix} \mathbf{0} \\ \mathbf{0} \\ \mathbf{0} \end{bmatrix}$$

which clearly violates the rigid-body condition, then eq (9) leads to an inconsistent equation system producing, from the first equation

$$\Delta \theta_3 = \frac{\sqrt{3}}{3} \xi$$

and, from the fourth equation,

$$\Delta 0_3 = 0$$

which contradicts the former one . This is due to the fact that Δf does not lie in the range of J because it does not comply with the rigid-body condition.

Thus, except for singular configurations, i.e. those leading to a rank -defficient Jacobian matrix - one for which rank (J) $<5-, \Delta u^k$ can be solved for from eq (9). Since this equation contains, in general, 5 consistent linearly independent scalar equations, but one does not know which one is redundant, 'Gauss' algorithm or, equivalently, the LU decomposition [7, pp 27-33] cannot be applied to it directly. The solution proposed here is to regard eq (9) as a 6x5 overdetermined linear system and compute its "least-squere" solution. This will be, in fact, *its solution*, given the consistency of the involved system of equations. There are several ways of computing the least-square solution of eq (9), but the one preferred by the author is using Householder reflections [8], which produces implicitly the pseudo-inverse [9, pp 103-113] of J, J^I. Thus

$$\Delta \sigma^{k} = -\sigma^{I} (\theta^{k}) f (\theta^{k})$$
 (10a)

with

$$J^{I}(\theta^{k}) = [J^{T}(\theta^{k}) J(\theta^{k})]^{-1} J^{T}(\theta^{k})$$
 (10b)

Pseudo-inverses have been already used in connection with the analysis of multiple-degree-of-freedom manipulators [10], though their use is not very popular.

11.

In this section, a system of six linearly-dependent consistent equations is derived that allows the computation of \dot{e} , given the velocities of A and B, henceforth denoted by the 3-dimensional vectors v_A and v_B , respectively. These are, clearly

$$\ddot{\mathbf{x}}_1 = \mathbf{v}_{\mathbf{A}}, \ \dot{\mathbf{y}}_1 = \ddot{\mathbf{v}}_{\mathbf{B}} \tag{11}$$

with

$$\dot{x}_{1} = \frac{\partial x_{1}}{\partial \theta} \dot{\theta}, \quad \dot{y}_{1} = \frac{\partial y_{1}}{\partial \theta} \dot{\theta}$$
(12)

which thus lead to

$$J(0) \hat{0} = v$$
 (13)

with the 6-dimensional vector v given by

$$v = \begin{bmatrix} v_A \\ -v_B \end{bmatrix}$$
 (14)

Recalling eq (10), 8 can then be obtained as

$$\dot{\theta} = J^{I} v$$
 (15)

 J^{I} being given as in eq (10b). It is pointed out here that, since J^{I} was already computed in computing 0, it need not be recomputed. In fact, using subroutines HECOMP and HOLVE [11] to obtain the least-square solution (10), HECOMP is first applied to J in order to rendering it upper triangular. Next, HOLVE is applied to the transformed J and the right-hand side of eq (9) to produce solution (10). This means that HECOMP is applied only once, whereas HOLVE twice. Host of the operations involved in the aforementioned procedure 13 are performed in HECOMP; thus, \tilde{O} is obtained at virtually no additional cost.

Acceleration analysis

Next the computation of $\ddot{\theta}$ is outlined. This is clearly obtained by differentiating eqs (11) with respect to time:

$$X_{j} = a_{A}, \quad Y_{j} = a_{B} \tag{16}$$

where a_A and a_B denote the accelerations of A and B, respectively. Furthermore, from eqs (12),

$$\ddot{x}_{j} = \frac{\partial x_{j}}{\partial \theta} \ddot{\theta} + \left(\frac{\partial^{2} x_{j}}{\partial \theta^{2}} \ddot{\theta}\right) \dot{\theta}$$
(17)

with a similar one for \ddot{y}_{f} . Substituting eq (17) and its \ddot{y}_{f} counte<u>r</u> part in eqs (16), one obtains

$$J(\theta) \ddot{\theta} = a$$
 (18)

with

÷

.

Hence, 8 is obtained as

$$\ddot{\mathbf{B}} = \mathbf{J}^{\mathbf{I}} \left(\mathbf{0} \right) \mathbf{a}$$
 (20)

and, as said before in connection with $\dot{\theta}_{*}$ at virtually no additional cost.

Example 1. Analysis of a 6R single-degree-of-freedom closed kinematic 15chain. The parameters defining the architecture of the chain are given below. Regarding links 6 and 5 as the fixed and the input link, respectively, determine the histories $\theta_i(t)$, $\theta_i(t)$, $\theta_i(t)$, $\theta_i(t)$, $\theta_i(t)$, $\theta_i(t)$, $\theta_i(t) = 0$, $\theta_i = 1$ rad/s and $\theta_i = 0$. The parameters are:

$$a_i = 1, b_i = 0, \alpha_i = 90^\circ, i = 1, \dots, 5, \alpha_6 = -90^\circ.$$

Fig 2 shows the kinematic chain described by the foregoing parameters, for the configuration

$$\theta_1 = -\theta_3 = \theta_5 = 120^\circ, \quad \theta_2 = \theta_4 = \theta_6 = 0^\circ$$

Solution:

Due to the symmetries of this linkage,

 $\theta_{1} = -\theta_{3} = \theta_{5}, \quad \theta_{2} = -\theta_{4} = \theta_{6}$ Vectors $x_{l_{2}} \quad (l_{2} = 4, 3, 2, 1)$ are readily computed as $x_{4} = a_{3} = [c_{4}, s_{4}, 0]^{T} = [c_{2}, -s_{2}, 0]^{T}$ $x_{3} = a_{3} + \theta_{3} \quad x_{4} = [c_{1}(1+c_{2}), -s_{1}(1+c_{2}), -s_{2}]^{T}$ $x_{2} = a_{2} + \theta_{2} \quad x_{3} = \begin{bmatrix} c_{2} \quad [1+c_{1}(1+c_{2})] - s_{2}^{2} \\ s_{2} \quad [1+c_{1}(1+c_{2})] + c_{2}s_{2} \\ -s_{1}(1+c_{2}) \end{bmatrix}$ $x_{1} = a_{1} + \theta_{1} \quad x_{2} = \begin{bmatrix} c_{1}\{1+c_{2}(1+c_{1}(1+c_{2})] - s_{2}^{2} \} - s_{1}^{2}(1+c_{2}) \\ s_{1}\{1+c_{2}(1+c_{1}(1+c_{2})] - s_{2}^{2}\} + c_{1}s_{1}(1+c_{2}) \\ s_{2}[1+c_{1}(1+c_{2})] + c_{2}s_{2} \end{bmatrix}$

Vectors y_k (k = 5, 4, 3, 2, 1) are similarly computed as

$$y_{5} = a_{5} = [c_{5}, s_{5}, 0]^{T} = [c_{1}, s_{1}, 0]^{T}$$

$$y_{4} = a_{4}^{'} + Q_{4} y_{5} = [c_{2} (1+c_{1}), -s_{2}(1+c_{1}), s_{1}]^{T}$$

$$y_{3} = a_{3} + Q_{3} y_{4} = \begin{bmatrix} c_{1} [1+c_{2} (1+c_{1})] - s_{1}^{2} \\ -s_{1} [1+c_{2}(1+c_{1})] - c_{1} s_{1} \\ -s_{2}(1+c_{1}) \end{bmatrix}$$

$$y_{2} = a_{2} + Q_{2}y_{3} = \begin{bmatrix} c_{1}c_{2}[c_{1}+1+c_{2}(1+c_{1})] - s_{2}^{2}(1+c_{1}) \\ c_{1}s_{2}[c_{1}+1+c_{2}(1+c_{1})] + c_{2}s_{2}(1+c_{1}) \\ - s_{1}[c_{1}+1+c_{2}(1+c_{1})] \end{bmatrix}$$

$$y_{1} = \begin{bmatrix} c_{1}\left\{1+c_{1}c_{2}[c_{1}+1+c_{2}(1+c_{1})] - s_{2}^{2}(1+c_{1})\right\} - s_{1}^{2}[c_{1}+1+c_{2}(1+c_{1})] \\ s_{1}\left\{1+c_{1}c_{2}[c_{1}+1+c_{2}(1+c_{1})] - s_{2}^{2}(1+c_{1})\right\} + c_{1}s_{1}[c_{1}+1+c_{2}(1+c_{1})] \\ s_{2}c_{1}[c_{1}+1+c_{2}(1+c_{1})] + c_{2}s_{2}(1+c_{1}) \end{bmatrix} \end{bmatrix}$$

The displacement equations are now obtained equating x_j with r_A and y_j with r_B , r_A and r_B being the position vectors of points A and B, respectively, in X_j , Y_j , Z_j coordinates. Hence

$$c_1 \left\{ 1 + c_2 \left[1 + c_1 \left(1 + c_2 \right) \right] - s_2^2 \right\} - s_1^2 \left(1 + c_2 \right) = - (1 + c_2)$$
 (1)

$$s_1 \left\{ 1 + c_2 \left[1 + c_1 (1 + c_2) \right] - s_2^2 \right\} + c_1 s_1 (1 + c_2) = 0$$
 (ii)

$$s_2 [1+c_1(1+c_2)] + c_2 s_2 = s_2$$
 (iii)

$$c_{1}\left\{1+c_{1}c_{2}\left[c_{1}+1+c_{2}(1+c_{1})\right] - s_{2}^{2}(1+c_{1})\right\} - s_{1}^{2}\left[c_{1}+1+c_{2}(1+c_{1})\right] = -1 \quad (iv)$$

$$s_{1}\left\{1+c_{1}c_{2}\left[c_{1}+1+c_{2}(1+c_{1})\right] - s_{2}^{2}(1+c_{1})\right\} + c_{1}s_{1}\left[c_{1}+1+c_{2}(1+c_{1})\right] = 0 \quad (v)$$

$$c_{1}s_{2}\left[c_{1}+1+c_{2}(1+c_{1})\right] + c_{2}s_{2}(1+c_{1}) = 0 \quad (vi)$$

From eq (iii), with $s_2 \neq 0$ and $c_1 \neq -1$.

$$(1+c_1)(1+c_2) = 1$$
 (1/0)

which is the input-output equation, valid for $c_1 \neq -1$. For $c_1 = -1$, $s_2 = 0$.

Now, let { } and [], be the corresponding brackets appearing in eqs (1-1ii). { }, and [], are correspondingly defined for eqs (iv - vi). If (1/0) is substituted into these, one has

$$[] = \frac{1+2c_1}{1+c_1}, [] = 1$$
 (vii)

$$()_{j} = -\frac{c_{j}}{1+c_{j}}, \qquad = -c_{j} \qquad (viii)$$

Substitution of relations (vii is viii) into eqs (i - vi) renders both sides identical, thus verifying the validity of the input-output equation. From (I/O) it is clear that neither θ_1 nor θ_2 can attain the value π . In fact, a mobility analysis yields the following mobility ranges:

$$-120^{\circ} \leq \theta_{1} \leq 120^{\circ}, \lambda \neq 1, \dots, 6$$

The program implementing the algorithm presented in this paper was tested with this linkage, its input-output equation (I/O) being plotted in Fig 3a. The program output was compared with the foregoing closed-form solution and the displacement, velocity and acceleration errors were plotted as shown in Figs 3b, 3c and 3d, respectively.

As in the case of seven-link chains [1], the displacement, velocity and acceleration errors grow unbounded as the Jacobian matrix approaches a singularity which, in this case, manifests itself as a deficiency of its rank. This occurs at singular configurations of the linkage under study, i.e. at dead-point positions of its input link, $\theta_A = \theta_7 = \pm 120^\circ$.

Example 2. Synthesis of a spatial trajectory.

Given the SR manipulator whose parameters appear in Table 1, determine the histories $\theta_{i}(t)$, $\dot{\theta}_{i}(t)$, for $0 \le t \le T$, i = 1, ..., 5, in order to guide line AB, as shown in Fig 4, through a given spatial trajectory. The location of point B, in fifth-link coordinates, is given by $[r]_{5} = [a_{5}, 0, 0]^{T}$, whereas the trajectory is chosen as the upper branch of the intersection of the sphere (S) with the cylinder (C) given below:

$$(x - a)^{2} + (y + b)^{2} + z^{2} = 15a^{a}$$
 (5)
 $x^{2} + z^{2} = a^{2}$ (C)
upper branch: $y > b$ (B)

the said surfaces being given in fixed coordinates, i.e. in X_j , Y_j , Z_j coordinates. It is required, moreover, that

```
x = -\alpha \cos \beta
y = \alpha \sin \beta
```

with β being chosen as a smooth-enough function of time, i.e. one for which at least $\dot{\beta}$ and $\ddot{\beta}$ be continuous functions of time, in the interval [0,T]. Additionally, the following is imposed on β :

$$\beta(0) = 0, \beta(T) = 2\pi, \beta(0) = \beta(T) = 0, \beta(0) = \beta(T) = 0$$

Such a β function can be readily synthesized using spline functions, as shown in [12, 13]. The following parameter values were assumed:

 $\alpha = 300 \text{ mm}, b = 2 220 \text{ mm}, T = 60 \text{ s}$

The orientation of AB was specified as follows. Let unit vectors

 e_{a} , e_{n} , e_{b} denote the tangent, normal and binormal vectors of the trajectory. Then AB is to be orientated so that:

 $r_{\rm B}$ - $r_{\rm A}$ = $a_5 e_n$

Table 1

(lengths in mm, angles in degrees)

a ₁ = 0	6 ₁ = 0	a _j = 90
$a_{2} = 0$	6 ₂ = 479	° ₇ ≈ 90
a ₃ = 0	6 ₃ = 0 ·	α ₃ = 90
$a_4 = 35.3$	b ₄ = 1016	α ₄ ≈ -90
a ₅ = 146	δ ₅ = 0	α ₅ = -90,

These values were taken from. The sixth-degree-of-freedom manipulator described in [14], from which the sixth revolute pair was removed.

Solution:

The reference configuration was chosen to be the following:

 $\theta_1 = 90^\circ, \ \theta_2 = 180^\circ, \ \theta_3 = 90^\circ, \ \theta_4 = \theta_5 = 0^\circ$

In order to guide line AB from its position in the reference configuration to the initial configuration along the prescribed trajectory, determined by B=0, continuation was used, as donde in [1]. To this end, point 8 was made to trace a straight path between its two positions, that in the reference one, B^{θ} , and that in B=0, B^{T} . This path was divided into 10 segments in order to ensure that the initial "guess" for the Newton-Raphson procedure lie close enough to the solution sought, which guarantees and accelerates its convergence. In fact, four interations were needed, at most, in order to reach convergence, in this stage.

Along the prescribed trajectory, one whole kinematic analysis was performed every 0.5s, i.e. 120 points on the given trajectory where fully analyzed. At virtually all of these points, convergence was reached after 3 iterations, which was the largest number of iterations required, for a tolerance of 10^{-6} imposed on $\Delta \theta$. At a few points about $\beta=0$, convergence was reached after only two iterations.

Computed values of θ_{i} , $\dot{\theta}_{i}$ and $\ddot{\theta}_{i}$ (*i*+1,...,5) are plotted in Figs 5-9. From these figures it becomes apparent that the smoothness of function β_{i} synthesized with the aid of a cubic periodic spline, is reflected in the smoothness of the θ_{i} functions obtained.

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Conclusions

The method presented here, aimed at the complete analysis of five-link revolute-coupled kinematic chains, differs basically from other methods intended to solve the same problem. Contrary to the usual practice of analyzing open kinematic chains (man ipulators) as closed ones by introducing a "fictitious" closing link, the approach followed here is oriented to the analysis of open chains, regarding closed ones (linkages), as particular cases of the former. It resorts to efficient numerical methods of solution of nonlinear overdetermined algebraic systems. The system obtained here is derived disregarding the orientation of the body meant to be quided, while intending to position two of its points, which is also an original approach. The results show quick convergence, which suggests the applicability of the algorithm, and the program implementing it, to the real-time kinematic control of robot manipulators of the topology assumed here, i.e. 5R, but of arbitrary architecture otherwise ...

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Fig 1 Architecture of a general fifth-degree-of-freedom revolute-coupled manipulator

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Fig 2 Layout of a 6R-single-degree-of-freedom linkage

F2

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Fig 4 Fifth-degree-of-freedom manipulator.



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à, (5")+ ë, (5") de (rad) -2 FIG 9



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OPTIMAL SYNTHESIS OF TRANSLATIONS - FOLLOWER CAN MECHANISM WITH PRESCRIBED FUNCTIONAL CONSTRAINTS

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JUNI0,1984.

OPTIMAL SYNTHESIS OF TRANSLATING-FOLLOWER CAN DECHAPISHS WITH DEDSCRIDED PERCEIDENCE CONSTRAINTS

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Abstract

A covel approach is presented, that allows the computer synthesiz of which mustize cat profiles, which are to produce a given rollower-displacement programs while exceting preactibed functional constraints. These can be imposed either on the previous angle or on the eccentricity of the contact point depending upon the type of follower that is being synthesised. The paper is limited to the synthesis of cum profiles for translating roller followers, but the mate approach has been raccessfully applied to flat-face followers, as reported proviously. The results obtained for the example included show the applicability of the procedure to the automatic decign of this type of can follower mechanisms.

Keywords; optimal synthesis, can mechanisms, optimum design, automatic design

1. INTRODUCTION

The synthesis of can bechanises to produce a given displacement programme of the follower is divided into two stages [1]. The first one is meant to produce an caty-to-handle smooth-enough function of the angle of rotation of the can disk, representing the displacement of the follower. The second one deals with the synthesis of the can profile. The synthesis of the displacement programme has been based traditionally on a limited set of functions giving rise to parabolic, harmonic, cycloidad, trapezoidal and polynomial notions. Based on the last type, Thompoulus and Knowles [2] proposed a novel method of displacement-programme synthesis using linear programming, but it is not until very recently that a totally new approach has been introduced using spline tonerions [3,4].

As to the profile synthesis, the traditional approach, based on graphical methods, has been abandoned in favour of numerical methods, that have been called for given the extensive use of computers and NC-machine tools [5]. The literature on graphical methods of profile synthesis is rather broad, as can be seen in [6], that on computer-oriented methods becoming abounding, particularly since the 1970's [6-14]. The approach introduced in this paper, regarding the optimal synthesis of can profilem, is that of mathematical programming, but, as seen in the discussion, it does not require the application of sophisticated and time-consuming optimisation methods.

2. THE TOLLOWER-DISPLACEMENT PROGRAMME

In this Section a method is described that allows to synthesise follower-displacement programmes, statting from an acceleration programme that is specified over a discrete set of values, $\{\psi_j\}_{j=1}^{n}$ of the angle $|\Psi|$ measuring the rotation of

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the disk can involved. Let the displacement of the follower by given by c+i(1), where c is a constant representing the minimum value of that displacement, O(q) being a positive definite function whose minimum value is zero, its maximum val we being h, the rise of the follower. Contant 2 is computed from considerations concerning the profile synthesis, as will be shown later, the objective of this Section being the synthesis of O(9). This will be performed using a section of a basic function t(y) which is synthesised using spline functions as shown in [4]. 1(y) is evathesised by prescribing a set of deceleration values, $\{1^{n}c, j\}_{j=1}^{n}$ so that $1_{j=1}^{n}c_{j}^{n}$, in order to obtain a periodic function $\tau(j)$. To this end, periodic splines are used. Thus, the set $\{\tau_j\}_{j=1}^n$ $\{\tau_j(\phi_j)\}_{j=1}^n$ is computed by solving a linear system of equations.For instance, 1". 2-1,2,..., a can be specified as having an harachie distribution, van ishing at ψ_1 and ψ_2 , the extremes of the interval of interest, which can be specified in turn as [0.2m]. Horeover, t" is also specified to vanish at inner points of this futerval. in order to the sure that the displacement function will have a zera acceleration at its higher point, the comfor which the displacement attains the value off. Furthermore, in order to meet the condition that the velocity be a continuous function of Q. symmetrics are introduced that guarantee the sanishing of the velocity at the two dwell intervals of the follower. Thus the is specified such that $\tau_1^{\alpha}=0$, the visite function, shown in Fig. 1, being assumed to be periodic and symmetric with respect

A lf c+v(k) represente the displacement of the follower them its velocity and acceleration are given by 0'(.), and v(.); 'to"(.), respectively, k and \tilde{v} being the velocity and the acceletation of the can disk.

to line $\psi_{0}\psi_{0}$ ==""", the midpoint of the interval [0,2*]. If the function is further assumed to be odd with respect to lines $\psi_{-}\psi_{-}\pi/2$ and $\psi_{-}3m-2=3\pi/2$, then "" will sho vanish at B and C. Horeover, this way, t" will sho vanish at A, B and C. The basic function to produces (ψ) can be obtained, in turn, from the curve shown in Fig.1 by shifting the ψ_{-} axis to the line $\tau_{-}\tau_{0}$. A proper way to prescribe t" is as follows:

1"-ksiny, i-1,...,m

 $\psi_{i} = \frac{\pi}{2(n-1)}$ (i-1) (1b)

 (1_{4})

where constant k is chosen so as to normalize τ_i i.e. such that $\tau_0 \approx 1/2$. Using a periodic spline, with $m \approx 11$, the curve appearing in Fig.1 was obtained with k=0.64191.

The displacement programme for the follower of a can nechanism is now assumed to have a lower dwell in the interval $[0, \frac{1}{2}]$, then a rise in the interval $[\psi_1, \psi_2]$, next a higher dwell in the interval $[\psi_2, \psi_3]$ and, finally, a return in the interval $[\psi_3, 2\pi]$. Let

$$\Delta \psi_1 - \psi_1 + \Delta \psi_2 - \psi_2 - \psi_1 + \Delta \psi_3 - \psi_3 - \psi_2 + \Delta \psi_4 - 7 = -\psi_3 (2a)$$

The foregoing displacement programme can now be synthesized from the basic function t vs ψ of Fig.1 by properly scaling separately its intervals $[0,\pi]$. To this end, redefine ψ as

$$\psi := \frac{n\psi_2}{2}\psi , \quad 0 \le \psi \le x \qquad (2b)$$

$$\psi := \frac{n\psi_4}{\pi} (\psi - \eta), \ \pi \le \psi \le 2\pi \qquad (2c)$$

Now Cunction O(¢) is defined as

A subprogram was written that allows the aucommutic synthesis of follower-displacement programmes by inputting the following data: ψ_{+} , ψ_{+} , ψ_{+} , and h. Its output is $\phi(\phi)$, $\sigma'(\phi)$ and $\sigma''(\psi)$. the procedure introduced to compute the basic function $\tau(\psi)$, on the other hand, need not be changed if an acceleration distribution different from the one given in (188b) is chosen. In fact, only the set $(T^{\mu}_{i})^{\mu}_{i}$ has to be changed as data input to the spline subroutines synthesising function $\tau(\psi)$. Moreover, if the follower is of the translating type, O(\$) and hence c and h will have units of length. They will be measured in radians if the follows, is of the oscillating Lype. The synthesis of the cam profile for trans lating followers of the roller type is next discussed. That of the cam profile for translating followers of the flat-face type is reported elsewhere [15].

3. OFTIMAL SYNTHESIS OF CAM PROFILES FOR TRANSLATING MOLLER FOLLOWERS

and hence

In this Section 12 is assumed that the displacement programme is specified up to a constant t, i.e. function $O(\psi)$ is known, whereas constant c is to be determined. Let the displacement of the follower be denoted by $u(\psi)$, i.e.

$$\mathbf{s}(\boldsymbol{\psi}) = \mathbf{c} + \sigma(\boldsymbol{\psi})$$
 (4)

$$=^{1}(\psi) - J^{*}(\zeta), = L^{*}(\zeta) - J^{*}(\zeta)$$
 (5)

Reference is note now to Fig.2, showing the layout of a case with a translating knite-cdge to lower. The interest of synthesising such a case profile lies in the fact that this is the pitch curve for a roller followers taking identical (b). In Fig.2, line P is the path of the fallower, lines T and N are the tangent and the nonmal to the case profile, given by its palat equation p=p(0), at the contact point Q. Horeaver, lines OF and OC are fixed to the trate and to the case disk, respectively, a being the eccentricity at the follower path. Angle u is the pressure angle. From the geometry of Fig. 2,

$$|\psi\rangle + \rho(\phi) + i\mu(0)\psi\rangle$$
 (6.4)

α=υ+⊈+;−π

$$\phi = \tan^{-1} \frac{\rho(0)}{\rho'(0)}$$
 (6d)

In order to keep the mechanical advantage of the mechanism within acceptable limits, a is usually bounded properly. Since a can be either pasitive or negative, its absolute value is bounded as

where D_{i} is normally chosen to lie "close" to 0°, for a "value of 90° would render the mechanical advantage zero. A value that is widely accepted is 30°. Finding the value of ψ at which |h|| attains its extrema is not so simple, for this function is not differentiable at the origin, where it attains its minimum. Hence a different even function of a that be smooth enough has to be extremised. A good candidate is cosh. Using relations (6c6d), this turns out to be

$$\frac{\cos \omega = \frac{s(\phi)}{\left\{ \left[\frac{1}{2} \left(\phi \right) - c \right]^2 + s^2(\phi) \right]^{\frac{1}{2}}}$$
(8)

Now cose will be kept within bounds as

$$cosa 2 cosa - c (9)$$

1.*.

where

The extrems of (8) are now found by receiving its derivative with respect to ζ . This is readily obtained as*

$$\frac{\mathrm{deoso}}{\mathrm{d}\psi} + \frac{(s^3 - e)[(s^3 - e)s^3 - ss^3]}{\{(s^3 - e)^2(s^3)\}^2} \tag{11}$$

which vanishes under either of the next two conditions:

(12a)
(11)
$$(s^{1}-e)s^{1}-ss^{n}$$
 or $\frac{s^{1}-e}{s} = \frac{s^{n}}{s^{1}}$ (12b)

Condition i) leads to the maximum +), as can be readily verified from eq(8), whereas condition ii) leads to the minimum c_{H} . Let τ_{0} be the value of ψ , not an yet determined, producing the minimum. From eq(8), then

* Henceforth, a prime on a variable means its derivative with respect to ψ, i.e. s'ids/dφ. Siuj larly, s"=d²s/dφ.

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with

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$$\frac{|v_{i}|^{-\alpha}}{v_{0}} = \tan_{H}$$
(13)

with s_{i+1} , s_{i+1}^{+} and s_{i+1}^{+} defined as the values attained by $s(\psi)$, $s_{i+1}^{+}(\psi)$ and $s_{i+1}^{+}(\psi)$ at $\psi = \psi_{\sigma}$. Relation (13) with already obtained by Chicurel (17). Since $s_{i+1}^{+}(\psi) = s_{i+1}^{+}(\psi) = s_{i+1}^{+}(\psi)$, eqs(12b) and (13) lead to $s_{i+1}^{+}(\psi) = s_{i+1}^{+}(\psi) = s_{i+1}^{+}(\psi)$.

$$\left|\frac{\partial}{\partial r} \frac{(2\pi)}{(2\pi)}\right| = \tan \frac{1}{2}$$
(14a)

$$\left[\left[\mathbf{o}^{n}(\varphi_{g}) \right] + t \operatorname{ang}_{H} \left[\mathbf{o}^{*}(\varphi_{g}) \right] \right]$$
(14b)

Substitution of eqs(455), evaluated at $\varphi \circ \psi_{P_1}$, into eq(13) and the assumed positive definiteness of $O(\psi)$ and c lead to

οr

$$[o_{0}^{*}-v] - \operatorname{cano}_{H}(c+J_{0})$$
 (15a)

On the other hand, during the rise phase the follower velocity is positive, as shown in Fig.3. From those plots it can also be seen that eq(14b) holds at two distinct values of σ^n , namely σ^n_i and σ^n_i . The positive value leads to $\sigma^n\sigma^1$ whereas the negative one leads to $\sigma^n\sigma^1$. Substitution of these values into eq(15s). together with the condition

c > 0 (15b)

define the set of values along which condition (10) holds during the rise phase. This is plotted in Fig.4. In that figure, this set is formed by the two branches of lines 1 and 2 making an angle og with wither direction of the claxis, the dashed section being excluded, for it violates (15b).

Since cost also attains stationary values at the lowest follower position (i.e. during the dwell phase D_1), the following must hold (or any set of values $\{c,e\}$:

Eq(16) is also plotted in Fig.4. From that figure, it is clear that points f and 0 satisfy eqs(15abb and 16) simultaneously. Taking the radius of the base circle, r, as a measure of the area, which is plousible because r is the lowest value of p(0), the distance OP leads to the mininum value of r.

To this end it was assumed that Ψ_0 was known. In fact, from the foragoing discussion, such value of Ψ_0 satisfying eq(14b), corresponds to the one which zeroes the function

$$f(\psi) = \sigma''(\psi) + t_{\text{ADD}_{\omega}}\sigma^*(\psi) \tag{17}$$

This function is plotted in Fig.5.

Values of ψ_0 for different pressure angles were obtained and plotted in Fig.6. From that plot the designer can select the appropriate percentage for calculating ψ_0 for the problem at hand.

Once the value of ψ_0 has been determined one can proceed to determining the values of σ_0 , σ'_0 and σ''_0 and thence the values of c_0 and c_0 which in turn will determine the minimum radius of the base circle. From Fig.4,

$$c_{\sigma} = \frac{1}{2} \overline{OA}$$
, $e_{\sigma} = \frac{1}{2} \overline{OB}$ (18a)

$$\overline{\mathbf{DA}} = \frac{\mathbf{O}}{\mathbf{CA}} - \mathbf{O}_{\mathbf{A}} - \mathbf{O}_{\mathbf{A}}$$
(16b)

and

$$\mathbf{0}\mathbf{B} = \mathbf{0}_{\mathbf{0}}^{*} - \mathbf{0}_{\mathbf{0}}^{*} \mathbf{u} \mathbf{u}_{\mathbf{N}}$$
(18c)

The can profile is new resultly obtained by combining eqs(Gasb) to yield

which produces a value of p for each siven value of ψ_1 lies it produces the function $z_{-1}(z)$. In order to obtain the function $z_{-1}(v)$, eqs() thy are again combined to produce

$$\theta = \tan^{-1}\left[\frac{\pi(\xi)}{2}\right] + \psi \qquad (156)$$

which thus produces a function 6-2(d).

Regarding now ψ as a parapeter between functions $\rho(\psi)$ and $\psi(\psi)$, the function $\gamma \circ \rho(\tilde{\psi})$ is readily obtained. The foregoing procedure clearly produces a discrete set of a pairs $\{(\phi_{i}, \phi_{i})\}$, which then have to be interpolated in order to produce a continuous function $\rho \circ \gamma(\psi)$. that would allow both the plotting of the can profile and the punching of the tape guiding a SC-machine tool used for its automatic production. Out of the different procedures to interpolate the obtained not, the one that is proposed here is based upon periodic parametric mplines, which have proved [4,16,pp.27-28] to have the following advantates: i) they require a low number (spall) value of all

- of sample values $\{(o_{1}^{-}, \theta_{2}^{-})\}_{j=1}^{m}$ to produce a given imposed accuracy
- ii) their parameters are readily computed by solving a system of linear equations that is symmetric, tridingonal and diagonally deminant

iii) as a consequence of if), the arising system of equations is well conditioned. . .

A subprogram cas written that, for given calues of $\alpha_{\rm H}$, σ_0 and σ_0^* , produces the optimising calues, e_0 and c_0^* given by eqs($|f_0\rangle$) and for a given value of \mathbf{n}_i , produces the set $((e_1^*, q_1^*))_{I_i}^*$, which is used next to interpolate the cam profile with periodic parametrice splines. The subprogram produces also the interpolated values of e_1^*, e_n^* and \mathbf{r}_i representing the unit tendent and normal vectors to the cam profile and its radius of curvature, respectively. Now, specifying a value a of the radius of the roller, the profile of the cogresponding cam can be obtained from

$$r_{\rm p} = r_{\rm h} = 4e_{\rm H}$$
(20)

where $r_{\rm f}$ and $r_{\rm h}$ are the position vectors of points on the cam profile for the coller and for the knife-edge followers, respectively. This way, knowing the cam profile for the knife-edge follower, as well as the radius d of the coller, eq(20) allows to obtain the ten profile for the toller follower. This would complete the optimal synthesis of the cam profile Soupht, if no e-ditional constraints were to be tepsed. There are, lowever, two more iteras that need be taken into account, namely 1) the maximum allowable value of the radius of the roller, 4, to avoid the phoneemon known as productions, and 11)the tink of allowable value of the radius of curvature of the tam profile for the roller follower, rer, which is necessary to specify in order to avoid too large values of the contact stress.

Let κ_k and κ_p be the curvature of the can profile for the knite-edge islawer and for the roller follower, respectively. Since the emit tangent vectors to both profiles are identical, one readily obtains κ_k

$$\kappa_{\rm E} = \frac{\kappa_{\rm E}}{1 - a\kappa_{\rm K}^2} \tag{21}$$

If the denominator of the right-hand side of eq(21) vanishes, the curvature of the ruller-follawer comprofile, $s_{\rm P}$, will become infinity. This means that, at values of ϕ where that denominator worthes, the sold comprofile has a cusp, which defect is known as analyzeathing. Horeover, both the petch curve and the roller-follower can profile should have, at corresponding points, ise, at the same value of ϕ , curvatures with the same sign. This is thus attained if and only if.

Relation (22) holds, in turn, it

$$t < (r_{ck})$$
 (22b)

or, equivalently, if a is specified as a given function f of (rek) min. i.e.

The minimum value of $r_{\rm ck}$ is now computed. To this end, $r_{\rm ck}^{\prime}$ is zeroed and the values of $r_{\rm ck}$ at stationary points are computed. It can be readily shown that $r_{\rm ck}$ attains stationary values at the dwell phases. These, however, are not, in general, global extrema, for which reason the global minimum is sought both in the rise and in the return phases. Both $r_{\rm ck}$ and $r_{\rm ck}^{\prime}(\psi)$ are given by [19]

$$\mathbf{r}_{ck} = \frac{\mathbf{N}(\mathbf{r})}{\mathbf{D}(\mathbf{\psi})} \tag{23}$$

and

$$\mathbf{r}_{ck}' = \frac{1}{D(\psi)} \{ N^{*}(\psi) - \mathbf{r}_{ck} D^{*}(\psi) \}$$
(24)

with

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$$\begin{split} &\mathsf{N}(\psi) = \left(\mathbf{s}^2 + (\mathbf{s}^1 - \mathbf{c})^2\right)^{\frac{1}{2}} & (25a) \\ &\mathsf{D}(\psi) = \mathbf{s}^2 + (\mathbf{s}^1 - \mathbf{c})(2\mathbf{s}^1 - \mathbf{c}) + \mathbf{s}^{-1} & (25b) \\ &\mathsf{N}^1(\psi) = 3\left[\mathbf{s}\mathbf{s}^1 + \mathbf{s}^{-1}(\mathbf{s}^1 - \mathbf{c})\right]\left[\mathbf{s}^2 + (\mathbf{s}^1 - \mathbf{c})^2\right]^{\frac{1}{2}} & (25c) \\ &\mathsf{D}^1(\psi) + \mathbf{s}(2\mathbf{s}^1 - \mathbf{s}^{-1})^2 + 3\mathbf{s}^{-1}(\mathbf{s}^1 - \mathbf{c}) & (25d) \end{split}$$

The zeros of r_{ck}^{c} can be readily computed with the aid of Subroutine ZEROIN [20] or any other efficient subprogram intended to find the zeros of a numlinear function of a real argument. The global minimum of r_{ck} is then substituted into eq(22c), thus producing the desired value of a, thereby completing the synthesis of the proposed mechanism. The suftware realising the described synthesis produces values of technical interest concerning the can profile, as well. These are: area, location of its centroid, principal moments of inertia. All these values are produced in order to ease the static and dynamic analysis of the overall mechanism.

At this point it is worth contioning the impossibility of lines 1 and 2, Fig.4. of passing

...

through the origin, which in turn, would lead to $\tau=0$. In fact, this would imply

However, from eq(15a) and taking intraccount the populative definitioners of G Juring the rise phase, one has

Nithout tops of ponerality, use can assure that f_{1} is a supporting point, f_{1} , but the spline curve. From the relationships littles the sets $\{0,1\}^{n}_{1}, \{0,1\}^{n}_{1}$ and $\{0,1\}^{n}_{1}$ for cubic splines $\{18,1\}^{n}_{1}$, 27-25, $e_{1}(27)$ would loss to

which is impossible to happen, given the way the set $\{\tau_{ij}^{ij}\}_{j=1}^{ij}$ can avecified in eqs(Lakb).

In the case of a case with a radial follower, e=0, and cause station stationary values either when s'=0 or size when s'=+us". The tirst condition leads to a minimum value of the pressure way give, whereas the second one leads to the maximum. Recalling that n=e+0, $s'=0^{\circ}$ and $s''=0^{\circ}$, one has

$$a^{+2} = (a+e)a^{\mu}$$
 or $\frac{a^{+}}{a^{+}e} = \frac{a^{\mu}}{a^{+}}$ (29)

Substitution of eq(14x) into $c_0/(29)$ and the fact that 0/20 during the rise phase lead to $p_0'=tand_H(p_0+c)$ (30a)

where ρ_0 , σ_0^{\dagger} and σ_0^{\dagger} are defined as before. From eq(30a) one obtains the radius of the base circle as σ^{\dagger} -0 taxa.

$$c_{\theta} = \frac{c_{\theta} c_{\theta} c_{\theta} c_{\theta} c_{\theta}}{c_{\theta} c_{\theta} c_{\theta}}$$
(30b)

4. EXAMPLE

Obtain the cam profile of the translating roller-follower cam rechamism producing the displacement programme given in Table 1. The maxicum pressure angle is limited to 30°, and f=0.50. The subprogram produced the cam profile appearing in Fig.7, with a value of a of 15.00.7.

	Angle of	Displacement
Phase	rotation(V)	(046)
0_{-} (d. (11)		0
A ⁴ (rise)	72	+50.0
D_ (dw211))	144	O D
R ² (return)	103*	

Table 1 Follower-displacement programme

The geometric properties of the cum profile produced by the subprogram are:

Area=19530.3Amm² Centroid=(-10.89,-19.53). E1 and E2 are the principal akes of inertia at the centroid, the corresponding moments of inertia being $I_2=40192170.0mm^2$ and $I_2=30070254.0mm^2$.

5. CONCLUSIONS

An automatic procedure, implemented with the aid of several computer subprograms, was developed. This procedure allows the digital and graphical production of the minimum-size cam profile that generates a given displacement programme for a ruller follower while observing bounds on the pressure angle. Radial and off set followers were rensidered, herein. The software presented here yields, additionaly, the radius of the roller as a fraction of the minimum value of the radius of curvature of the pitch curve, in order to avoid undercutting. Noncover, it produces relevant geometric properties of the profile, such as its area, its controld location, the orientation of its principal area and values of its principal comments of inettia. The software is user oriented and requires no deep knowledge of the algorithm described herein. It is a part of a wider program system intended for the synthesis of cam profiles of various types of followers, out of which the one corresponding to flatface followers was presented proviously.

6. ACKNOW MAGEMENTS

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Fig.1 Periodic function generating the followerdisplacement programme



Fig.2 Geometry of a can mechanism with a translating knife-edge follower



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Fig.4 Plots of eqs(15a616)



Fig.6 1 of the rise phase vs. pressure angle









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OPTIMAL SYNTHESIS OF OSCILLATING ROLLER-FOLLOWER CAN MECHANISMS WITH PRESCRIBED FUNCTIONAL CONSTRAINTS

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ABS:RACT

The synthesis of the can profile producing a given displacement program of its oscillating roller follower, while enclosing a minimum area and having a prescribed maximum pressure angle, is presented. The displacement program of the follower is synthesized, in turn, using cubic periodic splines. While the mathematical programming approach is followed throughout, it is shown that the problem can be solved using a simple root-finding routine applied to a nonlinear equation in one unknown. The procedure is illustrated with a fully-solved example showing its applicability to the automatic demign of cam-follower mechanisms.

INTRODUCTION

While the literature on can mechanisms has widespread since the 1970's, incorporating computer-oriented methods of analysis, synthesis and manufacture [1-17]. can mechanisms with oscillating followers have received very little attention, except for [2] and some other references not fully devoted to this type of mechanisms . This paper follows two provious ones [18.19] concerning the optimal synthesis of cum mechanisms with translating flat-face and coller followers. The approach introduced in those papers is now applied to the synthesis of disk cans with oscillating coller followers producing the desired angular motion of the follower, while enclosing a minimum area and having a prescribed maximum pressure angle. Departing from the usual practice of prescribing the follower motion via harmonic, cycloidal or polynomial functions, all of which contain a limited number of free parameturs whose values are chosen so as to match the different motion phases (lower dwell, rise, upper dwall and return), the authors use cubic periodic splines. To this end, the concept of function and curve synthesis introduced in [20,21] is resorted to. This way, a computer program interacting with the user enables the latter to obtain a smooth follower motion with

e desired phases. A complete turn of the cam plate is but divided into the four phases having angular lengths by, for 2+1,2,3,4. The user can supply these lengths wither in degrees or as percentages of the total turn. In order to offer the user a visual verification of the follower displacement program, the software graphical capabilities provide a display either on a CKT or on a plotter. Once the follower-displacement program is symthesized, the method presented here proceeds to determine the can disk position at which the pressure angle attains its maximum absolute value. This is done by finding the roots of a nonlinear equation in one single unknown, namely the variable defining the can disk ;osition. This value then provides the geometric parameters of the can mechanism. All in all, a visual verification of intermediate results is possible, but the software is so designed as to enable a totally automatic mode of operation.

SYNTHESIS OF THE FOLLOWER-DISPLACEMENT PROGRAM

The angular displacements of the cam plate and the follower are denoted by ψ and ψ , respectively. Moreover, ϕ is assumed to be the sum of a constant c, as yet to be determined, and a positive definite function $u(\psi)$, showe minimum value is 0, its maximum value being A, the omplitude of the follower oscillation, i.e.

$$\Phi(\psi) = c + \sigma(\psi) \tag{1}$$

Function $O(\psi)$ is synthesized using section PQ of the cubic periodic spline $\tau(\psi)$ shown in Fig.1. This is



synthesized, in turn, by specifying a hermonic distribution of $T^{*}(\psi)$ at a set of equally spaced H points in the incerval [0,2m]. (The relationship between $T^{*}(\psi) = \tau_{1}^{*}$ nd t(W2)=t2; for 2-1,2,...,# is linear [22]. Hence, as unknown ordinates of the supporting points of the spline can be obtained as the solution to a linear aystem, of equations, as shown in [20621]. Horeover, in order to obtain a vanishing slope of the resulting epline, at both P and Q. 1(\$) is prescribed to be add with respect to U-W and even with respect to U=#/2, 3#/2. The introduction of the said symmetries results in a smaller number of independent ordinates T2, in fact, only (0+3)/4. The resulting system of equations is sp-sociated to a fixff matrix that is symmetric, positive definite and tridiagonal, all properties of which render it well conditioned and very simple to handle in obtaining the Joknown ordinates. A proper scaling of the section PQ of Fig 1, plus a rigid body translation permit the obtention of the rise phase, R_1 , of function $\sigma(\psi)$, of Fig 2. Finally, the return phase, R_2 , of $\sigma(\psi)$ is synthemized by first reflecting section PQ of $\tau(\psi)$ with .respect to $\psi = \pi$, then scaling it and shifting it exactly as in synthesizing $R_{1^{-1}}$ Constant c of function $\phi(\gamma)$ is next retermined so as to produce a prescribed maximum pressure angle by, while making the area of the cam disk a minisum. This procedure is outlined in next Section.



'Fig.2 Follower displacement program

SYNTHESIS OF THE CAM PROFILE

The layout of the mechanism that is being synthesized is shown in Fig 3. From that figure, the following relations are readily derived:



Fig 3 Can mechanism with oscillating roller-follower

In order to obtain a good torque transminsion, jube pressure angle 0, shown in Fig 3, 1s bounded as

However, determining the exact position at which |a| attains its maximum value is not straightforward, given the lack of smoothness of the absolute-value function at the origin. Hence, the extreme of an even function of a that be smooth enough are sought. As already done in [1819], the values ψ_0 of Ψ , at which cosa sttains its minima are now found. From Fig 3, cosa is obtained from the inner product of the vector connecting Ψ with T and the unit tangent to the pirefacence shown dotted in that figure. This is, in turn, the trajectory traced by the center of the toller on the cam elsk. Thus,

$$\frac{c_{0+0} - \frac{s_{100}}{(1+\beta^2)^2 - 2\beta(1+\beta^2)\cos^2\beta^2}}{(5)}$$

(6)

with

At this moment it is pointed out that the radical of the right-hand side of eq(5) is positive definite. In fact it equals the squared length of a triangle having sides 1 and $B(1+\phi')$, both making an angle ϕ .

Zeroing of decoso/do leads to

$$CB^2 + DB + E = 0 \tag{7a}$$

where

$$C = \phi^{\dagger} \{1 + \phi^{\dagger}\}^{2} \cos \phi - \phi^{\prime\prime} \{1 + \phi^{\dagger}\}^{2} \sin \phi$$
 (7b)

E-¢'cos¢ (7d)

Now, $\{\alpha\}$ estains its maxima at values of ϕ where cost attains its minima. Let

 ψ_0 being the particular value of ψ as which commu-

Fig. 4 shows a mechanism configuration at which the follower is at its lower dwell. Bence,

$$\cos c = \beta$$
 (9)
 $\sin c = \frac{p_0}{a} = (1-\beta^2)^{\frac{1}{2}}$ (10)



Fig 4 Follower at its lower dual)

$$FB^{*} + 2CB^{*} + H^{2} = 0$$
 (11a)

$$\mathbf{x}_{1}^{*} = \mathbf{x}_{0}^{*} - \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}_{0}^{*} + \mathbf{x}$$

$$K_2 = (1+0^{\circ}) m^{\circ} = c_0$$
 (111)

 $v_{0} - \sigma(\phi_{0}), \sigma_{0}^{*} = \sigma_{0}^{*}(\phi_{0}), c_{0} = \cos \sigma_{0}, s_{0} = \sin \sigma_{0}$

From eq(11s) one obtains β as a function of ψ_0 , which then is substituted into eq(7a) together with eqs(1,8,9510), thus producing a nonlinear equation, $f(\psi_0)=0$, in one single unknown, ψ_0 .



Fig. 5 Plot of function f(\$)

Fig 5 shows a plot of $f(\psi)$ for the conditions given in the Example. The value of ψ at which $f(\psi)$ vanishes, ψ_0 , can be obtained either visually from that plot or iteratively, using a suitable algorithm. As previously done [18619], subrouting ZEROIM [23] is used within the program written for the synthesis of this type of follower. Once ψ_0 is determined, σ_0 and hence σ_0^* and σ_0^* can be readily determined from function $\sigma(\phi)$, of fig 2. With o_0 and o_0 already determined, parameter 8 is obtained from eq(11a). Constant c can thus be computed from eq(9). Function $\phi(\psi)$ is therefore totally determined. Substituting these values into eqs(261), for a discrete set of values of ψ_{i} (ψ_{i}); a set of pairs $\{\rho_{i}/a, 0_{i}\}_{i}^{n}$ can be readily computed. These determine in points of the pirch curve, which are now used as the supporting points of a cubic periodic parametric spline used to trace the normalized sold curve, i.e. for a-1. The actual size of this curve is then obtained once pa-Tameters a and r, shown in Fig 3, are defined. The radius of the coller for a unit length a, r/a, can be computed from the pitch curve for a unit length a, by prescribing r/s to be a given fraction q of the minimum radius of curvature of the sold pitch curve. In order to -void undercutting. The said fraction can be determined,

turn, from contact stress considerations, as pointed .4 in [2]. The radius of curvature of the pirch curve, rp, and the pressure angle, 0, are given by [2]

$$\frac{e \sin^2 \phi}{\left(1-\phi^1\right) \cos^2 \alpha \left(\frac{\sin \phi}{\cos \phi \phi} + \frac{w'}{1-\phi''} \sin \left(0+\phi\right) \frac{\phi^{(1-\phi'')}}{\left(1-\phi''\right)^2} \cos \left(0+\phi\right)}\right)}$$
(12)

$$= \tan \frac{-1}{\left[\cot \phi - \frac{\beta(1-\phi^*)}{\sin 2} \right]}$$
(13)

The softwarm realizing the method presented here produces $r_p(\psi)$ both numerically and graphically. Moreover, it yields the extrema of this function with the aid of subroatine MIN [2]]. Fig 6 shows a plot of r_p/a for the conditions given for the Example. Let r_m be the minimum value of r_p/a , p_p , p_c and n denoting the position vactors of corresponding points on the pitch curve and on the cam profile, as well as the unit normal vector of both curves at corresponding points, respectively. Thu cam profile is thus synthesized from the relation [19]:

$$\mathbf{p}_{\mathbf{p}} = \mathbf{p}_{\mathbf{p}} = \mathbf{rn} \tag{14}$$

The can profile obtained from eq(14) corresponds to a unit value of parameter a. This can be chosen, in turn, from considerations of space availability together with maximum allowable value of the contact stress. Used this parameter is determined, the actual cam profile is determined by scaling the foregoing normalized parameters, which is done by a simple multiplication.

DIAMPLE

a

Synthesize a can follower mechanism for an owcillating roller follower, which will produce the follower displacement program appearing in Table 1. The pressure angle is to attain a maximum absolute value of J0°, the amplitude of the follower oscillations being prencribed as 45°

Phase	Angle of rotation (ψ)	Displacement
D ₁ (dwell)	36*	0"
X (rise)	72*	+45"
D ₂ (dwell)	144*	0"
R ² (return)	108*	-45"

Table 1 Prescribed angular-displacement program



/here

wîth

The synthesis was executed with the software developed for the implementation of the method presented here. The values obtained, for a unit value of a, were ψ = 82.61⁴

Prescribing q=0.75 the radius of the roller was thus set as

r/a = 0.285

Function f(0), and the radius of curvature of the pitch curve are plotted in Figs 3 and 6. respectively. The synthesized cam profile is shown in Fig 7. Finally, the software produced the following geometric parameters: Arus of the cam disk = 2.0963

Centroid coordinates: x = -0.1877, y = -0.0805 Principal moments of

inertia of the centroid: $I_1 = 0.4544$, $I_2 = 0.3596$

The principal axes of inertia, E1, E2, corresponding to I₁ and I₂, are shown in Fig 7.



CONCLUSIONS

The method presented here implicitly produces a minimum-fize cam disk for an oscillating toller-follower. moving according to a prescribed angular-displacement program. In fact, by imposing the condition that the maximum absolute value of the pressure angle attains a given value a, the procedure produces the minimum-size can disk. The use of cubic periodic splines, for both the synthesis of the angular-displacement program of the follower and the synthesis of the pitch curve, and hence for that of the comprotile, allows a relatively simple computation of the geometrical parameters of the entire can mechanism. The software implementing the foregoing method gives the designer the freedom to choose the redius of the roller by allowing him/her to specify it as a fraction, q. freely chosen. Moreover, the designer

can determine freely the overall size of the mechanips by properly choosing parameter a. He/she can do this considering space availability and maximum contact stress. Since the paper is concerned with the pure geometric synthesis of the Mechanism, such considerations were left aside here. The noftware can be integrated. however, to a more general CAB program * enabling the designer to couple it to a Fill package allowing him/her to consider such effects as contact stress, stress concontrations and failure criteria. Finally, the software implementing this method produces geometric paramyters such as area, controld location and principal coments of inertia, that are necessary for a gratic and dynamic analysis of the mechanism.

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DIVISION DE EDUCACION CONTINUA FACULTAD DE INGENIERIA U.N.A.M.

DISENO CINEMATICO DE MAQUINARIA

DISENO AUTOMATIZADO DE MECANISMOS DE LEVA DE DISCO CON SEGUIDOR TRASLACIONAL DE CARA PLANA

DR. JORGE ANGELES ALVAREZ DR. CARLOS S. LOPEZ CAJUN

JUNIO, 1984.

OFFERD AUTOMATIZADO DE RECAMISHOS DE LEVA DE DISCO CON SECULIDOR TRASLACIONAL DE CARA PLANA

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Resumen

Se presenta la sintetis automótica de meanismos de leva con seguidor traslacionol de cara plana. La curva de dettlazemiento del seguidor se sintetiza mediante curvas spline presidicas, las cuales satisfacen condiciones preseritas de aceleration. En base a lo anterior se determina el radio óptimo del circulo base para un descentratiento máximo dado. Finalmente, el perfil de la levi se obtiene codjante curvas spline paramétriras periódicis, que interpolan los puntos generados por las ecuaciones de sintesis.

Abstract

The computer-aided synthesis of disk can mechanisms with translational flat face follower is presented. The follower displacement program is synthesised by periodic spline curves, which satisfy prescribed esceleration conditions.Bused on the toregoind, the optimum radius of the base circle is determined for a given maximum contact point recentricity. Finally, the cam profile is obtained by fitting periodic parametric spline curves through the points generated by the synthesis equations.

Introducción

Con el adventatento de las máguinas horramienta de contro) munărico, la manufactura de levas es mis confiable y practa [1]. For otra parte, durante las últimas décadas se hun realizado diversos estudios sobre el análisis y la sintesis de mecanismos de leve [2], muchos de los cuales involucran el uso de la computadora. En particular, sobre diseño óptimo de'levas, pueden citorse los tracajos de Chicurel [J], Hischke [4] y Angeles y Arteuga [5], entre otros. En usto trabajo so presenta un caso particular de síntesis de levas: pero, a diferencia de la práctica tradicional, que se base en el uso de un número limitado de funciones que contienen a su vez un número Finitedo de parámetros independientes, aquí se carestra el uso de curvas spline periódicas sintetizadas en terma tal que satisfacen condiciones de acoleración prescritat y que mirantizan continuidad on la volocidad del segundar [0,7]. La intro- ducción de functenes spline pomite conter con un número artitrario de parámitros independientes (lor coeficientes de la soline), que permiten satisticer un nimero igualmente arbitrario de condictimes soore al desplayamento del seculdor. las curvas spline han sido utilizadas por Kó Sánches y García de Jalón [6] para obtenir los programme del segundar. Sin enbando, un osto tryba-Jo. Vichas culvas se encuentran almae madas, en un banes de datos y, mediante escalusionica (y/o reflexión), pueden generarse los programas de desplazamiento de seguidar requeridos, de los evales se pueden obtener los puntos de velecidad máxica y por lo tanto, el radio óntimo del circulo base para un descentraciento táxico permitido. Por último, consciendo el radio del circulo base el perfil de la leva se obtieno al interpolar los puntos generados por las equeciones de sintenis, cadiante curvas spline pararitericas períodicas. Las funciones solino beriódices usadas atus para representan el desplazamiento del seguidor son de la forma

 $s(t) = a_k (n + n_k)^n + b_k (y + y_k)^2 + c_k (y + t_k) + d_k$ math $y_k = y_{k-1} + 1$ donds $a_k + b_k + c_k (1 + 1, ..., n)$ so obtions do las
condiciones de periodicidad $s(0) = s(2^+) t^*(0) = s^*$ (2*), $s^*(0) = s^*(2^+)$ y de continuidad on + 1, (2, ..., en tanto en $s(1^+)$ coro en $s^*(1^+)$ y on $s^*(1^+)$, para
satisfacer valores presentos de $t^*(0)$ on t_1 , $y_2, ..., y_n$. (Los dotalles pueden versen en (2^+)).

Definición dol Problema

Dados los intervalos $t_1, t_2, t_3 y t_4$, est como la eltvación del seguidar, h, que definen el programa de exsplazaciento de éste, ver fig 1, estenar el perfil de la leva de área minimo, que genere est programa y que tença un descentramiento máximo dado del punto de contacto, PA, de la fig 2.

Sintesis del programa de desplazamiento del seguidor

Se requiere sinterizar la curva víblivs, e de la Fig. 1. comerciando las fasos de reposo Ry y Ro. de forma tal que E y B sman tangentos a P1 y H2 en +1, 12, 03. y re. Esto es para curantizar que la volcoidad del seguidar sea continua en los puntos de comerción, Adresis, dado que la ineleración del seguidor es una tunción lineal de s" (3) y con e) fin de unentizar la continuidad de diche accleración e pita, 15 924, s"(*) debe anu-larsu en estos puntos. Las aún, se desea que la aceleration cambie successence entrony ply y en-tre %3 y t4. Lo anterior to puede lograr [7] rediante el uscalamiento adecuado del tramo PQ de la cuiva spline perticica de la Fig. 3, la cual se eleventra almacenada en un banco de datos. La statutis de la curva 2. definida catheris yba. se logia mediante la reflexión y esexiamiento del Sista Cramo PQ,

Sintesis del porfi) de la leva-

Considerando la Fin. 2, sean N y L times fijas al marco del mocanismo y a la leva, respectivacente. El finjulo le mide la rotación de la leva con respecto al marco, en tanto que s(c) puede considerarse como la suma de una constante 2, el radio del circulo base, más una función positiva definida a(+), cuyo valor mínimo es coro, siendo su valor máximo igual a h.

De la geometría de la Fig..2, se tiene:

Igualando las velocidadas en la dirección vertical del punto A, se tiene:

chadu v_{AL} es la componente vertical de la velocidad de A como punto de la leva y v_{AS} es la velocidad de A como punto del seguidor; por lo tento:

¢

De (1) y (3b) se tiene:

$$s^{2}(*) = s^{2}(*) + s^{*2}(*)$$
 (4)

Combinando (4) y (5) se obtiene el perfil de la leva dado por p = p(1),

Determinación del radio óptimo del circulo base

On the geometria de la Fig. 2, se tiene:

$$\pi X = a \cos(e + \phi)$$
 (6)

ustituyendo (6) cn (3a) y despejando 87:

 $BA = \frac{s}{2} = \frac{ds}{dy}$ (7)

conde.

sierdo Emár el máximo descentramiento (por unidad de longitud de radio del circulo base) permitido, se tiene:

$$\frac{\pi \pi}{y_{s}} = \frac{1}{2} \ln \pi x$$
 (9)

$$s^{*}(s) \leq c_{w\delta x} c^{*}$$
 (10)

de nonde puelle obtenerse el radio optimo del circuto base

$$c_{\text{det}} = \frac{H5 \times I_{\text{S}} + (\nu) I}{5 \mu \delta x}$$
(11)

Las velocidades máximus correspondenta los puntos de inflexión de las curvas Ely B, esta esta esta puntos donde la aceleración se anula y que por construcción de las curvas correspondenta los valores $t = (s_1 + s_2)/2$ y $s = (s_3 + s_4)/2$

Descripción del algoritro

- Lee los intervalos t1,t2,t3,t4, la elevación h y el máximo descentraciento permitidatera
- Genera las curvas de elevación y descenso del seguidor cediante escalemiento adocuado. Grafica está curva
- 3. Obtione Copt

t

- 4. Obtiene θ(e) y ρ (*)
- S. Genera p (*). Lo grafica reciante curvas splina paramétricas periòdicas

Ejempio

Obténgase el carfil de la leva que produzca el movipiento del seguidor mostrado en la Tabla 1 para un descontramiento máxico de 50%, y una elevación h=5 unidades de locjitud. El perfil obtenido se muestra en la fig. 4.

<u>Conclusiones</u>

Se mostró un procedimiento para obtener, en forse automítica, tanto el programa de desclazamiento del seculdor como el pertit de la leva de área mínima con un descentramiento máxima prescrito del punto de contecto. El algoritmo utilizado se realizó en un programa de computatora, al cual tiene acceso al usuario en fuma conversacional y propursiona los resultados tento en forma numérica como gráfica, in el último caso, se pueden obtener éstos en partalla o en copia dura mediante el uso de un gráficador. Se utilizaron para este fin las instalaciones del Laboratorio de Cálculo Automáticado para el Olseno de la División de Estudios de Posgrado de la Facultad de Ingeniería, UNAM.

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	TABLA 1					
Intervalo	Tipo de movimiento					
t, = 36*	reposo (R ₁)				
t, = 108*	elevación (E)				
t - 144*	reposo (R ₂)				
t. 72	descenso (1)	1.				
	Intervalo $t_1 = 36^{\circ}$ $t_2 = 108^{\circ}$ $t_3 = 144^{\circ}$ $t_4 = 72^{\circ}$	IntervaloTipo de movimie $t_1 = 36^{\circ}$ reposo(R_1) $t_2 = 108^{\circ}$ elevación(E $t_3 = 144^{\circ}$ reposa(R_2) $t_4 = 72^{\circ}$ descenso(B)				



Fig. 1 Programa de desplazamiento del seguidor











Nota final: El perfil de la leva se obtuvo con 37 puntos de apoyo distribuídos de la siguiente manera: 9 puntos de apoyo en cada período de reposo y 10 puntos de apoyo en los períodos de ascenso o descenso.

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